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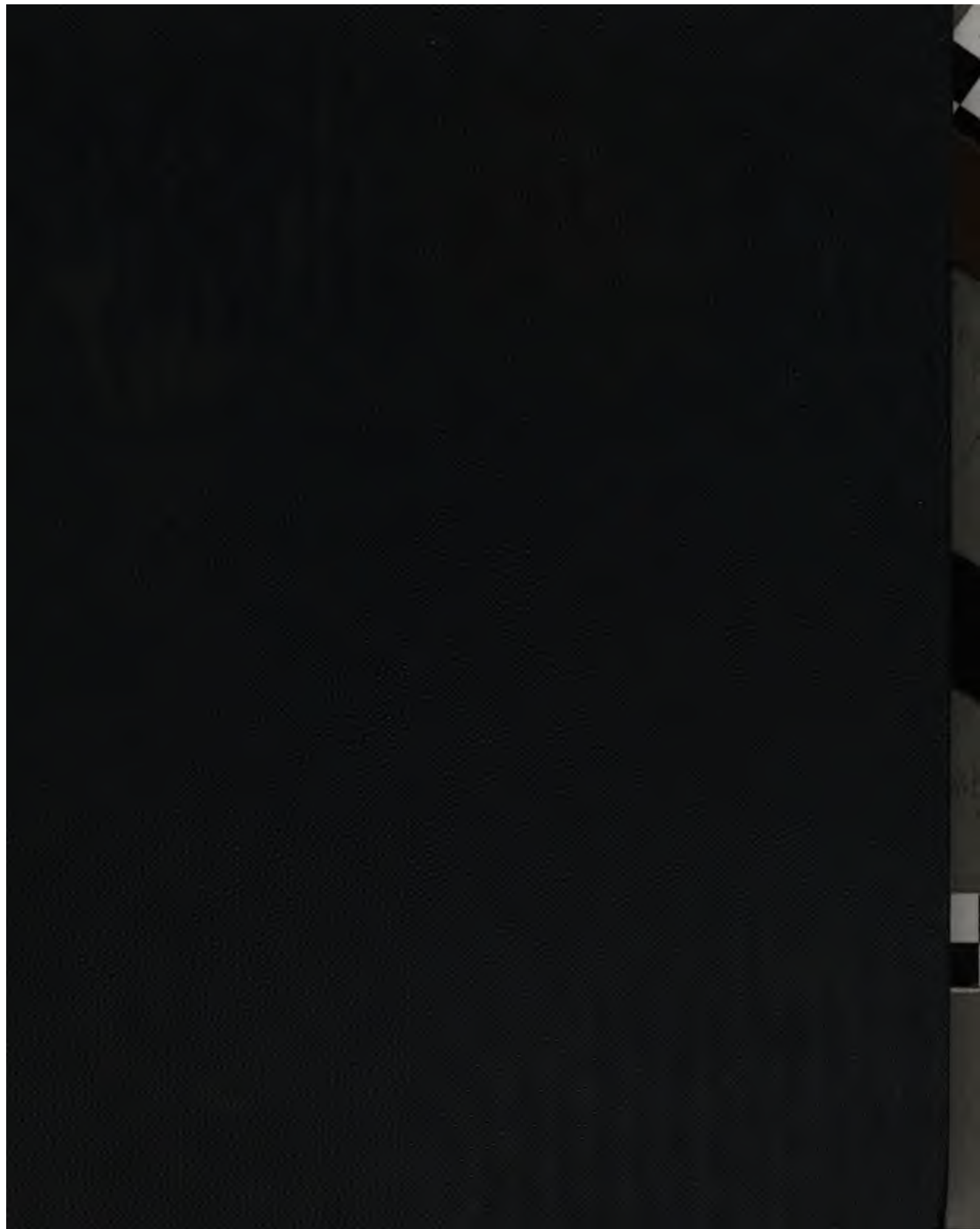
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THE  
ELEMENTARY PARTS  
OF  
*Robert*  
Dr. SMITH'S COMPLEAT SYSTEM OF OPTICKS,  
SELECTED AND ARRANGED  
FOR THE USE OF STUDENTS AT THE UNIVERSITIES:  
TO WHICH ARE ADDED  
IN THE FORM OF NOTES  
SOME  
EXPLANATORY PROPOSITIONS FROM OTHER AUTHORS.

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M. DCC. LXXVIII.

1854. d. 22



*December 1777, St. John's College, Cambridge.*

T H E  
P R E F A C E.

WHEN the Editor of the following Treatise drew the outlines of his plan, he had no Intention either of printing so large a volume or of publishing any. His work was then calculated to facilitate the progress of science in this society alone, by accommodating the students with those parts of Dr. SMITH'S OPTICKS, which are explained to them in the lectures on that Branch of philosophy. The scarcity of that Author's SYSTEM had long been a subject of complaint; and it was at length become impossible to procure a number of copies sufficient for the wants of so large and industrious a society. But the Editor had not proceeded far, before he was diverted from his first design, by being repeatedly informed that the same difficulty prevailed in almost every other College, and that no person could be found, who was willing to hazard a Republication of the whole of Dr. SMITH'S Treatise. He therefore determined, though already oppressed by accumulated business, to extend his plan; rather than suffer the scarcity complained of to be an Impediment to Industry, or become an Excuse for Idleness. To remove the former of these ill consequences, and obviate the latter in the most effectual manner possible, he communicated his design to the Tutors of other Colleges, and added to the stock, he had formerly collected, all the parts of Dr. SMITH'S Treatise, in which he found their lectures differed from his own.

But though the Editor was possessed of the requisite materials, the chief difficulty still remained. For the primitive form of these de-ranged Elements resembled more the chaos of the poets, than a well-connected series of philosophical Reasonings. His next attempt was therefore to reduce the whole into System, by classing its parts and



ranging the classes in such order as should best correspond with the plan of Lectures given by the Tutors in this university, whose exemplary diligence in the cultivation of science entitles them to this and every other assistance. — After extricating these materials from disorder and confusion, in which his success has far exceeded his expectations; he tried also to efface every other mark of deformity, and to give them the appearance of a new creation. But here he met with an insuperable obstacle. In some chapters of the Original the subject is popularly explained, and in others geometrically demonstrated. It was therefore impossible to diffuse over the Whole an uniformity of Style without endless Interpolations and a proportionable Delay, neither of which was consistent with the Reasons, which first induced him to undertake the Work.

Although the following Treatise was intended in its origin, and calculated in its progress, for the sole use of the Academick, the Editor now ventures, on a review of its contents, to recommend it to the use of others, who desire to be instructed in the first principles of Opticks. These principles are demonstrated in it with as much easiness and perspicuity, if not with so much elegance and accuracy, as in other works of the same nature; with so much easiness, indeed, as to be intelligible to every Reader, who is previously furnished with the mere Rudiments of Geometry. And the order in which they are ranged, though peculiarly adapted to the original design for which they were selected, is in the Editor's opinion the most natural, in which the subject can be treated. The following is a slight sketch of this arrangement. In the three first chapters are digested the articles which contain Dr. SMITH's Explanation of the general properties of *Light*, and Sir ISAAC NEWTON's Experiments to prove that Light consists of different *Colours*. In the next chapter is traced the motion of a single *Ray* of Light in its passage through refracting surfaces of a spherical figure, and in the four subsequent chapters the motion of a *Pencil* of Rays, till they unite again, after Reflection or Refraction, to form an *Image* of the object from whence they proceeded. And as we are made to see external objects by means of their Images formed upon the *Retina*, in the two next chapters are explained the whole structure of the human *Eye*, the whole process of *Vision* with the naked Eye, what assistances the sight receives from  
*Telescopes,*

*Telescopes, Microscopes, and Spectacles*, and the method of *constructing* these and other optical Instruments. The eleventh Chapter treats of the *Imperfection* in Telescopes and Microscopes, which is caused by the *aberration* of Rays from their geometrical focus in consequence of the spherical figure of a Lens and the different Refrangibility of different kinds of Light: and the twelfth chapter contains the known *Rules*, which, in constructing Telescopes, it is necessary to observe on account of these aberrations. The Elements of the science being demonstrated, they are applied in the two last chapters to solve the phænomena of *the Rainbow*, and *the annual aberration of the fixed stars*.

The only modern authors, who have treated geometrically the principles of this science, and whose works are still in print, are Mr. EMERSON and Mr. HARRIS. But the philosophical writings of a superior genius are seldom adapted to the capacity of an unassisted Learner; and the demonstrations of the former of these authors are not yet adopted by the Tutors in our universities. The elementary Treatise of the latter is entirely silent upon some of the most important subjects belonging to this science, such as the construction of Telescopes and the phænomena of the Rainbow. Had Mr. Harris lived to finish his Work, it would have precluded the necessity of this and every other publication of the same extent.

It is not intended, in this character of the writings of others, to intimate in the slightest degree that the following Treatise is faultless. It contains many Inaccuracies and even some Errors, of which the Editor was fully sensible before he sent it to the press, but was restrained from correcting them by the dread of Reprehension. The only method of correction was a compleat commentary on the Text, or frequent alterations of it. But, besides that such a commentary would have been as tedious and troublesome as a new Treatise on the subject, there were other objections against it too obvious to be mentioned: and to have erased and corrected the Text of an eminent Writer however judiciously, might have been deemed by some an impertinent presumption and an unjust Treatment of the author.

These Inaccuracies might indeed have been prevented, and an uniformity introduced by composing a new system of the same materials. But as the want of Dr. SMITH'S OPTICAL ELEMENTS was become  
too

too pressing an Inconvenience to allow sufficient time for executing a regular and well-digested plan, the Editor was reduced to the alternative either of garbling the works of that Author or of publishing some crudities of his own. Besides, he had neither health nor leisure to engage in any publication more laborious than the present: He prepared it for the press without the Trouble of copying any part of its contents, except a few propositions, which he has borrowed from other Books of eminence not easily to be procured, and which the Tutors in this university have introduced into their lectures for a more ample explanation of what Dr. SMITH has but slightly touched upon. To have circumscribed his collection of notes within a narrower limit would hardly have been possible; and he judged it more advisable that the Learner should be left to consult the entire works of *modern* authors, than that this volume should be swelled with Extracts from their writings; it being of the greatest assistance to the student, during his noviciate in philosophy, to have the same Truth represented to him in a variety of lights. For the principal notes, which are subjoined to the following pages, the Reader is indebted to Dr. BARROW and DES CARTES.

Such is the Nature and Intention of the following treatise, and such the Editor's apologies for presuming to publish so irregular and incorrect a composition. But if these Reasons be insufficient to defend him against private cavils, he hopes the following considerations will secure him from public Censure:—that it could not be a desire of Fame, which induced him to undertake the mechanical office of an *Editor*; or the Hope of Profit, to be the Instrument of a publication, the Expence of which must be great, and the purchasers few; that he could have no View to his own Improvement in forming a system of principles, which it has been his business for several years successively to explain to others; and lastly, had amusement been his object, that he certainly would have directed his attention to some other province in the intellectual world less frequented by him than the present, less barren and more beautiful. His only motives were public utility and a deep sense of the duty incumbent on every member of these Societies to promote the designs of those venerable Benefactors, whose Endowments they have the honour and happiness to participate.

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T H E

# T H E C O N T E N T S.

CHAP. I.	<b>C</b> ONCERNING <i>Light.</i>	—	—	Page 1
	II. Concerning the <i>Origin and Cause of Colours.</i>			8
	III. Concerning the <i>Causes of Refraction, Reflection, Inflection and Emission of Light, and concerning Transparency, Opacity, and Colours in Bodies.</i>	—	—	16
	IV. Concerning the <i>Refractions of a single Ray of Light in its Passage through a Prism, Globe, or Lens.</i>	—		27
	V. <i>To find the Focus of Rays reflected from any given Surface</i>			32.
	VI. <i>To determine the Place, Magnitude, and Situation of Images formed by reflected Rays.</i>	—	—	36
	VII. <i>To find the Focus of Rays falling almost perpendicularly upon any refracting Surface, Sphere or Lens.</i>	—	—	38
	VIII. <i>To determine the Place, Magnitude, and Situation of Images formed by refracted Rays.</i>	—	—	49
	IX. Concerning the <i>Eye and Manner of Vision.</i>	—		51
	X. Concerning <i>Vision with Glasses.</i>	—	—	62.
	XI. <i>To determine the Aberration of Rays from the geometrical Focus caused by their unequal Refrangibility, and also by the Sphericalness of the Figure of reflecting and refracting Surfaces.</i>	—	—	80
	XII. <i>A refracting or reflecting Telescope being given, whose Aperture and Eye-glass are adjusted by Experience, to determine the Length, Aperture and Eye-glass of another Telescope, through which an Object shall appear as bright and distinct as in the given one, and magnified as much as shall be required.</i>	—	—	88
	XIII. Concerning the <i>Rain-bow.</i>	—	—	98
	XIV. <i>Telescopical Discoveries in the fixt Stars.</i>	—		110

ERRATA.



## E R R A T A.

### Article

- 40. line 6. *for* reflected, *read* refracted.
- 47. l. 8. *for* the refracting, *r.* their refracting.
- 64. In the margin *for* Fig. 56 to 59. *r.* Fig. 56 & 59.
- 70. at the beginning of the 2d part, the word DEFINITIONS is wanting.
- 79. line 7. *for* the Focus *Q*, *r.* the Focus *q*.
- 129. l. 1. *for* *E q*, *r.* *E Q*.
- 140. l. 34. *for* perhaps a hundred miles or more, *r.* some hundreds of miles.
- 172. l. 8. *for* *p l*, *r.* *q l*.
- 181. l. 21. *for* *T c s*, *r.* *T c S*.
- 228. l. 1. *di*le and double microscopes.
- 244. in the 4th column of the Table, *for* 199 and 398, *r.* 200 and 400.
- 258. in the margin, *for* Art. 31. *r.* Newton's Opt. p. 114. 8vo.

### Plate

- 5. Fig. 67. *C* is wanting at the vertex of the parabola.
- 9. Fig. 127. *r* is wanting at the upper extremity of *p q r*.
- 11. Fig. 157. *C* is wanting where *E O* cuts the circle.

T H E

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# O P T I C K S.

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## CHAP. I. CONCERNING LIGHT.

1. **W**HOEVER has considered what a number of properties Light consists and effects of light are exactly similar to the properties of parts. and effects of bodies of a sensible bulk, will find it difficult to conceive that light is any thing else but very small and distinct particles of matter \*: which being incessantly thrown out from shining substances, and every way dispersed by reflection from all others, do impress upon our organs of seeing that peculiar motion, which is requisite to excite in our minds the sensation of light. But for the present purpose it is sufficient to observe that light consists of parts, both successive in the same lines and contemporary in several lines: because in the same place, you may stop that which comes one moment, and let pass that which comes presently after; and at the same time, you may stop it in one place, and let it pass in another. For that part of the light which is stoppt cannot be the same with that which is let pass.

2. The least light or part of light, which may be stoppt alone without the rest of the light, or propagated alone, or do or suffer any thing alone, which the rest of the light doth not or suffereth not, is called a Ray of light †. That rays of light are straight, is evident enough from the shadows of bodies; or from the appearance of light passing through little holes into a dark room full of dust or smoke; or because bodies cannot be seen through the bore of a bended

A ray of light  
what and how  
considered.

\* Newt. Opt. Qu. 29. p. 345. 8<sup>o</sup>. Edit.

† Newton's definition. Opt. p. 1.



pipe; or because they cease to be seen by the interposition of other bodies, as the fixt stars by the interposition of the moon and planets; and the parts of the sun by the interposition of the Moon, Mercury or Venus. Rays of light may therefore be represented by straight lines, not Mathematical but Physical, which are described by the motion of the parts or particles of light: and the point which a ray possesses in falling upon any surface may be considered as a Physical Point.

The manner  
of reflection  
and refraction  
of a ray de-  
scribed.  
Fig. 1.

3. When a ray of light falls obliquely upon a smooth polished surface, it is turned out of its way either by reflection or refraction in the following manner. Imagine the paper upon which this figure is drawn to be perpendicular to the surface of stagnating water, and to cut it in the line  $RS$ , and that a ray of light, coming in the air along the line  $AC$ , falls upon  $RS$  at the point  $C$ . Then supposing the line  $PCQ$  to be perpendicular to the surface of the water, if the ray be reflected, or turned back at  $C$  into the air again, it will describe a straight line  $CB$ , inclined to the perpendicular  $CP$  at an angle  $PCB$  exactly equal to the angle  $PCA$ .

Fig. 2.

But if the ray that came along  $AC$  goes into the water at  $C$ , it will not proceed straight forward, but being refracted or bent at  $C$ , it will describe another straight line  $CE$  inclined to the perpendicular  $CQ$  at a lesser angle  $ECQ$  than the angle  $ACP$ ; and the line  $CE$  will always be so situated, that when any circle, described about the center  $C$ , cuts the line  $CA$  in  $A$  and  $CE$  in  $E$ , the perpendiculars  $AD$  and  $EF$ , drawn from  $A$  and  $E$  to the line  $PQ$ , shall always bear the same proportion to each other; whatever be the magnitude of the angle  $ACP$ . In water the line  $EF$  is always three quarters of  $AD$ .

Angles and  
sines of inci-  
dence and re-  
fraction what.  
Fig. 1, 2.

4. In both these cases the line  $AC$  is called the Incident Ray,  $CB$  the Reflected Ray,  $CE$  the Refracted Ray,  $C$  the point of incidence,  $PCQ$  the perpendicular (at the point) of incidence, the angle  $ACP$  the Angle of Incidence,  $BCP$  the Angle of Reflection,  $ECQ$  the Angle of Refraction; the line  $AD$  the Sine of Incidence, that is, of the angle of incidence; and  $EF$  the Sine of Refraction, that is, of the angle of refraction.

A medium  
what.

5. Empty space, or any transparent body, is called a Medium; and mediums are denser in proportion as they are heavier bulk for bulk; and their power to reflect and refract light is found to be greater in proportion as they are denser, very nearly\*.

Laws of reflection and refraction what.

6. The foregoing properties of Reflection and Refraction being discovered and established by repeated experiments upon light and bodies of all sorts both fluid and solid, without any exception yet

\* Newt. Opt. p. 245. 8°.

known;



known; and being the principal foundation of the whole science of Opticks, are called the Laws of Reflection and Refraction; and are expressed by Sir *Isaac Newton* in the following words.

7. *The angles of reflection and refraction lye in one and the same plane with the angle of incidence*; that is, in the plane drawn through the incident ray and the perpendicular at the point of incidence, as represented in the figures, 1, 2. First law.

8. *The angle of reflection is equal to the angle of incidence*. Second law.

9. Hence it follows that the incident and reflected rays are equally inclined to the reflecting plane; that is, the angles *ACR* and *BCS* are equal; as appears by taking the equal angles *PCA* and *PCB* from the equal angles *PCR* and *PCS*. First consequence.  
Fig. 1.

10. It follows also that when the incident ray is perpendicular to the reflecting surface, it shall be reflected directly back along the same perpendicular; as appears by diminishing the equal angles of incidence and reflection till the rays *AC*, *CB* coincide with the perpendicular *CP*. Second consequence.

11. *If the reflected or refracted ray be returned directly back to the point of incidence, it shall be reflected or refracted into the same line before described by the incident ray*. Third law.

12. *Refraction out of a rarer medium into a denser<sup>a</sup> is made towards the perpendicular*; that is, so that the angle of refraction be less than the angle of incidence. Fourth law.  
<sup>a</sup> Art. 5.

13. *The sine of incidence, AD, is to the sine of refraction, EF, either accurately or very nearly in a given ratio*; that is, supposing any other incident ray *aC* to be refracted into the line *Ce*, and the sines *ad* and *ef* to be drawn perpendicular to *PQ*, the ratio of *ad* to *ef* is the same as the ratio of *AD* to *EF*. It is found by experience, that if the refraction be made out of air into water, the sine of incidence of red light is to the sine of its refraction as 4 to 3: if out of air into glass as 17 to 11, or nearly as 3 to 2. In light of other colours the sines have other proportions, but the difference is so little that it seldom need be considered. Fifth law.  
Fig. 4.

14. Hence it appears by inspection of the figures (2, 3, 4.) that when the angle of incidence *ACP* is increased, the corresponding angle of refraction *ECQ* will also be increased; because the ratio of their sines, *AD*, *EF*, cannot continue the same unless they be both increased. Consequently if two angles of incidence be equal to each other, the angles of refraction will also be equal to each other. On the contrary, when the angle of incidence is diminished, the angle of refraction will also be diminished; insomuch that when one of these angles becomes infinitely small the other also becomes infinitely small. First consequence.

Second consequence.

15. And so it comes to pass that when the incident ray coincides with the perpendicular to the refracting surface, it will proceed straight forward into the other medium without any bending at all.

Third consequence.

16. From which it is reasonable enough to conclude back again, that while the angle of incidence is continually increasing, the refracted ray will be continually more and more bent and diverted from the course of the incident ray produced: I mean if  $AC$  be continued to  $G$ , the arc  $EG$  and the angle  $ECG$  will continually increase \* : especially considering that when the angle of incidence in air becomes very nearly a right one, and consequently the incident ray goes almost parallel to the surface of the water, this ray is as much bent at  $C$  into the line  $CE$  as the 3<sup>d</sup> figure represents. In which

Fig. 2.

$EF$ , the sine of refraction, being always three quarters <sup>a</sup> of  $AD$ , is now three quarters of the radius of the circle. Hence we find † that this angle of refraction,  $ECQ$ , is about  $48\frac{1}{2}$  degrees: and so the angle  $ECS$  (being its complement to 90 degrees) is about  $41\frac{1}{2}$  degrees; which in this case measures the deviation of the ray from its first course along the surface of the water. The deviation at the surface of glass is greater than at the surface of water; the ratio of the sines.

Fig. 3.

<sup>a</sup> Art. 13.

\* PROPOSITION. *The greater the angle of incidence is, the greater will be the angle of deviation.*

Fig. 1.

<sup>a</sup> Art. 8.

1. In the case of reflection, the angle of deviation is that angle which is contained by the incident and reflected rays. But if the angle of incidence  $ACP$  increases, the double of that angle, or  $ACB$ <sup>a</sup>, must increase also.

Fig. 5.

2. In the case of refraction, the angle of deviation is that angle which is contained by the refracted ray and the course of the incident ray produced. Draw  $QBP$  perpendicular to the refracting plane  $EBF$ ; let  $ABG$ ,  $DBH$ , be two incident rays, of which  $AB$  falls more obliquely than  $DB$ ; and let  $Ba$ ,  $Bd$ , be the directions in which they move respectively after refraction: I affirm that the angle  $GBa$  is greater than  $HBd$ .

<sup>b</sup> Euc.III.31.

In the perpendicular  $QB$  produced take any point  $P$ , and upon the diameter  $BP$  describe the semicircle  $BGP$ , cutting  $AB$ ,  $DB$  produced in the points  $G$ ,  $H$ , and  $Ba$ ,  $Bd$  in the points  $a$ ,  $d$ : Join  $PG$ ,  $Pa$ ,  $PH$ ,  $Pd$ . And  $PGB$ ,  $PaB$ ,  $PHB$ ,  $PdB$  being right angles<sup>b</sup>,  $PG$ ,  $PH$  are the right sines of the angles of incidence  $PBG$ ,  $PBH$ , to the radius  $BP$ ; also  $Pa$ ,  $Pd$  are the right sines of the angles of refraction  $PBa$ ,  $PBd$ , to the same radius. Next, make the angle  $GBK$  equal to  $HBd$ , or the arc  $GK$  equal to the arc  $Hd$ ; and join  $Gd$ , cutting  $PK$  in  $X$ : lastly, draw the chords  $Ga$ ,  $Hd$ . Now the angles  $PGd$ ,  $PHd$  being equal<sup>c</sup>, and the angles  $GPK$ ,  $HPd$  being also equal, by construction, the triangles  $GPX$ ,  $HPd$  are similar. Therefore  $PG$  is to  $PX$  as  $PH$  to  $Pd$ : but, by the law of refraction<sup>d</sup>,  $PH$  is to  $Pd$  as  $PG$  to  $Pa$ ; consequently  $PG$  is to  $PX$  as  $PG$  to  $Pa$ , and therefore  $PX$  equals  $Pa$ . But  $PX$  is less than  $PK$ , because the chord  $Gd$  lies wholly within the circle; and therefore  $Pa$  is less than  $PK$ : consequently  $PK$  cuts the angle  $GPa$ , and the arc  $Ga$  is greater than the arc  $GK$ , or  $Hd$ . Wherefore the angle  $GBa$  is greater than  $HBd$ . Q. E. D.

<sup>c</sup> Euc.III.27.

<sup>d</sup> Art. 13.

In the preceding demonstration, the angle of incidence is supposed to be greater than the angle of refraction: But the truth of the proposition may be easily deduced from it, when the ray passes out of a denser medium into a rarer. For if  $aB$ ,  $dB$  are made the incident rays,  $BA$ ,  $BD$  become the refracted rays<sup>d</sup>; and therefore  $aBG$ ,  $dBH$  are still the angles of deviation.

† By a Table of sines.

being



being greater, that is, as 3 to 2, or nearer as 31 to 20. Hence we find that the angle  $ECQ$  is about 40 and  $ECS$  about 50 degrees.

17. The bending and deviation is the same when the ray goes back again along the same lines  $EC, CA$ ; and if an angle of incidence  $eCQ$  be any thing greater than about  $48\frac{1}{2}$  degrees in water, or any thing greater than about 40 in glass, this ray  $eC$  will not be refracted into air, but will be reflected into the line  $Cf$ , making the angle of reflection  $Q Cf$  equal to the angle of incidence  $Q Ce$ .

Refraction  
changed into  
reflection.

18. The truth of these laws and of all the consequences drawn from them may be easily examined in the manner following. Upon a smooth board  $KLMN$ , about a center  $C$  with any opening of the compasses (the larger the better) describe a circle  $PRQS$ ; and having drawn two diameters  $PQ$  and  $RS$  perpendicular to each other, from the point  $P$ , with any opening of the compasses, cut off equal arches  $PA, PB$ , and draw the lines  $CA, CB$ ; then sticking three pins perpendicular to the board at the points  $A, B, C$ , dip the board into water as far as the line  $RS$ ; and holding it perpendicular to the surface of the water, look along the pins  $A, C$ ; and an image of the pin  $B$  will appear in the water in the line  $AC$  produced. Which shews that the ray which came from the pin  $B$  is reflected from the water, at the point  $C$ , along the line  $CA$  to the eye of the spectator. If the pin at  $C$  touches the water, it will disturb the smoothness of its surface; and therefore it is better not to place it in the center, but a little higher in the line  $CA$ . The event will be the same if the reflection be made by any other fluid or solid body, as may be tried by cutting off the lower semicircle, and by placing the diameter,  $RS$ , of the upper semicircle upon the surface of the solid.

Experimental  
proof of these  
Laws of reflection and refraction.  
Fig. 4.

Upon the same board draw the line  $AB$  cutting  $CP$  in  $D$ , and from the lines  $DB$  and  $CS$  cut off  $DH$  and  $CI$ , each equal to three quarters of  $DA$ , and through the points  $H, I$ , draw the line  $HIE$ , cutting the circumference in  $E$ ; and the perpendicular  $EF$  drawn from  $E$  upon  $PQ$  will be equal to  $DH$ , or three quarters of  $DA$ . Then stick another pin at  $E$ , and the board being dipped into water, as before, the pin at  $E$  will appear to the eye to be in the same line with the pins at  $A$  and  $C$ . Which shews that the ray which comes from the pin  $E$  is so refracted at  $C$ , as to advance to the eye along the line  $CA$ ; and therefore when the refraction is made out of water into air,  $EF$  the sine of incidence, is to  $AD$  the sine of refraction, as 3 to 4. If other pins be fixed any where in the line  $CE$ , they will all appear in the line  $AC$  produced: and the whole line  $CE$  will appear in the water as if it were a continuation of  $AC$  straight forward. Which shews that the ray which comes from the pin  $E$ , describes a straight line in the water; and that it is bent at the surface only. On the

con-



contrary, if an opportunity be taken when the Sun is just so high, that the shadow of the pin  $A$  shall coincide with the line  $AC$ , the refracted shadow will coincide with the line  $CE$ . Or whatever be the Sun's height, move the pin  $A$  higher or lower till the shadow falls upon the center  $C$ , and there fix it, suppose at  $a$ ; then sticking the compasses into any point of the refracted shadow, take up the board, and through this point and the center  $C$  draw a line  $Ce$ , cutting the circle in a new point  $e$ ; and the ratio of the new perpendiculars,  $ad$  and  $ef$ , will be the same as before; that is, as 4 to 3, as near as can be measured.

This proof  
applied to  
spherical sur-  
faces.

Fig. 6, 7.

<sup>a</sup> Art. 18.

An object  
what and how  
it radiates.

A focus, pen-  
cil, parallel  
rays what.  
Fig. 10.

19. Lastly it is to be observed, that a ray of light is reflected or refracted at a spherical surface according to the same laws as if it were reflected or refracted at a plane, touching the spherical surface at the point of incidence. Let  $AC$  be a ray of light falling upon any point  $C$  of a spherical surface  $MCN$ , represented by the arc  $MCN$ , whose center is  $O$ ; through the points  $O$  and  $C$  draw the line  $PQ$ , and the line  $RCS$  perpendicular to it, representing a plane surface touching the spherical surface at  $C$ . Now because a ray of light is considered as a physical line, and is refracted or reflected at a physical point<sup>a</sup>, which is common to both surfaces,  $MCN$  and  $RCS$ , it follows that the refracted or reflected ray will take the same course in both cases. And this argument is also confirmed by universal experience.

20. As rays of light are incessantly thrown out and dispersed in all possible directions from every point of a luminous body; so when they illuminate other bodies, on which they fall, they are also incessantly thrown back from every point of these bodies. For the points of opaque bodies so enlightened are visible to the eye at any point of space and in any point of time, as well as the points of the luminous body that enlightened them. The numberless rays which flow from all visible bodies, called objects, may be methodically distributed in this manner. The surface of the object is considered as consisting of physical lines, and these lines as consisting of physical points, and these points are conceived to radiate all manner of ways. It is usual to make use of nothing else for an object but a physical line. For by how much that line is increased or diminished in apparent magnitude or brightness or distinctness, so much the diameter or length of any object, in its place, would be increased or diminished.

21. The point  $Q$  from which rays diverge, or towards which they converge (being made to go back towards the same point though they may never meet at it) is called their focus. And in both cases any parcel of these rays, as  $QBC$ , or  $QBA$ , considered apart from the rest, is called a pencil of rays; and these rays are said to belong to that



that focus, whether it be near at hand or at an immense distance; and in the latter case the rays are called, and considered as, parallel or equidistant from each other; because the difference of their distances at any two given places is insensible.

22. The 8th and 9th figures represent a pencil of rays,  $QC$ , which falling in parallel lines upon a plane polished surface, represented by the line  $ACB$ , are reflected from it into as many other parallel lines,  $Cq$ ; because they are inclined to that plane just as much as the incident rays were inclined to it<sup>a</sup>.

Reflection of a pencil of parallel rays at a plane surface.

<sup>a</sup> Art. 9.

23. In the 11th figure  $QC$  represents a pencil of parallel rays falling obliquely upon a straight line  $ACB$ , or upon a plane surface represented by it, which after refraction are also parallel among themselves; every one being equally bent. Because when the angles of incidence are all equal among themselves the angles of refraction are also equal among themselves<sup>b</sup>. For the same reason if these rays be refracted again at another plane, either parallel or oblique to the former, they will still be parallel among themselves after every refraction. In strictness this is only to be understood of rays of the same colour: as will be explained in the next chapter.

Refraction of a pencil of parallel rays at a plane surface.

<sup>b</sup> Art. 14.

Fig. 11, 12.

24. Let the light which flows from a point  $A$  and passes through a square hole  $bcd$  be received upon a plane,  $BCDE$ , parallel to the plane of the hole; or if you please let the figure  $BD$  be the shadow of the plane  $bd$ ; and when the distance  $AB$  is double of  $Ab$ , the length and breadth of the shadow  $BD$  will each be double the length and breadth of the plane  $bd$ ; and treble, when  $AB$  is treble of  $Ab$ ; and so on: which may be easily examined by the light of a candle placed at  $A$ .

Experiment: To shew that the breadths of a pencil of diverging rays are as their distances from the focus.

Fig. 13.

25. Therefore the surface of the shadow  $BD$ , at the distance  $AB$  double of  $Ab$ , is divisible into four squares, and at a treble distance, into nine squares, severally equal to the square  $bd$ , as represented in the figure. The light then which falls upon the plane  $bd$ , being suffered to pass to a double distance, will be uniformly spread over four times the space, and consequently will be four times thinner in every part of that space, and at a treble distance it will be nine times thinner, and at a quadruple distance sixteen times thinner, than it was at first; and so on according to the increase of the square surfaces  $bcd$ ,  $BCDE$ , &c, or of the square surfaces  $Abfg$ ,  $ABFG$ , &c, built upon the distances  $Ab$ ,  $AB$ , &c. Consequently the quantities of this rarified light received upon a surface of any given size and shape whatever, removed successively to those several distances, will be but one quarter, one ninth, one sixteenth, of the whole quantity received by it at the first distance  $Ab$ . Or in general words the densities and quantities,

Hence the density and quantity of light received upon a given plane are reciprocally as the squares of its distances from the luminous body.

quantities of light, received upon any given plane, are diminished in the same proportion as the squares of the distances of that plane, from the luminous body, are increased: and on the contrary, are increased in the same proportion as those squares are diminished. For the lights of the several points of the body, which severally follow this rule, will compose a light which will still follow the same rule.

## CHAP. II.

## CONCERNING THE ORIGIN AND CAUSE OF COLOURS.

Design.

THE preceding chapter contains such properties as belong to all kinds of light; in this are related some of the experiments by which Sir *Isaac Newton* discovered that light consists of different colours.

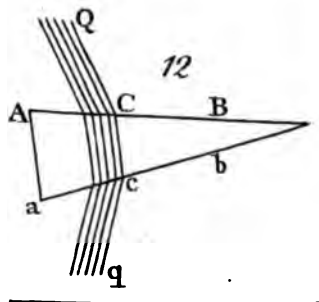
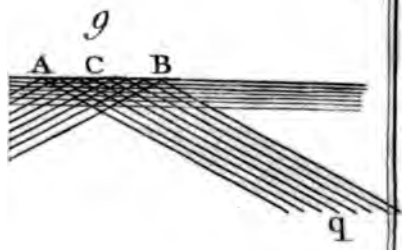
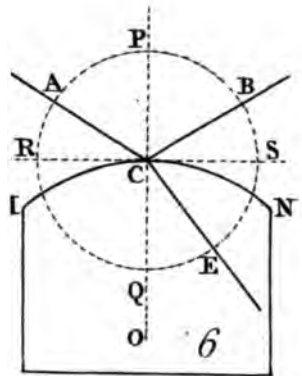
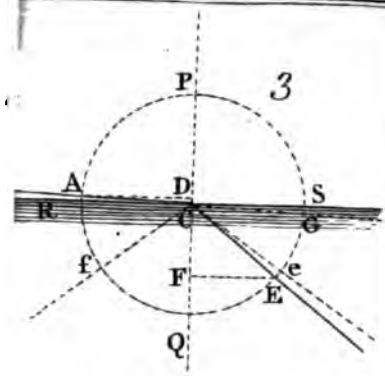
A prism  
what.  
Fig. 14, 15.

26. A glass prism is a body, shaped like a wedge, that has three edges, being bounded with two equal and parallel triangular ends  $ABC$  and  $abc$ , and three plane and well polished sides, which meet in three parallel lines  $Aa$ ,  $Bb$ ,  $Cc$ , running from the three angles of one end to the three angles of the other: and when it is viewed endways it is represented only by a triangle  $ABC$ , as in the 15th figure.

I.  
Experiment.  
A description  
of the sun's  
image made  
by a prism.  
Newt. Opt.  
p. 21.  
Fig. 15.

27. In a very dark chamber at a round hole  $F$ , about one third of an inch broad, made in the shut of a window, I placed a glass prism  $ABC$  whereby the beam of the sun's light  $SF$ , which came in at that hole, might be refracted upwards, toward the opposite wall of the chamber, and there form a coloured image of the sun, represented at  $PT$ . The axis of the prism, (that is the line passing through the middle of the prism, from one end of it to the other end, parallel to the edge of the refracting angle) was in this and the following experiments perpendicular to the incident rays. About this axis I turned the prism slowly, and saw the refracted light on the wall, or coloured image of the sun, first to descend, and then to ascend. Between the descent and ascent when the image seemed stationary, I stopped the prism and fixt it in that posture.

Then I let the refracted light fall perpendicularly upon a sheet of white paper  $MN$ , placed at the opposite wall of the chamber, and observed the figure and dimensions of the solar image,  $PT$ , formed on the paper by that light. This image was oblong and not oval, but terminated by two rectilinear and parallel sides and two semicircular ends. On its sides it was bounded pretty distinctly, but on its







its ends very confusedly and indistinctly, the light there decaying and vanishing by degrees. At the distance of  $18\frac{1}{2}$  feet from the prism the breadth of the image was about  $2\frac{1}{2}$  inches, but its length was about  $10\frac{1}{2}$  inches, and the length of its rectilinear sides about 8 inches; and  $ACB$  the refracting angle of the prism, whereby so great a length was made, was  $64$  degrees. With a less angle the length of the image was less, the breadth remaining the same. It is farther to be observed that the rays went on in straight lines from the prism to the image, and therefore at their going out of the prism had all that inclination to one another from which the length of the image proceeded. This image  $PT$  was coloured, and the more eminent colours lay in this order from the bottom at  $T$  to the top at  $P$ ; red, orange, yellow, green, blue, indigo, violet; together with all their intermediate degrees in a continual succession perpetually varying.

28. Our author concludes from this experiment, and some others to be mentioned hereafter, that the light of the sun consists of a mixture of several sorts of coloured rays, some of which at equal incidences are more refracted than others, and therefore are called more refrangible. The red at  $T$ , being nearest to the place  $T$ , where the rays of the sun would go directly if the prism was taken away, is the least refracted of all the rays; and the orange, yellow, green, blue, indigo and violet are continually more and more refracted, as they are more and more diverted from the course of the direct light. For by mathematical reasoning he has proved, that when the prism is fixt in the posture above mentioned, so that the place of the image shall be the lowest possible, or at the limit between its descent and ascent, the figure of the image ought then to be round like the spot at  $T$ , if all the rays that tended to it were equally refracted. Therefore seeing by experience it is found that this image is not round, but about 5 times longer than broad, it follows that all the rays are not equally refracted.

Hence the  
sun's rays are  
differently re-  
frangible.

For the discovery of this fundamental property of light, which has opened the whole mystery of colours, we see our author was not only beholden to the experiments themselves, which many others had made before him, but also to his skill in geometry; which was absolutely necessary to determine what the figure of the refracted image ought to be upon the old principle of an equal refraction of all the rays: but having thus made the discovery he contrived the following experiment to prove it at sight.

In the middle of two thin boards,  $DE, de$ , I made a round hole in each, at  $G$  and  $g$ , a third part of an inch in diameter; and in the window-shut a much larger hole being made, at  $F$ , to let into my  
B darkened

II.  
Experiment.  
Newt. Opt.  
p. 37.  
Fig. 16.



darkened chamber a large beam of the sun's light, I placed a prism, *ABC*, behind the shut in that beam, to refract it towards the opposite wall; and close behind this prism I fixed one of the boards *DE*, in such manner that the middle of the refracted light might pass through the hole made in it at *G*, and the rest be intercepted by the board. Then at the distance of about 12 feet from the first board I fixed the other board *de*, in such manner that the middle of the refracted light, which came through the hole in the first board, and fell upon the opposite wall, might pass through the hole *g* in this other board *de*, and the rest being intercepted by the board might paint upon it the coloured spectrum of the sun. And close behind this board I fixed another prism *abc* to refract the light which came through the hole *g*. Then I returned speedily to the first prism *ABC* and by turning it slowly to and fro about its axis, I caused the image which fell upon the second board *de* to move up and down upon that board, that all its parts might pass successively through the hole in that board, and fall upon the prism behind it. And in the mean time I noted the places, *M*, *N*, on the opposite wall, to which that light after its refraction in the second prism did pass; and by the difference of the places at *M* and *N*, I found that the light, which being most refracted in the first prism *ABC*, did go to the blue end of the image, was again more refracted by the second prism *abc*, than the light which went to the red end of that image. For when the lower part of the light which fell upon the second board *de*, was cast through the hole *g*, it went to a lower place *M* on the wall; and when the higher part of that light was cast through the same hole *g*, it went to a higher place *N* on the wall; and when any intermediate part of the light was cast through that hole, it went to some place in the wall between *M* and *N*. The unchanged position of the holes in the boards made the incidence of the rays upon the second prism to be the same in all cases. And yet in that common incidence some of the rays were more refracted and others less: and those were more refracted in this prism, which by a greater refraction in the first prism were more turned out of their way; and therefore for their constancy of being more refracted are deservedly called more refrangible.

Our author shews also, by experiments made with a convex glass, that lights (reflected from natural bodies) which differ in colour, differ also in degrees of refrangibility<sup>a</sup>: and that they differ in the same manner as the rays of the sun do.

<sup>a</sup> Newt. Opt.  
p. 16.

Definitions  
Newt. Opt.  
p. 4.

The light whose rays are all alike refrangible I call simple homogeneous and similar, and that whose rays are some more refrangible than



than others I call compound, heterogeneous and dissimilar. The former light I call homogeneous not because I would affirm it so in all respects; but because the rays which agree in refrangibility agree at least in all their other properties which are considered in this chapter.

The colours of homogeneous lights I call primary, homogeneous and simple, and those of heterogeneous lights, heterogeneous and compound. For these are always compounded of homogeneous lights, as will appear in the following articles.

The homogeneous light and rays which appear red, or rather make objects appear so, I call rubrified or red-making; those which make objects appear yellow, green, blue and violet, I call yellow-making, green-making, blue-making, violet-making; and so the rest. And if at any time I speak of light and rays as coloured or endued with colours, I would be understood to speak not philosophically and properly but grossly, and according to such conceptions as vulgar people in seeing all these experiments would be apt to frame. For the rays to speak properly are not coloured. In them there is nothing else than a certain power and disposition to stir up a sensation of this or that colour. For as sound in a bell or musical string or other sounding body, is nothing but a trembling motion, and in the air nothing but that motion propagated from the object, and in the sensorium it is a sense of that motion under the form of sound; so colours in the object are nothing but a disposition to reflect this or that sort of rays more copiously than the rest; in the rays they are nothing but their dispositions to propagate this or that motion into the sensorium; and in the sensorium they are sensations of those motions under the forms of colours.

29. Homogeneous light is refracted regularly without any dilatation splitting or shattering of the rays, and the confused vision of objects seen through refracting bodies by heterogeneous light, arises from the different refrangibility of several sorts of rays. This will appear by the experiments which follow. In the middle of a black paper I made a round hole about a fifth or a sixth part of an inch in diameter. Upon this paper I caused the spectrum of homogeneous light described in the former article, so to fall that some part of the light might pass through the hole in the paper. This transmitted part of the light I refracted with a prism placed behind the paper, and letting this refracted light fall perpendicularly upon a white paper two or three feet distant from the prism, I found that the spectrum formed on the paper by this light was not oblong, as when it is made, in the first experiment, by refracting the sun's compound

Newt. Opt.  
p. 108.

III.  
Experiment.  
Homogeneous  
light is re-  
fracted regu-  
larly, &c.  
Newt. Opt.  
p. 62.



light, but was (so far as I could judge by my eye) perfectly circular, the length being no where greater than the breadth; which shews that this light is refracted regularly without any dilatation of the rays; and is an ocular demonstration of the mathematical proposition mentioned in the 28th article.

IV.  
Experiment.  
Newt. Opt.  
p. 63.

In the homogeneous light I placed a paper circle of a quarter of an inch in diameter; and in the sun's unrefracted, heterogeneous, white light I placed another paper circle of the same bigness; and going from these papers to the distance of some feet I viewed both circles through a prism. The circle illuminated by the sun's heterogeneous light appeared very oblong, the length being many times greater than the breadth. But the other circle illuminated with homogeneous light appeared circular and distinctly defined, as when it is viewed by the naked eye; which proves the whole proposition mentioned at the beginning of this article.

V.  
Experiment.  
Ibid.

In the homogeneous light I placed flies and such like minute objects, and viewing them through a prism I saw their parts as distinctly defined as if I had viewed them with the naked eye. The same objects placed in the sun's unrefracted heterogeneous light which was white, I viewed also through a prism, and saw them most confusedly defined, so that I could not distinguish their smaller parts from one another. I placed also the letters of a small print one while in the homogeneous light and then in the heterogeneous, and viewing them through a prism they appeared in the latter case so confused and indistinct that I could not read them; but in the former they appeared so distinct that I could read readily, and thought I saw them as distinct as when I viewed them with my naked eye; in both cases I viewed the same objects through the same prism at the same distance from me and in the same situation. There was no difference but in the lights by which the objects were illuminated and which in one case was simple in the other compound; and therefore the distinct vision in the former case and confused in the latter could arise from nothing else than from that difference in the lights. Which proves the whole proposition.

The colour of  
homogeneous  
light cannot  
be changed  
by refractions  
nor by reflexions.  
Newt. Opt.  
p. 107.

30. In these three experiments it is farther very remarkable that the colour of homogeneous light was never changed by the refraction: and as these colours were not changed by refractions, so neither were they by reflexions. For all white, grey, red, yellow, green, blue, violet bodies, as paper, ashes, red lead, orpiment, indigo, bise, gold, silver, copper, grass, blue flowers, violets, bubbles of water tinged with various colours, peacocks' feathers, the tincture of *lignum-nephriticum* and such like, in red homogeneous light appeared totally red,

red, in blue light totally blue, in green light totally green, and so of other colours. In the homogeneous light of any colour they all appeared totally of that same colour, with this only difference, that some of them reflected that light more strongly, others more faintly. I never yet found any body which by reflecting homogeneous light could sensibly change its colour.

From all which it is manifest, that if the sun's light consisted of but one sort of rays, there would be but one colour in the world. Nor would it be possible to produce any new colour by reflexions and refractions: and by consequence that the variety of colours depends upon the composition of light.

31. Every homogeneous ray considered apart is refracted according to one and the same rule<sup>a</sup>, so that its sine of incidence is to its sine of refraction in a given ratio: that is, every different coloured ray has a different ratio belonging to it. This our author has proved by experiment, and by other experiments has determined by what numbers those given ratios are expressed. For instance, if an heterogeneous white ray of the sun emerges out of glass into air, or which is the same thing, if rays of all colours be supposed to succeed one another in the same line *AC*, and *AD* their common sine of incidence in glass be divided into 50 equal parts, then *EE* and *GH* the sines of refraction into air, of the least and most refrangible rays, will be 77 and 78 such parts respectively. And since every colour has several degrees, the sines of refraction of all the degrees of red will have all intermediate degrees of magnitude from 77 to  $77\frac{1}{8}$ , of all the degrees of orange from  $77\frac{1}{8}$  to  $77\frac{1}{4}$ , of yellow from  $77\frac{1}{4}$  to  $77\frac{1}{2}$ , of green from  $77\frac{1}{2}$  to  $77\frac{3}{4}$ , of blue from  $77\frac{3}{4}$  to  $77\frac{7}{8}$ , of indigo from  $77\frac{7}{8}$  to  $77\frac{8}{8}$ , and of violet from  $77\frac{8}{8}$  to 78<sup>b</sup>.

32. Colours may be produced by composition which shall be like to the colours of homogeneous light, as to the appearance of colour, but not as to the immutability of colour and constitution of light. And those colours, by how much they are more compounded, by so much are they less full and intense; and by too much composition they may be diluted and weakened till they cease, and the mixture becomes white or grey. There may be also colours produced by composition, which are not fully like any of the colours of homogeneous light. For a mixture of homogeneous red and yellow compounds an orange, like in appearance of colour to that orange which in the series of unmixed prismatic colours lies between them. But the light of one orange is homogeneous as to refrangibility, that of the other is heterogeneous; and the colour of the one, if viewed through a prism remains unchanged, that of the other is changed and resolved.

The sine of incidence of every homogeneous ray is to its sine of refraction in a given ratio. Newt. Opt. p. 64. Fig. 17. Art. 13.

<sup>a</sup> Newt. Opt. p. 109. The different properties of simple and compound colours. Newt. Opt. p. 115.



solved into its component colours red and yellow. And after the same manner other neighbouring homogeneous colours may compound new colours, like the intermediate homogeneous ones: as yellow and green the colour between them both; and afterwards if blue be added there will be made a green, the middle colour of the three which enter the composition. For the yellow and blue on either hand, if they are equal in quantity, draw the intermediate green equally toward themselves, and so keep it as it were in æquilibrium, that it verge not more to the yellow on one hand, than to the blue on the other, but by their mixed actions remain still a middle colour. To this mixed green there may be farther added some red and violet, and yet the green will not presently cease but only grow less full and vivid; and by increasing the red and violet, it will grow more and more dilute, until by the prevalence of the added colours it be overcome and turned into whiteness or some other colour. So if to the colour of any homogeneous light, the sun's white light composed of all sorts of rays be added, that colour will not vanish or change its species, but be diluted, and by adding more and more white it will be diluted more and more perpetually. Lastly if red and violet be mingled there will be generated according to their various proportions various purples: such as are not like in appearance to the colour of any homogeneous light; and of these purples mixed with yellow and blue may be made other new colours.

VI.  
Experiment.  
Whiteness  
may be com-  
pounded of  
colours.  
Newt. Opt.  
p. 117.  
Fig. 18.

33. Whiteness and all grey colours between white and black, may be compounded of colours; and the whiteness of the sun's light is compounded of all the primary colours mixed in a due proportion.

For let the solar image *PT* fall upon a lens *MN* above four inches broad and about six feet distant from the prism *ABC*, and so figured that it may cause the coloured light which divergeth from the prism to converge and meet again at its focus *G* about 6 or 8 feet distant from the lens, and there to fall perpendicular upon a white paper *DE*. And if you move this paper to and fro, you will perceive that near the lens, as at *de*, the whole solar image, suppose at *pt*, will appear upon it intensely coloured after the manner above explained: and that by receding from the lens those colours will perpetually come towards one another, and by mixing more and more dilute one another continually, until at length the paper comes to the focus *G*, where by a perfect mixture they will wholly vanish and be converted into whiteness, the whole light appearing now upon the paper like a little white circle.

VII.  
Experiment.  
Newt. Opt.  
p. 129.

In the foregoing experiment I have produced whiteness by mixing the prismatic colours. If now the colours of natural bodies are to be



be mingled, let a little water thickened with soap be agitated to raise a froth, and after that froth has stood a little, there will appear to one that shall view it intently various colours every where in the surface of the several bubbles, but to one that shall go so far off that he cannot distinguish the colours from one another, the whole froth will grow white with a perfect whiteness.

34. The colours of natural bodies arise from hence, that some of them reflect some sort of rays, others other sorts more copiously than the rest. Minium reflects the least refrangible or red-making rays most copiously and thence appears red. Violets reflect the most refrangible most copiously, and thence have their colour: and so of other bodies. Every body reflects the rays of its own colour more copiously than the rest, and from their excess and predominance in the reflected light has its colour.

The permanent colours of natural bodies explained.

For if in the homogeneous lights obtained by the 2d experiment, you place bodies of several colours, you will find as I have done, that every body looks more splendid and luminous in the light of its own colour. Cinnaber in the homogeneous red is most resplendent, in the green light it is manifestly less resplendent, in the blue light still less. Indigo in the violet blue light is most resplendent, and its splendor is gradually diminished as it is removed thence by degrees through the green and yellow light to the red. By a leek the green light, and next that the blue and yellow which compound green, are more strongly reflected than the other colours red and violet, and so of the rest. But to make these experiments the more manifest, such bodies ought to be chosen as have the fullest and most vivid colours, and two of those bodies are to be compared together. Thus for instance, if cinnaber and ultra-marine blue, or some other full blue be held together in the red homogeneous light, they will both appear red; but the cinnaber will appear of strongly luminous and resplendent red, and the ultra-marine blue of a faint obscure and dark red. And if they be held together in the blue homogeneous light, they will both appear blue; but the ultra-marine will appear of a strongly luminous and resplendent blue, and the cinnaber of a faint and dark blue. Which puts it out of dispute that the cinnaber reflects the red light much more copiously than the ultra-marine doth, and the ultra-marine reflects the blue light much more copiously than the cinnaber doth. The same experiment may be tried successively with red and indigo or with any other two coloured bodies, if due allowance be made for the different strength or weakness of their colour and light.

VIII. Experiment. Newt. Opt. p. 157.

And that this is not only a true reason of their colours, but even the

<sup>a</sup> Art. 30.

the only reason, may appear farther from this consideration; that the colour of homogeneous light cannot be changed by the reflection of natural bodies<sup>a</sup>. For if bodies by reflection cannot in the least change the colour of any one sort of rays, they cannot appear coloured by any other means, than by reflecting those which either are of their own colour, or by mixture must produce it.

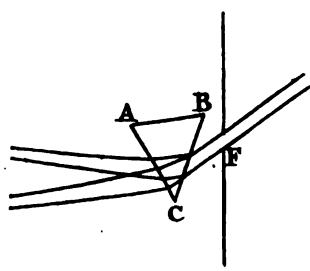
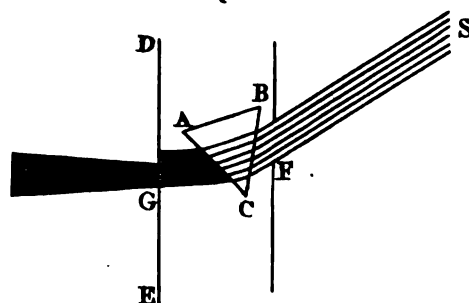
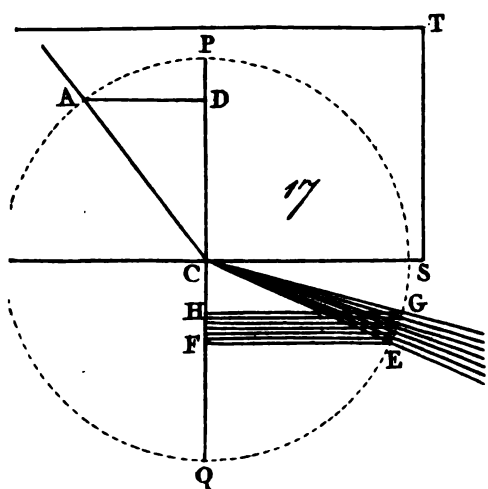
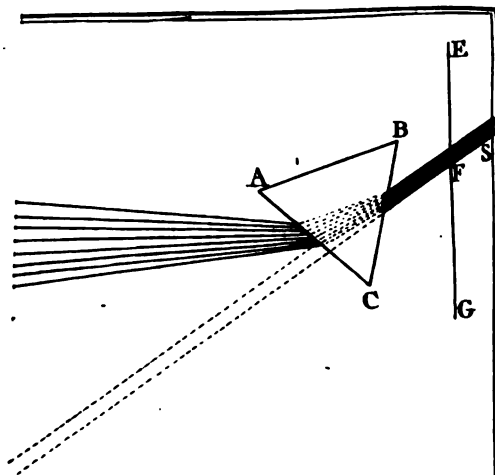
### CHAP. III.

CONCERNING THE CAUSE OF REFRACTION, REFLECTION, INFLECTION AND EMISSION OF LIGHT, AND CONCERNING TRANSPARENCY, OPACITY, AND COLOURS IN BODIES.

Reflexion not  
caused by the  
impinging of  
light upon  
the medium.  
Newt. Opt.  
p. 237.

<sup>b</sup> Art. 17.

35. **T**HAT the cause of reflection is not the impinging of light on the solid or impervious parts of bodies, as is commonly believed, will appear by the following considerations. First, that in the passage of light out of glass into air there is a reflection as strong as in its passage out of air into glass, or rather a little stronger, and by many degrees stronger than in its passage out of glass into water. And it seems not probable that air should have more reflecting parts than water or glass. But if that should possibly be supposed, yet it will avail nothing; for the reflection is as strong or stronger when the air is drawn away from the glass, by an air-pump, as when it is adjacent to it. Secondly, if light in its passage out of glass into air be incident more obliquely than an angle of 40 or 41 degrees, it is wholly reflected, if less obliquely it is in a great measure transmitted<sup>b</sup>. Now it is not to be imagined that light at one degree of obliquity should meet with pores enough in the air to transmit the greatest part of it, and at another degree of obliquity should meet with nothing but parts to reflect it wholly: especially considering that in its passage out of air into glass, how oblique soever be its incidence, it finds pores enough in the glass to transmit a great part of it. If any man supposes that it is not reflected by the air, but by the outmost superficial parts of the glass, there is still the same difficulty: besides that such a supposition is unintelligible, and will appear to be false by applying water behind some part of the glass instead of air. For so in a convenient obliquity of the rays suppose of 45 or 46 degrees, at which they are all reflected where air is adjacent to the glass, they shall be in a great measure transmitted where water is adjacent to it. Which argues that their reflection or trans-  
mission







mission depends on the constitution of air and water behind the glass, and not on the striking of the rays on the parts of the glass. Thirdly, if the colours made by a prism, placed at the entrance of a beam of light into a darkened room, be successively cast on a second prism placed at a great distance from the former, in such manner that they are all alike incident upon it, (as they will be when transmitted through the holes in the two boards made use of in the 2d Fig. 16. experiment,) the second prism may be so inclined to the incident rays, that those which are of a blue colour shall be all reflected by it, and yet those of a red pretty copiously transmitted. Now if reflection be caused by the parts of air or glass, I would ask why at the same obliquity of incidence, the blue should wholly impinge on those parts, so as to be all reflected, and yet the red find pores enough to be in a great measure transmitted? Lastly were the rays of light reflected by impinging on the solid parts of bodies, their reflections from polished bodies could not be so regular as they are. For in polishing glass with sand, putty or tripoli, it cannot be imagined that those substances can by grating and fretting the glass bring all its least particles to an accurate polish; so that all their surfaces shall be truly plane or truly spherical, and look all the same way, so as together to compose one even surface. This manner of polishing with powders can do no more than bring the roughness of the glass to a very fine grain, so that the scratches and frettings of the surface become too small to be visible. And therefore if light were reflected by impinging upon the solid parts of the glass, it would be scattered as much by the most polished glass as by the roughest. So then it remains a problem how glass polished by fretting substances can reflect light so regularly as it does.

36. And this problem is scarce otherwise to be solved than by saying that the reflection of a ray is effected not by a single point of the reflecting body, but by some power of the body which is evenly diffused all over its surface, and by which it acts upon a ray without immediate contact. For that the parts of bodies do act upon light at a distance, will appear by the following experiments.

37. The sun shining into my chamber through a hole a quarter of an inch broad, I placed at the distance of two or three feet from the hole a sheet of pastboard, which was blacked all over on both sides, and in the middle of it had a hole about three quarters of an inch square for the light to pass through. And behind the hole I fastened to the pastboard with pitch the blade of a sharp knife, to intercept some part of the light which passed through the hole. The planes of the pastboard and of the knife were parallel to one another

IX.  
Experiment.  
This power  
acts upon  
light at a  
distance by  
attracting and  
repelling it.  
Newt. Opt.  
p. 300.

other and perpendicular to the rays. And when they were so placed that none of the sun's light fell upon the pastboard, but all of it passed through the hole to the knife, and there part of it fell upon the blade of the knife, and part of it passed by its edge; I let this part of the light, which passed by, fall on a white paper two or three feet beyond the knife, and there saw two streams of faint light shoot out both ways from the beam of light into the shadow, like the tails of comets. But because the sun's direct light by its brightness upon the paper obscured these faint streams so that I could scarce see them, I made a little hole in the midst of the paper for that light to pass through, and fall upon a black cloth behind it, and then I saw the two streams plainly. They were like one another, and pretty nearly equal in length and breadth, and quantity of light. Their light at that end next the sun's direct light was pretty strong for the space of about a quarter of an inch or half an inch, and in all its progress from that direct light decreased gradually till it became insensible. The whole length of either of these streams measured upon the paper, at the distance of three feet from the knife, was about six or eight inches; so that it subtended an angle at the edge of the knife of 10 or 12, or at most 14 degrees.

X.  
Experiment.

I placed another knife by this, so that their edges might be parallel and look towards one another, and that the beam of light might fall upon both knives and some part of it pass between their edges. And when the distance of their edges was about the 400th part of an inch, the stream parted in the middle and left a shadow between the two parts. This shadow was so black and dark that all the light which passed between the two knives seemed to be bent and to be turned aside to the one hand and to the other. And as the knives still approached one another the shadow grew broader, and the streams shorter at their inward ends next the shadow, until upon the contact of the knives the whole light vanished and left its place to the shadow. And hence I gather that the light which is least bent, and goes to the inward ends of the streams, passes by the edges of the knives at the greatest distance, and this distance when the shadow begins to appear between the streams is about the 800th part of an inch. And the light which passes by the edges of the knives at distances still less and less is more and more bent, and goes to those parts of the streams which are farther and farther from the direct light. Because when the knives approach one another till they touch, those parts of the streams vanish last which are farthest from the direct light.

Our author has made it appear from these and some other experiments,



ments, that bodies act upon light in some circumstances by an attractive and in others by a repulsive power. For he found that the shadows of hairs, threads, pins, straws, and such like slender substances, placed in a slender beam of light let into a dark room, were considerably broader than they ought to be, if the rays of light passed on by these bodies in right lines. Particularly he found that the shadow of a hair of a man's head, at the distance of 10 feet from the hair, was 35 times broader than the hair itself<sup>a</sup>.

<sup>a</sup> Newt. Opt,  
p. 293.

38. That this power which acts upon light is infinitely stronger than the power of gravity will appear by the following argument. Sir *Isaac Newton* has demonstrated that all bodies attract one another by the force of gravity, and that the attractive forces of two homogeneous spheres, upon particles of matter placed very near their surfaces, are to each other in proportion as the diameters of the spheres<sup>1</sup>. That is to say, if a refracting medium be spherical and of the same density as the earth, the earth's force of attraction near its surface, will exceed the medium's force near its surface, as much as the diameter of the earth exceeds the diameter of the medium; or almost infinitely with respect to human conceptions. Yet we observe that a cannon-ball, just shot from the mouth of the cannon, is scarce sensibly deflected towards the earth by its attraction; and the least particle of the ball, if it was separate from the rest, would be no more deflected than the whole; because gravity makes bodies of all sorts and sizes descend with the same swiftness, by affecting them alike whether joined or separated. Therefore a particle of light which moves, I may say, infinitely quicker than a cannon-ball, would be infinitely less bent than the particle of the ball by the attraction of the whole earth, and still infinitely less, than this last bending, by the attraction of the spherical medium, which was shewn to be infinitely weaker than that of the earth. But in fact we find it is very sensibly bent or refracted; and therefore it must be affected by some other power of the medium, which near its surface is infinitely stronger than the power of gravity. And is infinitely stronger than the power of gravity.

39. It is difficult to determine the exact law of this refractive power, or the degrees of its force at given distances from the refracting surface. However since we find that the effects of gravity, which decrease as the squares of the distances from the center increase, are very sensible at great distances, we may conclude that the refractive power of a medium, which at its surface we find is infinitely stronger than gravity, and yet vanishes at a very small distance

And decreases  
much quicker.  
er.

<sup>1</sup> Princip. lib. 1. prop. 74. cor. 2. & lib. 3. prop. 8.



<sup>a</sup> Art. 37.

from it<sup>a</sup>, decreases much quicker or in a greater proportion than gravity does.

This one  
power both  
refracts and  
reflects light.  
Nemt. Opt.  
p. 244.  
<sup>b</sup> Art. 17.

40. It is reasonable to conclude that bodies reflect and refract light by one and the same power variously exercised in various circumstances; because when light goes out of glass into air as obliquely as it can possibly do, if its incidence be made still more oblique, it becomes totally reflected<sup>b</sup>: (for the power of the glass after it has reflected the light as obliquely as is possible, if the incidence be still made more oblique becomes too strong to let any of its rays go through and by consequence causes total reflections:) And for this other reason, that those surfaces of transparent bodies which have the greatest refracting power, reflect the greatest quantity of light, as will be shewn in the 47th article.

Its forces in  
different bo-  
dies are as  
their densities  
nearly.  
Nemt. Opt.  
p. 245.

41. From the different ratios of the sines of incidence and refraction in a great many different bodies, our author has also collected that the forces of bodies to reflect and refract light are very nearly proportionable to their densities, excepting that unctuous and sulphureous bodies refract more than others of the same density. Whence, he says, it seems rational to attribute the refractive power of all bodies chiefly, if not wholly to the sulphureous, oily particles with which they abound. For it is probable that all bodies abound more or less with sulphurs. And as light congregated by a burning glass acts most upon sulphureous bodies to turn them into fire and flame, so since all action is mutual, sulphurs ought to act most upon light. For that the action between light and bodies is mutual, may appear from this consideration; that the densest bodies which refract and reflect light most strongly grow hottest in the summer sun, by the action of the refracted or reflected light. If bodies be conceived to have certain densities exactly proportionable to their refractive powers, these may be called their refractive densities.

Refractive  
densities  
what.

This force  
acts in lines  
perpendicu-  
lar to the re-  
fracting sur-  
face.

<sup>c</sup> Nemt. Opt.  
p. 323, &c.

42. The direction of the refractive force of a medium, acting upon particles of light, is every where perpendicular to the refracting surface. For whether this force be a real attraction, or whether it be an impulse upon light, caused by the spring or elastick power of a subtil fluid which pervades the medium, and being gradually denser without than within it, may impel the light towards the medium by its greater elasticity without than within<sup>c</sup>; be this as you please, yet if the medium be uniform in all its parts, its immediate power upon the light it self, or upon the subtil fluid which acts upon it, will be equally strong in every point of a plane drawn parallel to the refracting surface; though its strength may be different in the next parallel plane, and still different in the next, and so on

as



as far as that power is extended on each side of the surface of the medium. The extent of this power will therefore be terminated by two planes, parallel to one another and to the refracting surface; and the space between them may be called the space of activity, whether the power attracts or repels. This being premised, I say the force of the medium will act upon light, either in attracting or repelling it, in lines perpendicular to its surface. For let  $p$  be a particle of light acted upon by any uniform power in the line  $de$  parallel to the refracting surface  $AB$ ,  $pc$  a line perpendicular to those parallels, cutting  $de$  in  $c$ ; it is evident that the force of the power at  $c$  will move the particle  $p$  in the line  $pc$ ; and taking any two points  $d, e$  at equal distances on each side of  $c$ , the powers at  $d$  and  $e$  being equal and acting at equal distances,  $pd, pe$ , equally inclined to  $pc$ , cannot move  $p$  in any direction but that of  $pc$ ; and what has been said of the equal powers in the line  $de$  is applicable to the powers in every line drawn parallel to  $AB$ , that is to the whole power of the refracting medium.

Space of activity what.

Fig. 19, 20.

43. Now when a ray of light falls perpendicularly upon the space of activity its particles will be accelerated or retarded in the same perpendicular direction, according as the power of the medium acts with or against the course of their motion; and when the particles are got through that space they will proceed with an uniform velocity. But if a ray  $op$  or  $sr$  falls obliquely upon the space of activity  $klmn$ , the force of the medium now acting sideways or obliquely upon the particles, will bend their course into a curve  $pqr$ , during their passage through that space. For as light has this property in common with all other bodies, of moving straight forwards, while its motion is not disturbed by any oblique force, so when it is disturbed, we may reasonably conclude, it will follow those other laws of motion, to which all other bodies are equally subject. The force of the medium acting sideways upon its oblique course, will therefore draw it perpetually out of one direction into another. But having passed through the space of activity, it will then proceed straight forwards; for being attracted or impelled every way alike, or else not at all if it be in empty space, it will have the same freedom of motion in both cases: just as an animal surrounded with air, though violently pressed on every side, feels no constraint, but has an equal facility of moving in any direction. Thus we see that the refraction of light is performed in the same manner as if a stone was thrown in the direction  $op$ , and its course was bent into a curve  $pqr$  by its gravity; or being thrown the contrary way in the direction  $sr$ , it was bent into the curve  $rqp$  in ascending: and supposing the

The manner of its operation in causing refractions and reflections.  
Fig. 21.



Fig. 22.

the attraction of the earth to reach no higher than the line  $kl$  the stone would from thence proceed in a straight line  $po$ . Now the gravity of the stone may be so great, or the force of projection so weak, or the direction of the motion so oblique to the horizontal line  $kl$ , that it cannot ascend so high as this line. In this case the stone will descend from the highest point of its course by the same degrees of curvity with which it ascended; and if its gravity be supposed to cease in all places below the line  $mn$ , the stone will go on in the direction of the last particle of the curve produced. This is a parallel case to that of reflections from the farther surface of dense mediums, when the incident ray is so much inclined to that surface as to be pulled back into the same medium. Hitherto I have supposed the refracting medium to be contiguous to empty space; but the manner of reflection and refraction is the same at the common surface of any two mediums. For since the separate forces of the mediums act in the same lines, perpendicular to their common surface<sup>a</sup>, and in contrary directions; the light will be affected with the difference of those forces in the same manner as before. And if the mediums have equal forces they will balance each other, without causing any reflection or refraction at all. It has been observed already that the perpendicular breadth of the space of activity is exceeding small, and consequently in physical experiments the incurvation of the ray may still be considered as performed in a physical point.

\* Art. 42.

And in causing the different refrangibility of rays.  
Newt. Opt.  
p. 347.

Fig. 23.

44. According to this theory nothing more is requisite for producing all the variety of colours and degrees of refrangibility, than that the rays of light be bodies of different sizes; the least of which may make violet, the weakest and darkest of the colours, and be more easily diverted by refracting surfaces from its right course; and the rest, as they are bigger and bigger, may make the stronger and more lucid colours, blue, green, yellow and red; and be more and more difficultly diverted. For particles of different sizes, that fall upon the space of activity  $klmn$  in the line  $op$ , having different forces, may describe different curves, as  $pa$ ,  $pb$ ,  $pc$ , and consequently will emerge from that space in different angles.

And in causing the angles of reflection to be the same in all sorts of rays.  
Fig. 23.

45. Thus may heterogeneous particles diverge from one another by refraction, though not by reflection. For if the line of their incidence  $op$  be so oblique to the space of attraction  $klmn$ , that all the particles are pulled back into the same medium, they will return in parallel lines  $rs$ ,  $tv$ ,  $xy$ , &c. inclined to that space in the same angles as the line of incidence  $op$  is inclined to it. Just as several balls of different sizes, shot with different forces out of a cannon

non

non  $op$  in any fixt position, will describe different curves, as  $pdr$ ,  $pet$ ,  $pfx$ , &c. yet in returning to the ground they will all strike upon it in equal angles, at  $r$ ,  $t$ ,  $x$ , &c. every one being equal to the angle of elevation at  $p$ . Now since the space of attraction is exceeding thin, the parallel lines  $rs$ ,  $tv$ , &c. will be so close together that the sense cannot perceive a distinct sensation of the separated particles, and consequently the reflected and incident light will appear of the same colour. And when the incident light consists of several rays, though the particles in each ray may be a little separated after reflection, and proceed in different lines, yet those several lines will be mixt together, and consequently the reflected light will appear white or of the same colour as the incident light.

46. Sir *Isaac Newton's* notion of the cause and manner of reflection from opake bodies, and from the first surface of transparent bodies, seems to be this that follows. Let the attractive power of the dense medium  $ABCD$  end at the line  $kl$ , and there let the repulsive power begin<sup>a</sup>, and let it end at the parallel line  $bi$ ; and when a ray  $op$  falls from air upon the space of repulsion  $bikl$ , it will be perpetually diverted from one direction into another by the opposition of the repulsive force, and so will describe a curve  $pqr$ , till it emerges from that space in the same angle at  $r$  with which it immersed at  $p$ , and then it will proceed in a right line  $rs$ . This will be the course of the ray if its progressive force be but weak, or the repulsive force be so strong as to hinder it from entering the space of attraction  $klmn$ . For if it enters this space, instead of being reflected, it will be refracted into the dense medium. And in reality some part of the incident light is always reflected and some refracted at all transparent surfaces; the cause of which our author has also considered<sup>b</sup>.

And in causing reflections from opake bodies and from the first surface of transparent ones.

Fig. 24.  
Art. 37.

<sup>b</sup> Newt. Opt. P. 253.

Hence it seems to follow that the repulsive power of a dense medium is less extended or else weaker than the attractive. For if the bending of a ray by the repulsive power, was not less than the contrary bending made by the attractive, the refraction into a dense medium could not always be made towards the perpendicular, as it always is. We may also observe that a refracted ray, in its passage through the surface of a transparent medium, is bent backward and forward with a motion like that of an eel; and our author takes notice of the same sort of motion in its passage by the edges and sides of bodies. It follows also that the repulsive power does not extend to a sensible distance from the medium; for if it did, it would be discovered by a sensible incurvation of the ray throughout that extent; contrary to experience.

47. Those



Stronger and  
weaker reflections how  
caused.  
Newt. Opt.  
p. 220.  
<sup>a</sup> Art. 41.

<sup>b</sup> Art. 17.

47. Those superficies of transparent bodies reflect the greatest quantity of light which have the greatest refractive power; that is which intercede mediums that differ most in their refractive densities<sup>a</sup>: and in the confines of equally refracting mediums there is no reflection. The analogy between refraction and reflection will appear by considering, that when light passes obliquely out of one medium into another which refracts from the perpendicular, the greater is the difference of the refracting densities, the less obliquity of incidence is requisite to cause a total reflection<sup>b</sup>. Those superficies therefore which refract most, do soonest reflect all the light which is incident upon them, and so must be allowed most strongly reflective. But the truth of this proposition will farther appear by observing, that in the superficies interceding two transparent mediums (such as are air, water, oil, common glass, crystal, metalline glass, island glasses, white transparent arsenick, diamonds, &c.) the reflection is stronger or weaker accordingly as the superficies hath a greater or less refractive power. For in the confine of air and sal-gem it is stronger than in the confine of air and water, and still stronger in the confine of air and common glass or crystal, and stronger in the confine of air and a diamond. If any of these and such like transparent solids, be immersed in water, its reflection becomes much weaker than before, and still weaker if they be immersed in the more strongly refracting liquors of well rectified oil of vitriol or spirit of turpentine. If water be distinguished into two parts by an imaginary surface, the reflection in the confine of these two parts is none at all; in the confine of water and ice it is very little; and in that of water and oil it is something greater; in that of water and sal-gem still greater, and in that of water and glass or crystal or other denser substances still greater, accordingly as those mediums differ more or less in their refracting powers. Hence in the confine of common glass and crystal there ought to be a weak reflection, and a stronger reflection in the confine of common and metalline glass, though I have not yet tried this. But in the confine of two glasses of equal densities, as of two object-glasses of long telescopes pressed gently together, there is not any sensible reflection. For objects may be seen by rays obliquely transmitted through a round black spot where the glasses touch one another, but not through other places where the light is reflected at the interval between the glasses. And the same may be understood of the superficies interceding two crystals, or two liquors, in which no reflection is caused. So then the reason why uniform pellucid mediums, such as water, glass, or crystal, have no sensible reflection, but in their external superficies,



perficies, where they are adjacent to other mediums of a different density, is because all their contiguous parts have one and the same degree of density: or this uniform density of their contiguous parts is a necessary condition of the transparency of the whole.

48. The least parts of almost all natural bodies are in some measure transparent: and the opacity of those bodies ariseth from the multitude of reflections caused in their internal parts. That this is so hath been observed by others, and will easily be granted by them that have been conversant with microscopes. And it may also be tried by applying any substance to a hole, through which some light is immitted into a dark room. For how opaque soever that substance may seem in the open air, it will by that means appear very manifestly transparent, if it be of a sufficient thinness. Only white metalline bodies must be excepted, which by reason of their excessive density seem to reflect almost all the light incident on their first superficies; unless by solution in menstruums they be reduced into very small particles, and then they become transparent.

Opacity caused by a multitude of internal reflections.  
Newt. Opt. p. 222.

49. Between the parts of opaque and coloured bodies are many spaces, either empty or replenished with mediums of other densities; as water between the tinging corpuscles wherewith a liquor is impregnated; air between the aqueous globules that constitute clouds or mists; and for the most part spaces void of both air and water, but yet perhaps not wholly void of all substance, between the parts of hard bodies. The truth of this is evident by the two preceding articles. For by the latter article there are many reflections made by the internal parts of bodies, which by the former article would not happen if the parts of these bodies were continued, without any such interstices between them: because reflections are caused only in superficies which intercede mediums of a different density by article 47.

The constitution of opaque and coloured bodies what.  
Newt. Opt. p. 223.

But farther, that this discontinuity of parts is the principal cause of the opacity of bodies, will appear by considering, that opaque substances become transparent by filling their pores with any substance of equal or almost equal densities with their parts. Thus paper dipped in water or oil, the *oculus mundi* stone steeped in water, linen cloth oiled or varnished, and many other substances soaked in such liquors as will intimately pervade their little pores, or separating parts, become by that means more transparent than otherwise. So on the contrary, the most transparent substances may by evacuating their pores, or separating parts, be rendered sufficiently opaque; as salts or wet paper or the *oculus mundi* stone by being dried; horn by being scraped; glass by being reduced to powder or otherwise

D

flawed;



flawed; turpentine by being stirred about with water till they mix imperfectly; and water by being formed into many small bubbles, either alone in the form of froth, or by shaking it together with oil of turpentine or oil of olive or with some other convenient liquor, with which it will not perfectly incorporate.

The constitu-  
tion of trans-  
parent bodies  
what.  
Newt. Opt.  
p. 225.

50. The parts of bodies and their interstices must not be less than of some definite bigness, to render them opaque and coloured. For the opaquest bodies, if their parts be subtilly divided, (as metals by being dissolved in acid menstruums, &c.) become perfectly transparent; and at the superficies of the object-glasses, mentioned in the 47th article, where they were very near to one another though they did not absolutely touch, there was no sensible reflection. And likewise if a bubble be blown with water first made tenacious by dissolving a little soap in it, and be covered with a clear glass, to defend it from being agitated by the external air, and be suffered to rest a while, till by the continual subsiding of the water it becomes very thin; at the top where it is thinnest, there will grow a round, black, spot (like that between the object-glasses) which will continually dilate it self more and more till the bubble breaks; now this spot appears black and transparent for want of a sensible reflection, whereas the sides of the bubble which are thicker than the top appear coloured and opaque by a strong reflection.

On these grounds I perceive it is that water, salt, glass, stones and such like substances are transparent. For upon diverse considerations they seem to be as full of pores or interstices between their parts as other bodies are, but yet their parts and interstices to be too small to cause reflections in their common surfaces.

## CHAP. IV.

CONCERNING THE REFRACTIONS OF A SINGLE RAY OF LIGHT  
IN ITS PASSAGE THROUGH A PRISM, GLOBE, OR LENS.

51. **A** RAY of light  $EF$  falling obliquely upon a flat piece of glass, or any medium terminated by two parallel planes represented by the lines  $AB, CD$ , will emerge from it after both refractions at  $F$  and  $G$  in a line  $GH$  parallel to the incident ray  $EF$ . For since any line  $FG$  which the ray describes in passing between the parallel planes, is equally inclined to them both<sup>a</sup>, it will be bent as much at  $G$  in going forward, as it would be at  $F$  in going backward<sup>b</sup>; and these equal bendings being made contrary ways, the incident and emergent rays  $EF$  and  $GH$  are therefore parallel.

52. The lines described by the incident and emergent rays  $EF$  and  $GH$ , being produced are closer together when the glass is thinner, and also when the ray falls less obliquely upon it; because the bendings at  $F$  and  $G$  are then less<sup>c</sup>: and in these cases if the glass be not flat but bent a little as represented in the 26th figure by two parallel arches  $AB, CD$ , the line  $EF, GH$  will still be nearly parallel. For the bended surfaces refract the ray  $EFGH$  just as much as two planes would do supposing they touched the surfaces at  $F$  and  $G$ <sup>d</sup>: and these planes will be nearly parallel when the line  $FG$  is but little inclined to the surfaces; being exactly so when it stands perpendicular to them both.

53. A thin piece of glass or of any transparent substance bounded on one side by a polished plane surface, represented by the line  $EF$ , and on the other side by a small portion of a polished spherical surface, represented by the arch  $ACB$ ; or bounded on both sides by spherical surfaces  $ACB, EDF$ , is called a lens or simply a glass; and is conceived by mathematicians to be generated or described by turning the figure  $ACBFDE$  round about the line  $CD$ , drawn through the middle of it perpendicularly to both its sides. This line  $CD$  produced is therefore called the axis of the lens; and it passes through  $G$  and  $H$ , the centers of its surfaces. The points  $C, D$  where it cuts the surfaces are called the vertexes of the lens, and the middle point between them is called its center. The 27th figure represents a plano-convex glass, the 28th a plano-concave, the 29th a double-convex, the 30th a double-concave, and the 31st and

Refraction of  
a ray through  
parallel plane  
surfaces.  
Fig. 25.

Eucl. I. 29.

Art. 11.

Refraction of  
a ray through  
parallel spheri-  
cal surfaces.  
Art. 16.

Art. 19.

A lens what.  
Fig. 27. to 32.



32d two concavo-convex glasses, whereof the first is called a meniscus, because it resembles a little moon. It must be remembered once for all, that the thickness  $CD$  of all these glasses is generally so small, that it seldom need be considered <sup>a</sup>.

<sup>a</sup> Art. 61.

Refraction of  
a single ray  
through a  
prism.

Fig. 33, 34,  
35.

54. When a ray of light  $EFGH$  is refracted at  $F$  and  $G$  in passing through the sides,  $AB, BC$ , of a prism, the course of the emergent ray,  $GH$ , always deviates from,  $EF$ , the course of the incident ray, towards the thicker part of the prism, more or less, as the refracting angle  $ABC$  is greater or smaller. And if the refracting angle be given (or invariable) and the refractions be but small, the quantity of deviation will also be given, though the position of the incident ray be varied at pleasure.

Fig. 33.

For supposing at first that the ray  $FG$ , within the prism, is equally inclined to its sides  $AB, BC$ , as in fig. 33, it is evident from the position of the perpendiculars to those sides at the points  $F$  and  $G$ , that both the refractions are made from the edge  $B$  towards the opposite side  $AC$  <sup>b</sup>.

<sup>b</sup> Art. 12.

Fig. 34.

Now let  $FG$  become unequally inclined to the sides  $AB, BC$ , by turning it gradually into the position  $fg$ ; and while it becomes less and less oblique to one side, suppose  $AB$ , it will become more and more oblique to the other side  $BC$ . Consequently supposing a ray to go both ways along this variable line  $fg$ , it will be more and more bent in going through the side  $BC$  and less and less in going back through the side  $AB$ ; so that the total bending of the ray, consisting of both its bendings, or angles  $efg$  and  $fgb$ , taken together, will continue to be much the same in all its positions. The circulation of the line  $fg$ , may be farther continued till it becomes perpendicular to the side  $AB$ ; and then the bending at this side is nothing: it may also be continued still farther till the bending at  $f$  is made the contrary way; which still takes off from the perpetual increase of the greater bending at  $g$  and keeps the total bending invariable.

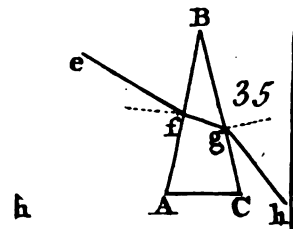
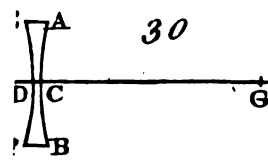
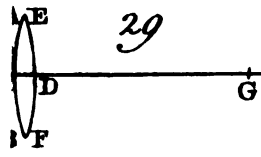
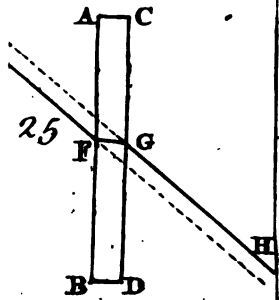
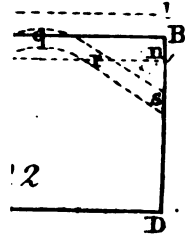
Fig. 35.

Fig. 34.

When  $fg$  is perpendicular to  $AB$ , let the latter plane  $BC$  be turned gradually towards the former  $BA$ , upon the edge  $B$ , and the ray that comes along  $fg$  will gradually fall less obliquely upon it; and consequently the bending at  $g$  will be gradually diminished <sup>c</sup>; and reduced to nothing when the refracting angle  $ABC$  vanishes. Lastly if several homogeneous rays be supposed to come parallel to one another they will all emerge parallel to one another <sup>d</sup>. Therefore the quantity of deviation of a ray does not at all depend upon its passage through a thicker or thinner part of the prism, nor upon its inclinations to the sides of the prism, but is proportioned to the quantity of the refracting angle  $ABC$ ; and the more exactly as this angle

<sup>c</sup> Art. 14, 15.

<sup>d</sup> Art. 23.







angle and the refractions at its sides are smaller. The truth of this conclusion is proved mathematically in the next article.

55. If the angles of incidence and refraction of a ray,  $\angle ACS$ , Fig. 36, 37. that passes through a very small angle of a prism,  $AIC$ , be so little as to be reckoned proportionable to their sines; the angle of deviation  $RFS$ , contained under the incident ray  $QAFR$  and the emergent ray  $SCFT$  produced, will be to the refracting angle  $AIC$ , as the difference of the sines of incidence and refraction to the lesser of them; and consequently the magnitude of the angle of deviation  $RFS$  will be invariable in all positions of the ray.

For let the perpendicular  $AB$ , to the first surface  $AI$ , cross  $CD$ , the perpendicular to the second, in  $E$ ; and supposing the ray  $AC$  to go both ways out of the prism, the angle of incidence  $ACD$  will be to the angle of emergence  $DCT$ , in the given ratio of the sine of incidence to the sine of refraction, that is of  $i$  to  $r$ ; and disjointly, we have  $ACD$  to  $ACT$  as  $i$  to  $r-i$ ; and the angle  $CAB$  is to  $CAR$ , in the same ratio, supposing the ray to go backward along  $CA$ ; and conjointly or disjointly we have  $ACD \pm CAB$  to  $ACT \pm CAR$ , that is  $AED$  or  $AIC$  to  $RFS$  in the same given ratio of  $i$  to  $r-i$ .  $\angle E. D.$

56. *Corol. 1.* Hence any two homogeneous emergent rays produced, will be inclined to one another in the same angle as the two incident rays are inclined to one another. For let the two incident rays  $QF, qf$  Fig. 38. (produced) meet in  $K$ ; and let the emergent rays  $SF, sf$  (produced) meet in  $L$ ; and let one of the incident rays cross the other emergent ray in  $M$ ; and since in the triangles  $KMF, LMf$ , the angles at  $M$  are equal and also those at  $F$  and  $f$  by article 55, it follows that the remaining angles at  $K$  and  $L$  are also equal.

57. *Corol. 2.* When the ray  $AC$  within the prism, coincides with Fig. 39. a perpendicular to either of the planes, as with  $AB$ ; one of the refractions will vanish at  $A$ ; and then the angle of deviation  $RFS$  made by the other single refraction, will continue the same in quantity as before, when it was made by two refractions; because the magnitude of the angle of deviation is invariable by article 55.

58. *Corol. 3.* Therefore when an heterogeneous ray is separated into coloured rays, by small refractions through a small refracting angle of a given quantity, the emergent rays of given colours will be inclined to one another and to the incident ray in certain given angles, in all positions of the incident ray. Because these inclinations made by two refractions, are every where equal to the inclinations made by a single refraction at the second plane, when the incident ray falls perpendicular upon the first plane.

59. When.



Refractions of  
a single ray  
through the  
edge of a  
lens, or the  
sides of a  
globe.

Fig. 40. to 47.

\* Art. 19.

<sup>b</sup> Art. 54.

59. When a ray of light  $EFGH$  passes through the edge of a convex or concave lens, or the sides of a globe, its emergent part  $GH$  always deviates from the course of the incident part  $EF$  towards the thicker part of the glass, if the medium in which they are placed, is rarer than the globe or lens. For the refractions at  $F$  and  $G$  are the same as if they were made by two planes  $FA, GC$ , that touch the spherical surface at  $F$  and  $G^a$ ; and so the sides of the glass may be considered as inclined to each other like the sides of a prism. But the course of the ray is bent towards the thicker part of the prism <sup>b</sup>, and therefore towards the thicker part of the globe or lens.

The 44th and 59th figures represent the refractions of a ray passing through a sphere placed within a medium denser than itself; the bending of which ray may, in like manner, be proved to be towards the thinner part of the sphere.

Refractions  
of a single ray  
through the  
middle of a  
lens.

Fig. 48. to 55.

60. From the same method of reasoning it follows, that the deviation of the course of the emergent ray from that of the incident ray is gradually diminished as the ray goes nearer and nearer to the middle of the glass; till, when it goes through the middle, its emergent and incident parts are either parallel to each other, or else are one continued line, when the ray coincides with the axis of the glass. For the angle made by the touching planes,  $FA, GC$ , is gradually diminished as the ray  $FG$  approaches to the middle: till at last it vanishes when they become parallel, as in the 51st article \*.

This ray is  
considered as  
straight, and  
is called the  
axis of a pen-  
cil.

<sup>c</sup> Art. 60.

61. When a pencil of rays falls upon any glass, that ray which passes through its center, or middle point, is called the axis of the pencil. And because its incident and emergent parts  $EF$  and  $GH$ , are either one continued line or two parallel lines <sup>c</sup>, its whole course in optical experiments may be always taken for one straight, physi-

\* The author's demonstration of this article being applicable to those cases only, in which the angle at the vertex of the touching prism and the angles of incidence are very small <sup>a</sup>, the following demonstration of the manner in which a ray is refracted by the sides of a globe is inferred from Dr. Barrow's Opt. Lect.

<sup>a</sup> Art. 55.

PROPOSITION. *The farther a ray is distant from the center of a globe through which it passes, the greater will be its deviation.*

Fig. 60.

Let  $NRSP$  represent a globe, which differs in density from the medium in which it is placed;  $RS, NP$ , two rays passing through it, of which  $RS$  is the more remote from the center  $C$ , than  $NP$  is. Draw the radii  $CN, CP, CR, CS$ , and the angle  $RCR$  will be less than the angle  $NCP$  <sup>b</sup>. Therefore the sum of the angles  $CRS, CSR$  is greater than the sum of the angles  $CNP, CPN$  <sup>c</sup>; and each of the angles  $CRS, CSR$  is greater than each of the angles  $CNP, CPN$  <sup>d</sup>. But the angle of refraction at  $R$  exceeding that at  $N$ , the angle of incidence at  $R$  will also be the greater <sup>e</sup>. Since therefore the angles of incidence at  $R$  and  $S$  are greater than those at  $N$  and  $P$ , the deviation of  $RS$  is greater than that of  $NP$  <sup>f</sup>.  $\text{Q. E. D.}$

<sup>b</sup> Euc. III. 15.  
& I. 25.

<sup>c</sup> Euc. I. 32.

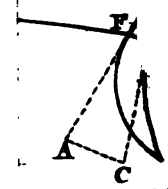
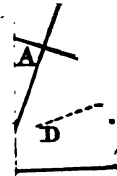
<sup>d</sup> Euc. I. 5.

<sup>e</sup> Art. 13.

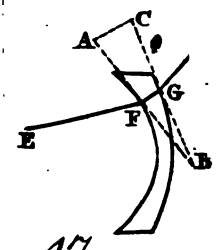
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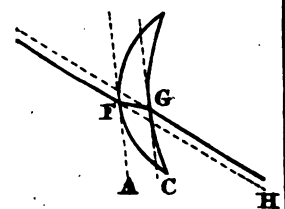
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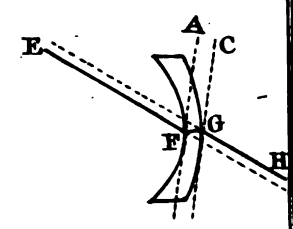
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47



51



55





cal line: from which it differs insensibly when the thickness of the glass is small, and when the pencil falls not too obliquely upon it. Because the parallel lines  $EF$  and  $GH$  produced, go closer together in proportion as the line  $FG$  is shorter, and as the bendings at  $F$  and  $G$  are smaller.

62. When two homogeneous rays are refracted through the same point of any lens, whose thickness is inconsiderable, the angle contained under their incident parts is equal to the angle under their emergent parts. For the thickness of the lens being very small, if the rays have a common point of incidence, their points of emergence will be very near one another; or if the point of emergence be common to both rays, their points of incidence will be very near one another; and still nearer if the rays cross one another within the lens; consequently the refractions through the lens will be nearly the same as through two planes, that touch its surfaces at two given points, near the points of incidence and emergence, and contain a given angle with one another<sup>a</sup>.

Refractions  
of two homo-  
geneal rays  
through the  
same point of  
any lens.

Fig. 57. 58.

<sup>a</sup> Art. 56.

63. An heterogeneous ray, refracted through a given point in a lens, has the same property as in the prism; that is, the emergent rays of given colours are inclined to one another and to the incident ray in given angles in all positions; for the reasons mentioned in the last article and in article 58.

Refractions  
of an hetero-  
geneal ray  
through the  
same point in  
any lens.

64. All rays, as  $EFGH$  and  $efgh$ , which cross each other in a refracting globe, and pass through it at equal distances from its center, so as to touch a concentric globe, are equally bent. For in this case the chords  $FG$ ,  $fg$  being equal, their obliquities to the surface of the globe are also equal, and consequently the bendings of the ray  $EFGH$  at  $F$  and  $G$ , both severally and together, are equal to the bendings of the ray  $efgh$  at  $f$  and  $g$ : as is evident by conceiving the rays to go both ways along the chords  $FG$ ,  $fg$ . Therefore the angle made by the incident and emergent parts of one ray, produced till they meet, will be equal to the angle made by the incident and emergent parts of the other ray produced till they meet; which is what I mean when I say the rays are equally bent.

Rays are  
equally bent  
which pass at  
equal distan-  
ces from the  
center of a  
globe.  
Fig. 56. to 59.

65. All rays, as  $EFGH$ ,  $efgh$ , which cross each other at any given point of a lens, or which pass through it at equal distances from its center, are equally bent, provided they do not fall very obliquely upon it. Imagine a line  $FG$  within the glass, at first to be equally inclined to its sides, and then to be turned a little about any point of it, till it comes into the position  $fg$ ; and while it becomes more and more oblique to one side of the glass, suppose  $Ff$ , it will become

And from the  
center of a  
lens.  
Fig. 57. 58.



<sup>a</sup> Art. 16.<sup>b</sup> Art. 54.

become less and less oblique to the other side  $Gg$ . Consequently if a ray be supposed to go both ways along this variable line  $fg$  it will be more and more bent in going through the side  $Ff$ , and less and less bent in going through the other side  $Gg$ <sup>a</sup>: so that the total bending of the ray consisting of both its bendings, or angles  $efg$ ,  $fgb$ , taken together, will continue to be much the same in all its positions. The circulation of the line  $fg$ , about the given point, may be farther continued till the bending at  $g$  is diminished to nothing; and still farther till it be made the contrary way; (as was explained in the 54th article;) which still takes off from the perpetual increase of the greater bending at  $f$  and keeps the total invariable. To keep the same bending it is only necessary that the rays  $FG$ ,  $fg$  should keep at equal distance from the axis of the lens as near as possible: and nothing alters the total bending but the alteration of that distance<sup>b</sup>; because the inclination of the tangent planes, like the refracting angle of a prism, will then only be altered.

## CHAP. V.

## TO FIND THE FOCUS OF RAYS REFLECTED FROM ANY GIVEN SURFACE.

## PROPOSITION I.

Fig. 61.

66. **L**ET  $ACB$  be a reflecting plane, and  $Q$  the focus of the incident rays, and  $QC$  a perpendicular to that plane; and if this perpendicular be produced to  $q$ , so that  $qC$  be equal to  $QC$ , the point  $q$  shall be the focus of the reflected rays.

<sup>a</sup> 8th Ax. Euc.<sup>b</sup> Art. 9.<sup>c</sup> Art. 11.

For let  $QA$  be any incident ray; draw  $qA$  and produce it towards  $O$ , and  $CA$  towards  $D$ . Then because  $Cq$  is made equal to  $CQ$ , the triangles  $CAq$ ,  $CAQ$  will be equal<sup>a</sup>. And consequently the angle  $DAO$ , which is equal to the opposite angle  $CAq$ , is also equal to the angle  $CAQ$ . Therefore  $AO$  is the reflected ray<sup>b</sup>.  $Q. E. D.$

67. *Corol.* Hence all the rays that flow towards  $q$ , will flow to  $Q$  after reflection<sup>c</sup>.

## LEMMA.

68. *Quantities and their proportions, which so approach to a state of equality as to become equal at last, may be taken for equal in a state immediately*



diately preceding the last; and also in a state somewhat remote from the last without sensible error in physical subjects: and the same may be said of figures which continually approach to a state of similitude; especially if these errors, when computed, are found inconsiderable.

The meaning of the lemma will appear very plain when applied to the following propositions.

## PROPOSITION II.

69. When parallel rays as  $DA$ ,  $EC$  fall almost perpendicularly upon a spherical surface  $ACB$ , the focus,  $T$ , of the reflected rays will bisect that semidiameter  $EC$ , which is parallel to the incident rays. Fig. 62, 63.

For drawing  $EA$ , it will be perpendicular to the spherical surface at  $A$ , and since  $EC$  is in the same plane as the angle of incidence  $DAE$ , the reflected ray  $Aq$  (produced backwards in fig. 63.) will meet  $EC$  somewhere in  $q^a$ ; so that the angle of reflection  $EAq$  may <sup>a</sup> Art. 7. be equal <sup>b</sup> to the angle of incidence  $EAD$ , or to the alternate angle <sup>b</sup> Art. 8.  $AEq$ . The two sides  $Aq$ ,  $Eq$  of the triangle  $AqE$  are therefore equal to each other <sup>c</sup>; and consequently each of them greater than <sup>c</sup> Euc. I. 6. half the third side  $EA$ , or than  $ET$  by construction. But as the point of incidence at  $A$  approaches towards  $C$ , the lines  $Eq$ ,  $ET$ , continually approach towards equality, and become equal when the triangle  $AEq$  is vanishing: and so the focus of rays falling almost perpendicularly on the surface, or the nearest to the point  $C$ , is to be reckoned at  $T^d$ .  $\text{Q. E. D.}$  <sup>d</sup> Art. 68.

70. *Corol.* Hence if  $T$  be the focus of the incident rays, the reflected ones will go parallel to the line  $TE^e$ . <sup>e</sup> Art. 11.

The point  $T$  is called the principal focus of the reflecter  $ACB$ , and  $TE$  its focal distance. And in general, when the rays come parallel to each other, that point to which they converge or from which they diverge after reflection or refraction, is called the *principal focus* of the reflecter or refracter. And the distance of that point from the center of the reflecter or refracter is called the *focal distance* of the reflecter or refracter, and by some authors its *focal length*; in figures 89, 90, 95, &c. to 100,  $F$  is the principal focus, and  $FE$  the focal distance.

## PROPOSITION III.

71. Let  $ACB$  be a reflecting surface of any sphere whose center is  $E$ . Fig. 64, 65. Bisect any radius thereof, suppose  $EC$ , in  $T$ ; and if in this radius, on the same side of the point  $T$ , you take the points  $Q$  and  $q$ , so that  $TQ$ ,  $TE$  and  $Tq$  be continual proportionals; and the point  $Q$  be the focus of the incident rays, the point  $q$  shall be the focus of the reflected ones.

E

Let



Let  $QA$  and  $Aq$  be an incident and reflected ray (produced) making equal angles with the perpendicular  $AE$ ; and the reflected ray  $Aq$  (produced) will cut  $QE$  (produced) somewhere in  $q$ , as being in the plane of incidence  $EAQ$ <sup>a</sup>. Draw  $EG$  parallel to  $Aq$ , and let it meet  $AQ$  in  $G$ ; and also  $Eg$  parallel to  $AQ$ , meeting  $Aq$  in  $g$ : then because the angles  $EAG$ ,  $EAg$  are equal<sup>b</sup>, it follows that the triangles  $EAG$ ,  $EAg$  are equiangular at their common base  $AE$ <sup>c</sup>; and therefore equicrural<sup>d</sup>; and also equal to each other; and consequently each side of the equilateral figure  $AGEg$ , in its vanishing state when  $A$  comes to  $C$ , will be equal to half its diagonal  $AE$ <sup>e</sup>, or by construction to  $ET$ . Now because the triangles  $GQE$ ,  $gEq$  are equiangular<sup>f</sup>, it will be as  $GQ$  to  $GE$ , so  $gE$  to  $gq$ <sup>g</sup>; that is, when the point  $A$  is coinciding with  $C$ , and consequently the points  $G$ ,  $g$  with  $T$ , as  $TQ$  to  $TE$  so  $TE$  to  $Tq$ <sup>h</sup>.  $Q.E.D.$

72. *Corol. 1.* If  $q$  be the focus of incident rays,  $Q$  will be the focus of the reflected ones<sup>i</sup>. And if either of these focuses recedes from  $T$ , the other will approach towards it. For the middle term  $TE$  of the continual proportionals  $TQ$ ,  $TE$ ,  $Tq$ <sup>k</sup>, being invariable in the same reflecter, the rectangle under the extremes is also invariable<sup>l</sup>; and therefore  $TQ$  varies as  $Tq$  inversely. These focuses will meet each other at the center  $E$ , and at the surface  $C$ . For at these points, the three terms in the proportion are equal<sup>m</sup>.

73. *Corol. 2.* The rays that belong to  $Q$  may be reckoned parallel when the distance  $TQ$  is infinite, and then by this proposition its reciprocal  $Tq$  becomes nothing; which is the second proposition.

Fig. 65. 74. *Corol. 3.* Hence also we may deduce the first proposition; for supposing  $Q$  the focus of incident rays upon the convex surface  $AB$ ; since  $TQ$ ,  $TC$ ,  $Tq$  are continual proportionals, it is well known that their differences  $CQ$ ,  $Cq$  must become equal when the lines themselves are infinitely great; that is when the surface becomes a plane by removing its center to an infinite distance.

The figures serve for the cases of a convex surface supposing the incident rays to go backwards in the same lines produced through the surface.

75. By the demonstrations of the two last propositions, it appears that the focus of reflected rays there determined, is nothing else in strictness of geometry, but the intersection of the axis of the surface, that is of the ray passing through its center, and of the nearest rays to it: and also that other rays intersect the axis in different points farther and farther from that focus, as they fall farther and farther from the vertex of the surface. So that a spherical surface cannot possibly reflect all the incident rays to a single point.



point\*. Nevertheless when these aberrations of the remoter rays from the geometrical focus shall be considered, it will appear hereafter, that the density of their intersections, near that focus, is immensely greater than their density at any considerable distance from it. So that in physical things, the focus of all the rays, that fall almost perpendicular upon a spherical surface, may be considered as a physical point. And the same is to be understood of the focus of refracted rays, as will appear by the like sort of demonstrations<sup>a</sup>. \* Chap. XI.

76. Hence it appears that the focus of rays reflected from any curved surface whatever, must be reckoned the same as if they were reflected from a spherical surface of an equal curvity to that surface about the points of incidence. As if  $CD$  be any curve whatever,  $C$  Fig. 66. the point of incidence,  $CE \perp$  perpendicular to the curve, or to its tangent at  $C$ ,  $CE$  the radius of a circle  $ACB$  of the same degree of curvity at  $C$ ; the rays coming parallel to  $CE$ , will be reflected to the same focus  $T$  from either of the surfaces; and also the rays that flow from any point  $Q$ , will be reflected by either surface to the same focus  $q$ . Because we consider the focus of those rays only, that fall

\* All the incident rays belonging to the same pencil may be reflected to a single point by means of surfaces, which are generated by the revolution of some conic section round its transverse axis. For

1st, Let  $AC$  represent a parabola, the axis of which is  $CQ$  and focus  $T$ ; describe a Fig. 67. paraboloid  $ACB$  by turning  $AC$  round its axis; draw  $EF$  touching its surface in any point  $A$ ; draw also the diameter  $dAD$ , join  $AT$  and produce it to  $t$ . Now the angles  $DAF$ ,  $TAE$  are equal<sup>a</sup>, and therefore the angle  $dAE$  equals  $tAF$ . Since then the incident and reflected rays are equally inclined to the reflecting plane<sup>b</sup>, if  $DA$  and  $dA$  be incident rays,  $AT$  and  $At$  will be their respective reflected rays; or if  $TA$ ,  $tA$  be incident rays,  $AD$ ,  $Ad$  will be their respective reflected rays. <sup>a</sup> Hamilton's Conics, II. 15. <sup>b</sup> Art. 9.

Hence, all the rays which are incident on the concave side of this paraboloid parallel to its axis will be reflected converging to the focus: and all the rays, which diverge from the focus will be reflected parallel to the axis. The convex surface will make parallel rays diverge from a single point, and converging rays go parallel, after reflection.

2dly, Let  $ACQ$  represent an ellipse, and  $ACB$  an hyperbola, the transverse axis of which is  $QC$ , and  $T$ ,  $D$ , the foci; let a spheroid be generated by turning the ellipse  $QAC$  round  $QC$ , and an hyperboloid by the motion of  $AC$  round its transverse axis; draw  $EF$  touching the surface in any point  $A$ ; join  $TA$ ,  $DA$ , and produce them to  $t$  and  $d$ . Since the angles, which  $TAi$ ,  $DAi$ , make with  $FE$  are equal<sup>c</sup>, and since the incident and reflected rays are equally inclined to the reflecting plane<sup>b</sup>, it is evident that, if  $TA$  or  $tA$  be an incident ray,  $Ad$  or  $AD$  will be the direction in which they are reflected. <sup>c</sup> Hamilton's Conics, II. 16. <sup>b</sup> & 17.

Hence all the rays, which are incident diverging from one focus of the spheroid, will be reflected converging to the other focus; and all the rays, which are incident upon the convex side of the spheroid converging to one focus, will be reflected diverging from the other. In like manner, the convex side of the hyperboloid will make a pencil of diverging rays diverge from, and a pencil of converging rays converge to, a single point after reflection. It will make diverging rays diverge more, and converging rays converge less. And the concave side of this reflector increases the convergency, and lessens the divergency, of a pencil of rays.



upon the common points of both curves about  $C$ , all the rest being dispersed much thinner into other places.

77. In all these propositions when the focuses  $\mathcal{Q}$ ,  $q$  lye on the same side of the reflecting surface, if the incident rays flow from  $\mathcal{Q}$  the reflected ones will flow towards  $q$ ; and if the incident rays flow towards  $\mathcal{Q}$ , the reflected ones will flow from  $q$ ; and the contrary happens when  $\mathcal{Q}$  and  $q$  are on contrary sides of the surface. Because the incident and reflected rays go contrary ways.

## CHAP. VI.

TO DETERMINE THE PLACE, MAGNITUDE AND SITUATION OF IMAGES FORMED BY REFLECTED RAYS.

### PROPOSITION I.

78. **I**MAGES formed by reflections from a plane surface are similar and equal to the objects; and their parts have the same situation with respect to the backside of the plane as the parts of the object have with respect to its fore-side.

Fig. 70, 71. From any number of points  $P$ ,  $\mathcal{Q}$ ,  $R$  of an object in any situation, draw the perpendiculars  $PA$ ,  $\mathcal{Q}C$ ,  $RB$  to the plane  $ACB$ , and produce them through it to the points  $p$ ,  $q$ ,  $r$ , each as far behind the plane as  $P$ ,  $\mathcal{Q}$ ,  $R$  are before it. The points  $p$ ,  $q$ ,  $r$  being the respective focuses of the rays that belonged to  $P$ ,  $\mathcal{Q}$ ,  $R$ , and being evidently in the same order, together with infinite others, will constitute an image of the object, equal to it in the whole and in every corresponding part, and alike situated: as will appear by conceiving the surface of the object, and of its image, divided into corresponding lines, such as  $P\mathcal{Q}R$ ,  $pqr$ , by planes such as  $PprR$  perpendicular to the reflecting plane; and by folding up or doubling each plane in its line of intersection,  $AB$ , with the reflecting plane. For the parts of each plane on each side of  $AB$  will exactly cover each other, as appears by the construction.  $\mathcal{Q}$ ,  $E$ .  $D$ .

\* Art. 66.

### PROPOSITION II.

Fig. 72. to 75. 79. If an arch of a circle  $PQR$ , concentrick to a concave or convex spherical surface  $AB$ , be considered as an object, its image  $pqr$  will also be a similar concentrick arch, whose length will be to the length of the object, in the ratio of their distances from the common center  $E$ ; and its situation will be erect or inverted, according as it is on the same or the opposite side of the center to the object.

For

For as the focus  $\mathcal{Q}$  was found by making  $T\mathcal{Q}$ ,  $TE$ ,  $Tq$  continual proportionals in the line  $\mathcal{Q}E$  drawn through the center<sup>a</sup>; so the focus  $p$ , of rays that belong to any other point  $P$ , is found by drawing  $PEA$ , and bisecting  $EA$  in  $S$ , and by making  $SP$ ,  $SE$ ,  $Sp$  continual proportionals. The two first terms of one proportion are severally equal to the two first of the other; and consequently the third terms  $Tq$ ,  $Sp$  are equal; and thence  $Ep$  and  $Eq$  are equal. The same being true of every point of the circular object  $P\mathcal{Q}R$ , shews that its image  $pqr$  is a concentrick arch, similar to it, both being terminated by the same lines  $EPp$ ,  $ERr$ ; and consequently their lengths are in the same ratio as their semidiameters  $E\mathcal{Q}$ ,  $Eq$ . Lastly, according as the corresponding extremities  $P$  and  $p$ , of the object and image, are on the same or opposite sides of the center  $E$ , they are also on the same or opposite sides of their middle points  $\mathcal{Q}$ ,  $q$ ; that is, the image is accordingly erect or inverted.  $\mathcal{Q}.E.D.$  Art. 71.

80. *Corol.* The smaller the circular object is with respect to its radius or distance from the center, the nearer it approaches in shape to a straight line, and so does its similar image. Consequently a small straight object, placed at a good distance from the center of the glass, may be reckoned to have a straight image very nearly: though in strictness of geometry it is an arch of a conick section.

81. The images of all sorts of objects may be determined, by finding the images of their out-lines, by the foregoing propositions. For instance, if the plane of the figures  $PER$ ,  $pEr$  be turned round their common diameter  $\mathcal{Q}Eq$ , the circular surface generated by  $pqr$  will be the image of the circular object generated by  $P\mathcal{Q}R$ : and if the same figures  $PER$ ,  $pEr$  be moved a little about an axis  $EF$ , situated in their own plane, and perpendicular to the diameter  $\mathcal{Q}Eq$ , the curvilinear figure generated by this motion of  $pqr$ , will be the image of a similar figure generated by  $P\mathcal{Q}R$ . Because the reflecting arch  $ACB$  generates the reflecting spherical surface at the same time.

82. But if the whole figure  $PERrp$  be moved parallel to itself in a direction  $EF$ , now perpendicular to its own plane, so that the arch  $ACB$  may generate a portion of a cylindrical surface, the figure described by this motion of  $pqr$ , will still be the image of that described by  $P\mathcal{Q}R$ ; but will not be similar to it, except when they are placed at equal distances on each side the center  $E$ , and consequently are equal to each other: and their dissimilitude will be so much the greater as the disproportion between  $Eq$  and  $E\mathcal{Q}$ , or between their lengths  $pr$ ,  $PR$ , is greater; their breadths, described by the motion aforesaid, being always equal to each other.



## CHAP. VII.

TO FIND THE FOCUS OF RAYS FALLING ALMOST PERPENDICULARLY UPON ANY REFRACTING SURFACE, SPHERE OR LENS.

## DEFINITION.

Fig. 76, 77. 83. **T**HE sine of an angle  $ABC$ , or of an arch  $AC$  that measures that angle, is a line  $AD$  drawn from the extremity of one of the semidiameters,  $AB$ ,  $BC$ , perpendicular to the other, produced if the angle be obtuse. And therefore an angle  $ABC$  and its complement  $ABE$ , to two right angles, have each the same sine  $AD$ ; and when the sines of several angles are compared together, they are always understood to belong to the same or to equal circles.

84. The sines of very small angles, and of their complements, become at last insensibly different from the arches that measure them; and consequently are proportionable to the angles themselves.

## LEMMA.

Fig. 78. 85. *The sines of the angles of any triangle are proportionable to the opposite sides: as in the triangle  $ABC$ , the sine of the angle  $ABC$  is to the sine of the angle  $BCA$ , as  $CA$  to  $AB$ .*

For the perpendiculars  $CD$ ,  $BE$  upon those sides  $AB$ ,  $AC$  produced, are the sines of the angles  $ABC$ ,  $BCA$  or  $BCE$  with respect to circles whose radius is  $BC^a$ . And since the triangles  $CAD$ ,  $BAE$  are equiangular<sup>b</sup>, we have  $CD$  to  $BE$  as  $CA$  to  $AB$ . Q. E. D.

<sup>a</sup> Art. 83.

<sup>b</sup> Euc. I. 32.

Cor. 2.

86. *Corol.* Small angles, as  $BAC$ ,  $BCE$ , subtended by the same perpendicular  $BE$ , are reciprocally as their legs  $BA$ ,  $BC$  or  $EA$ ,  $EC$ . For the angle  $BAC$  is to  $BCE$ , when very small, as the sine of  $BAC$  to the sine of  $BCE$ <sup>c</sup>, or as  $BC$  to  $BA$ <sup>d</sup>, or as  $EC$  to  $EA$ <sup>e</sup>.

<sup>c</sup> Art. 84.

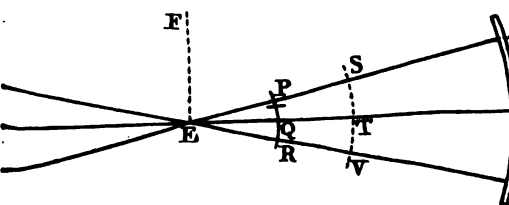
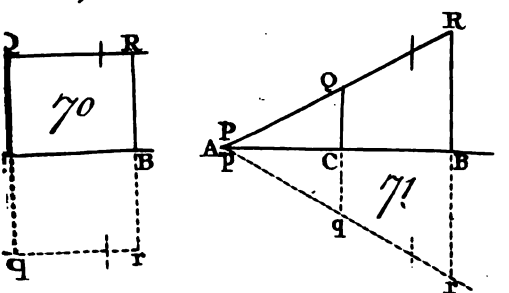
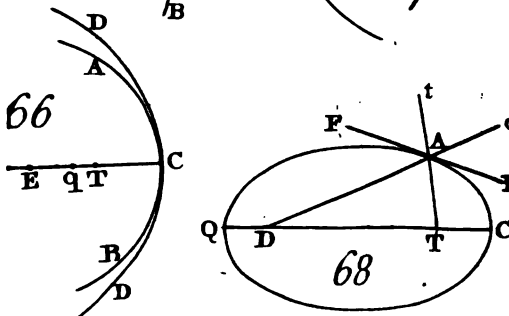
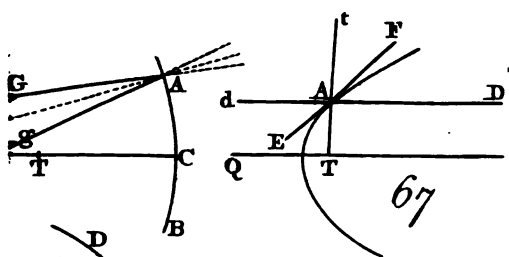
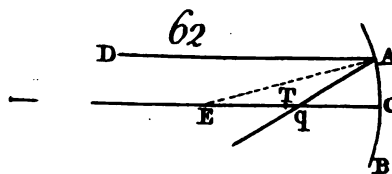
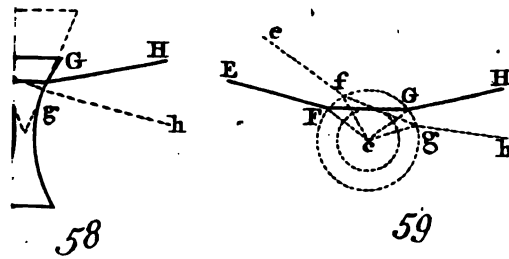
<sup>d</sup> Art. 85.

<sup>e</sup> Art. 68.

## PROPOSITION I.

Fig. 79. to 82. 87. *Let  $ACB$  be a refracting plane, and  $Q$  the focus of the incident rays, and  $QC$  a perpendicular to that plane; and if  $qC$  be taken in this perpendicular, on the same side of the plane as  $QC$ , and in proportion to  $QC$ , as the sine of incidence to the sine of refraction, the point  $q$  shall be the focus of the refracted rays.*

For







## PROPOSITION II.

Fig. 85. to 88. 92. Let  $ACB$  represent a refracting spherical surface whose center is  $E$ , and let the incident rays as  $DA$  come parallel to any semidiameter  $CE$  in which produced forward or backward, according as the denser medium is convex or concave, take  $CT$  to  $CE$  as the sine of incidence to the difference of the sines, and  $T$  will be the focus of the refracted rays.

For let the refracted ray  $AT$  (produced) cut the semidiameter  $CE$  produced, in any point  $T$  whatever; and since the semidiameter  $CE$  is perpendicular to the refracting surface at  $A$ , the angle of incidence will be equal to the angle  $AEC$ , and the angle of refraction, or complement to two right ones, will be  $EAT$ ; consequently the sine of incidence is to the sine of refraction, (as the sine of the angle  $AEC$ , to the sine of  $EAT$ , or as  $AT$  to  $TE$ <sup>a</sup>, that is, when  $A$  comes nearest to  $C$ , and so the incident rays are almost perpendicular to the surface,) as  $CT$  to  $TE$ <sup>b</sup>; and disjointly 'the sine of incidence is to the difference of the sines as  $CT$  to  $CE$ . Q. E. D.

<sup>a</sup> Art. 85.

<sup>b</sup> Art. 68.

<sup>c</sup> Euc. V. 17.

93. Corol. 1.  $CT$  is to  $TE$  as the sine of incidence to the sine of refraction.

94. Corol. 2. If this point  $T$  be the focus of incident rays, the refracted rays will go parallel to  $TE$ <sup>d</sup>.

<sup>d</sup> Art. 11.

## PROPOSITION III.

Fig. 89, 90. 95. When parallel rays fall upon a sphere, either denser or rarer than the ambient medium, in the diameter  $CD$  produced, which is parallel to the incident rays, as  $QA$ , let  $T$  be their focus after their first refraction at  $A$ ; and the point  $F$  which bisects  $TD$  shall be their focus after their second refraction at  $D$ .

For let the incident and emergent rays,  $QA$ ,  $FG$  produced, meet in  $H$ , and since the refractions at  $A$  and  $G$  are equal, as appears supposing a ray to go both ways along the chord  $AG$ , the triangle  $AHG$  is equiangular at its base  $AG$ , and therefore equicrural; and so is the similar triangle  $GFT$ , the lines  $AH$ ,  $FT$  being parallel. Therefore when  $A$  approaches toward  $C$  till  $G$  is coincident with  $D$  and the triangle  $GFT$  is vanishing, the leg  $GF$  will become equal to half the base  $GT$ ; that is,  $DF$  will become equal to  $DT$ <sup>f</sup>. Q. E. D.

<sup>e</sup> Euc. I. 6.

<sup>f</sup> Art. 68.

## LEMMA.

Fig. 91. to 94. 96. There is a certain point  $E$  within every double convex or double concave lens, through which every ray that passes, will have its incident and emergent rays parallel.



The lines  $QA$  and  $Aq$ , produced as in the figures, represent an incident and a refracted ray, cutting  $QC$  in any point  $q$  and since a perpendicular to the plane at  $A$  is parallel to the angle  $AQC$  will be equal to the angle of incidence, and the angle of refraction. Therefore since equal angles have the sine of incidence is to the sine of refraction (as the angle  $AQC$  to the sine of  $AqC$ , or as  $Aq$  to  $AQ$ , that <sup>a</sup> Art. 85.  
the ray  $QA$  is almost perpendicular to the plane  $AB$ ) as <sup>b</sup> Art. 68.  
 $Q, E. D.$

1. If the surface  $ACB$  is glass,  $qQ$  is half of  $QC$ , and a <sup>c</sup> Art. 13.  
 $C$ . For  $qC$  is to  $QC$  as 3 to 2<sup>c</sup>; and therefore  $qQ$  is to <sup>d</sup> Euc. V.  
2<sup>d</sup>. In like manner, if  $ACB$  is water,  $qQ$  is a third of <sup>d</sup> Def. 16.  
fourth of  $qC$ .

2. Hence the refractions of a pencil of rays through a <sup>Fig. 83.</sup>  
may be determined. Let  $Q$  be the focus of incident  
its first side  $AB$ , and  $QC$  perpendicular to  $AB$ . To  $QC$   
equal to half  $QC$ ; and  $T$  will be the focus of the rays  $QA$ ,  
after refraction at the surface  $AB$ ; and being also the <sup>c</sup> Art. 88.  
incident rays at  $a$  and  $b$  upon the second surface  $ab$ , from  
a perpendicular to  $ab$ , take away  $Tq$  equal to a third  
 $Cc$ ; and  $q$  will be the focus of the emergent rays  $qa, qb$

3. Hence the focuses of incident and emergent rays at  $a$   
are always very nearly at equal distances from it; provided the  
angles and the refracting angle be but small. For then the  
distances  $TC, Tc$  are nearly equal; and in glass  $QC$  and  $qC$   
are thirds of them respectively.

4. By proceeding in the same manner as in the 2d Co- <sup>Fig. 83, 84.</sup>  
found that, when the planes  $AB, ab$  are parallel, and  $TC$ ,  
side,  $Qq$  is one third of  $Cc$ , the thickness of the glass \*.

The distance  $Cc$ , between the parallel planes, and the ratio between the sines  
of incidence and refraction at the first surface,  $I$  and  $R$ , being given, the distance  $Qq$  may  
be found more readily in all cases by means of the following proposition.

*A pencil of converging or diverging rays passes through two parallel planes, the distance  
between the planes is to the distance between the foci of incident and emergent rays, as the sine  
of incidence on the first surface is to the difference between the sines of incidence and refraction.*

$ab$ , be the refracting planes;  $Q$  the focus of incident rays;  $QC$  a perpen- <sup>Fig. 84.</sup>  
dicular to the planes;  $T$  the point in this perpendicular where the directions of the re-  
fracted rays  $Aa, Bb$ , meet it;  $q$  the focus of emergent rays. And  $QA$  being parallel  
to the triangle  $Tqa$ ,

and inver.  $TQ : Qq :: TA : Aa$ , that is, as  $TC : Cc$  <sup>a</sup> Art. 51.  
 $TC : TQ :: Cc : Qq$  <sup>b</sup> Euc. VI. 2.  
 $TC : QC :: I : R$  <sup>c</sup> Art. 87.  
inver.  $TC : TQ :: I : I - R$   
 $Cc : Qq :: I : I - R. \quad Q, E. D.$

## PROPOSITION II.

Fig. 85. to 88. 92. Let  $ACB$  represent a refracting spherical surface whose center is  $E$ , and let the incident rays as  $DA$  come parallel to any semidiameter  $CE$ , in which produced forward or backward, according as the denser medium is convex or concave, take  $CT$  to  $CE$  as the sine of incidence to the difference of the sines, and  $T$  will be the focus of the refracted rays.

For let the refracted ray  $AT$  (produced) cut the semidiameter  $CE$  produced, in any point  $T$  whatever; and since the semidiameter  $EA$  is perpendicular to the refracting surface at  $A$ , the angle of incidence will be equal to the angle  $AEC$ , and the angle of refraction, or its complement to two right ones, will be  $EAT$ ; consequently the sine of incidence is to the sine of refraction, (as the sine of the angle  $AEC$ , to the sine of  $EAT$ , or as  $AT$  to  $TE$ <sup>a</sup>, that is, when  $A$  comes nearest to  $C$ , and so the incident rays are almost perpendicular to the surface,) as  $CT$  to  $TE$ <sup>b</sup>; and disjointly the sine of incidence is to the difference of the sines as  $CT$  to  $CE$ . Q. E. D.

93. Corol. 1.  $CT$  is to  $TE$  as the sine of incidence to the sine of refraction.

94. Corol. 2. If this point  $T$  be the focus of incident rays, the refracted rays will go parallel to  $TE$ <sup>d</sup>.

## PROPOSITION III.

Fig. 89, 90. 95. When parallel rays fall upon a sphere, either denser or rarer than the ambient medium, in the diameter  $CD$  produced, which is parallel to the incident rays, as  $QA$ , let  $T$  be their focus after their first refraction at  $AC$ ; and the point  $F$  which bisects  $TD$  shall be their focus after their second refraction at  $DG$ .

For let the incident and emergent rays,  $QA$ ,  $FG$  produced, meet in  $H$ , and since the refractions at  $A$  and  $G$  are equal, as appears by supposing a ray to go both ways along the chord  $AG$ , the triangle  $AHG$  is equiangular at its base  $AG$ , and therefore equicrural<sup>e</sup>; and so is the similar triangle  $GFT$ , the lines  $AH$ ,  $FT$  being parallel. Therefore when  $A$  approaches toward  $C$  till  $G$  is coinciding with  $D$  and the triangle  $GFT$  is vanishing, the leg  $GF$  will become equal to half the base  $GT$ ; that is,  $DF$  will become equal to half  $DT$ <sup>f</sup>. Q. E. D.

## LEMMA.

Fig. 91. to 94. 96. There is a certain point  $E$  within every double convex or double concave lens, through which every ray that passes, will have its incident and emergent



emergent parts  $QA, aq$  parallel to each other: but in a plano-convex or plano-concave lens that point  $E$  is removed to the vertex of the concave or convex surface; and in a meniscus and in that other concavo-convex lens, it is removed a little way out of them, and lyes next to the surface which has the greatest curvity.

For let  $REr$  be the axis of the lens joining the centers  $R, r$  of its surfaces  $A, a$ . Draw any two of their semidiameters  $RA, ra$  parallel to each other, and join the points  $A, a$ , and the line  $Aa$  will cut the axis in the point  $E$  above described. For the triangles  $REA, rEa$  being equiangular,  $RE$  will be to  $Er$  in the given ratio of the semidiameters  $RA, ra$ ; and consequently the point  $E$  is invariable in the same lens. Now supposing a ray to pass both ways along the line  $Aa$ , it being equally inclined to the perpendiculars to the surfaces, will be equally bent and contrary ways in going out of the lens; so that its emergent parts  $AQ, aq$  will be parallel. Now any of these lenses will become plano-convex or plano-concave, by conceiving one of the semidiameters  $RA, ra$  to become infinite, and consequently to become parallel to the axis of the lens, and then the other semidiameter will coincide with the axis; and so the points  $A, E$  or  $a, E$  will coincide. *Q. E. D.*

97. *Corol.* Hence when a pencil of rays falls almost perpendicularly upon any lens, whose thickness is inconsiderable, the course of the ray which passes through  $E$ , above described, may be taken for a straight line passing through the center of the lens, without sensible error in sensible things. For it is manifest from the length of  $Aa$  and from the quantity of the refractions at its extremities, that the perpendicular distance of  $AQ, aq$  when produced, will be diminished both as the thickness of the lens and the obliquity of the ray is diminished.

DEFINITION. The point  $E$  is called the center of the lens.

#### PROPOSITION IV.

98. *To find the focus of parallel rays falling almost perpendicularly upon any given lens.* Fig. 95. to 100.

Let  $E$  be the center of the lens,  $R$  and  $r$  the centers of its surfaces,  $Rr$  its axis,  $gEG$  a line parallel to the incident rays upon the surface  $B$ , whose center is  $R$ . Parallel to  $gE$  draw a semidiameter  $BR$ , in which produced let  $V$  be the focus of the rays after their first refraction at the surface  $B$ , and joining  $Vr$  let it cut  $gE$  produced in  $G$ , and  $G$  will be the focus of the rays that emerge from the lens.

For since  $V$  is also the focus of the rays incident upon the second

F

surface



<sup>a</sup> Art. 97.

surface  $A$ , the emergent rays must have their focus in some point of that ray which passes straight through this surface; that is in the line  $Vr$ , drawn through its center  $r$ : and since the whole course of another ray is reckoned a straight line  $gEG^a$ , its intersection  $G$  with  $Vr$  determines the focus of them all. *Q. E. D.*

<sup>b</sup> Art. 92.

99. *Corol. 1.* When the incident rays are parallel to the axis  $rR$ , the focal distance  $EF$  is equal to  $EG$ . For let the incident rays that were parallel to  $gE$  be gradually more inclined to the axis till they become parallel to it; and their first and second focuses  $V$  and  $G$  will describe circular arches  $VT$  and  $GF$  whose centers are  $R$  and  $E$ . For the line  $RV$  is invariable; being in proportion to  $RB$  in a given ratio of the lesser of the sines of incidence and refraction to their difference<sup>b</sup>; consequently the line  $EG$  is also invariable, being in proportion to the given line  $RV$  in the given ratio of  $rE$  to  $rR$ , because the triangles  $EGr$ ,  $RVr$  are equiangular.

100. *Corol. 2.* The last proportion gives the following rule for finding the focal distance of any thin lens. As  $Rr$ , the interval between the centers of the surfaces, is to  $rE$ , the semidiameter of the second surface, so is  $RV$  or  $RT$ , the continuation of the first semidiameter to the first focus, to  $EG$  or  $EF$ , the focal distance of the lens. Which according as the lens is thicker or thinner in the middle than at its edges, must lye on the same side as the emergent rays or the opposite side.

<sup>c</sup> Art. 92.

101. *Corol. 3.* Hence when rays fall parallel on both sides of any lens, the focal distances  $EF$ ,  $Ef$  are equal. For let  $rt$  be the continuation of the semidiameter  $Er$  to the first focus  $t$  of rays falling parallel upon the surface  $A$ ; and the same rule that gave  $rR$  to  $rE$  as  $RT$  to  $EF$ , gives also  $rR$  to  $RE$  as  $rt$  to  $Ef$ . Whence  $Ef$  and  $EF$  are equal, because the rectangles under  $rE$ ,  $RT$  and also under  $RE$ ,  $rt$  are equal. For  $rE$  is to  $rt$  and also  $RE$  to  $RT$  in the same given ratio<sup>c</sup>.

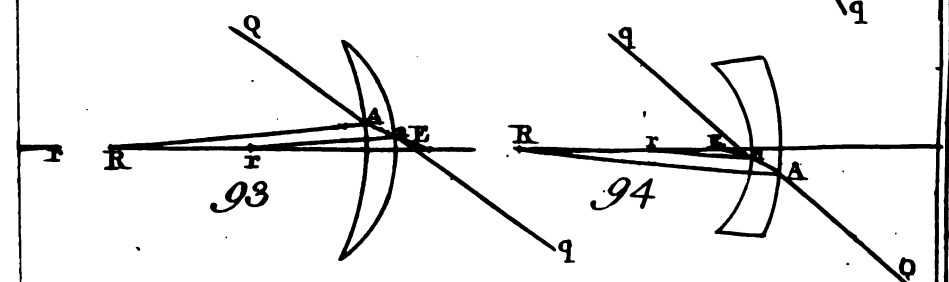
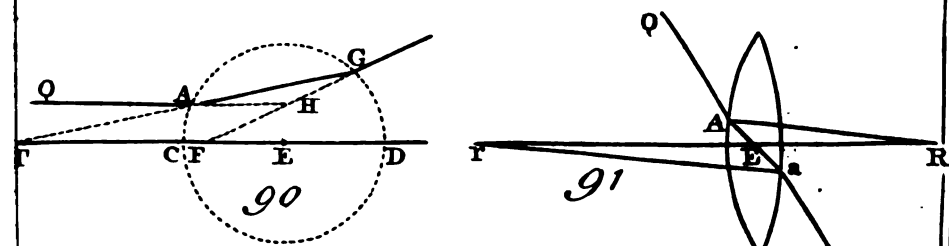
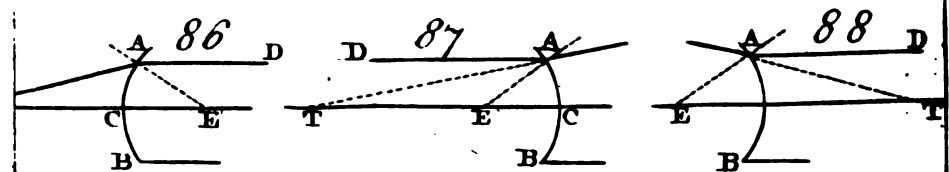
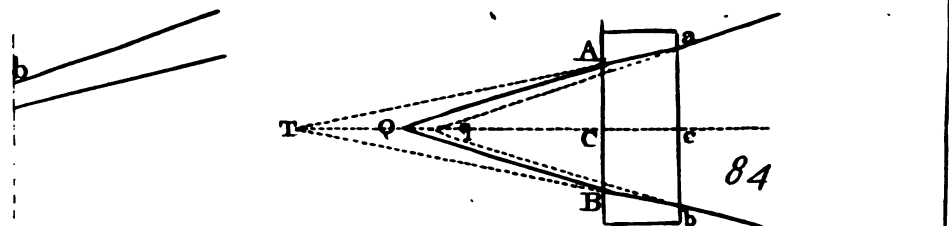
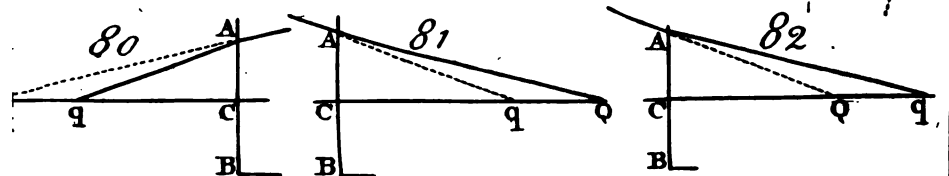
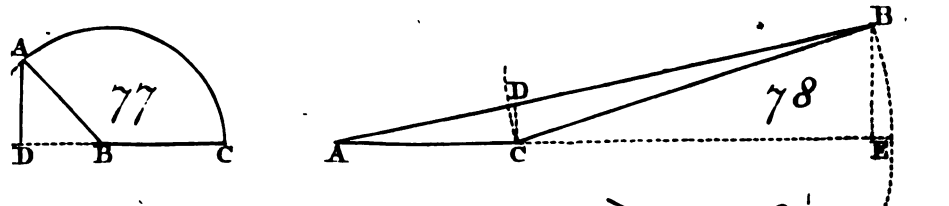
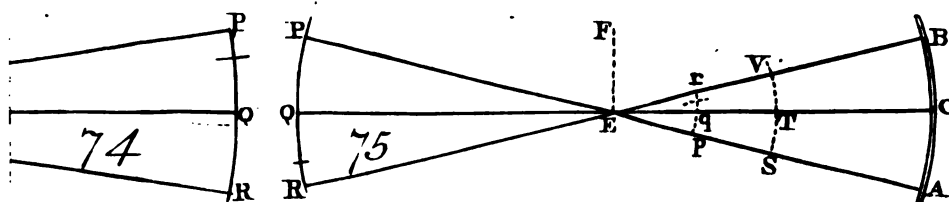
<sup>d</sup> Art. 93. 13.

102. *Corol. 4.* Hence in particular in a double-convex or double-concave lens made of glass, it is as the sum of their semidiameters (or in a meniscus as their difference) to either of them, so is double the other, to the focal distance of the glass. For the continuations  $RT$ ,  $rt$  are severally double their semidiameters: because in glass  $ET$  is to  $TR$  and also  $Et$  to  $tr$  as 3 to 2<sup>d</sup>.

103. *Corol. 5.* Hence if the semidiameters of the surfaces of the glass be equal, its focal distance is equal to one of them; and is equal to the focal distance of a plano-convex or plano-concave glass whose semidiameter is as short again. For considering the plane surface as having an infinite semidiameter, the first ratio of the last mentioned proportion may be reckoned a ratio of equality.

P R O -









## PROPOSITION V.

104. *The focus of incident rays upon a single surface, sphere or lens being given, it is required to find the focus of the emergent rays.* Fig. 101. to 106.

Let any point  $Q$  be the focus of incident rays upon a spherical surface, lens, or sphere, whose center is  $E$ ; and let other rays come parallel to the line  $QEq$  the contrary way to the given rays, and after refraction let them belong to a focus  $F$ ; then taking  $Ef$  equal to  $EF$  in the lens or sphere, but equal to  $CF$  in the single surface, say as  $QF$  to  $FE$  so  $Ef$  to  $fq$ , and placing  $fq$  the contrary way from  $f$  to that of  $FQ$  from  $F$ , the point  $q$  will be the focus of the refracted rays, without sensible error; provided the point  $Q$  be not so remote from the axis, nor the surfaces so broad as to cause any of the rays to fall too obliquely upon them.

For with the center  $E$  and semidiameters  $EF$  and  $Ef$  describe two arches  $FG$ ,  $fg$  cutting any ray  $QAaq$  in  $G$  and  $g$ , and draw  $EG$  and  $Eg$ . Then supposing  $G$  to be a focus of incident rays, (as  $GA$ ,) the emergent rays (as  $agq$ ) will be parallel to  $GE$ <sup>a</sup>; and on the other hand supposing  $g$  another focus of incident rays (as  $ga$ ,) the emergent rays (as  $AGQ$ ,) will be parallel to  $gE$ . Therefore the triangles  $QGE$ ,  $Egq$  are equiangular, and consequently  $QG$  is to  $GE$  as  $Eg$  to  $gq$ ; that is, when the ray  $QAaq$  is the nearest to  $QEq$ ,  $QF$  is to  $FE$  as  $Ef$  to  $fq$ <sup>b</sup>. Now when  $Q$  accedes to  $F$  and coincides with it, the emergent rays become parallel, that is  $q$  recedes to an infinite distance; and consequently when  $Q$  passes to the other side of  $F$ , the focus  $q$  will also pass through an infinite space from one side of  $f$  to the other side of it.  $Q. E. D.$  Art. 94. 99. 101.

105. *Corol. 1.* In refractions at a spherical surface  $AC$ , the focus  $q$  may also be found by this rule, as  $QF$  to  $FC$  so  $Cf$  to  $fq$ ; because  $FC$  and  $Ef$  and also  $FE$  and  $Cf$  are equal<sup>c</sup>.

106. *Corol. 2.* It may also be found by this rule, as  $QF$  to  $QE$  so  $QC$  to  $Qq$ ; placing  $Qq$  so that all the four distances from  $Q$  may lye on one side of it, or else two on each. For the triangles  $QGE$ ,  $QAq$  being equiangular we have  $QG$  to  $QE$  as  $QA$  to  $Qq$ .

107. *Corol. 3.* In a sphere or lens the focus  $q$  may be found by this rule, as  $QF$  to  $QE$  so  $QE$  to  $Qq$ , to be placed the same way from  $Q$  as  $QF$  lyes from  $Q$ . For let the incident and emergent ray  $QA$ ,  $qa$  be produced till they meet in  $e$ ; and the triangles  $QGE$ ,  $Qeq$  being equiangular, we have  $QG$  to  $QE$  as  $Qe$  to  $Qq$ ; and when the angles of these triangles are vanishing, the point  $e$  will coincide with  $E$ ; because in the sphere the triangle  $Aea$  is equiangular at the base  $Aa$ , and consequently  $Ae$  and  $ae$  will at last be-



come semidiameters of the sphere. In a lens the thickness  $Aa$  is inconsiderable.

108. *Corol. 4.* In all cases the distance  $f q$  varies reciprocally as  $F Q$  does; and they lye contrary ways from  $f$  and  $F$ ; because the rectangle or the square under  $EF$  and  $Ef$ , the middle terms in the foregoing proportions, is invariable. Hence if either of the corresponding focuses  $Q, q$  be put in motion along the axis of the pencil, the other focus will move the same way: and therefore if these focuses be on contrary sides of the glass, while one moves towards it the other will move from it; but if they be both on the same side of the glass, they will both move from it or both towards it; and will come nearer to each other as they come nearer to the glass, till when one coincides with its surface the other will do so too, in the single surface accurately<sup>a</sup>, and in the lens very nearly, provided the glass be very thin, and the distance of the ray from its axis be very small<sup>b</sup>. But these focuses cannot coincide at the surface of a globe; for  $QF$  and  $QE$  being finite,  $Qq$  cannot vanish<sup>c</sup>. They will coincide at the center of the single surface, because the rays fall perpendicularly, and therefore pass through without suffering any refraction<sup>d</sup>.

<sup>a</sup> Art. 105.

<sup>b</sup> Art. 104.

<sup>c</sup> Art. 107.

<sup>d</sup> Art. 15.

109. *Corol. 5.* Convex lenses of different shapes that have equal focal distances, when put into each others places, have equal powers upon any pencil of rays to refract them to the same focus. Because the rules abovementioned depend only upon the focal distance of the lens, and not upon the proportion of the semidiameters of its surfaces.

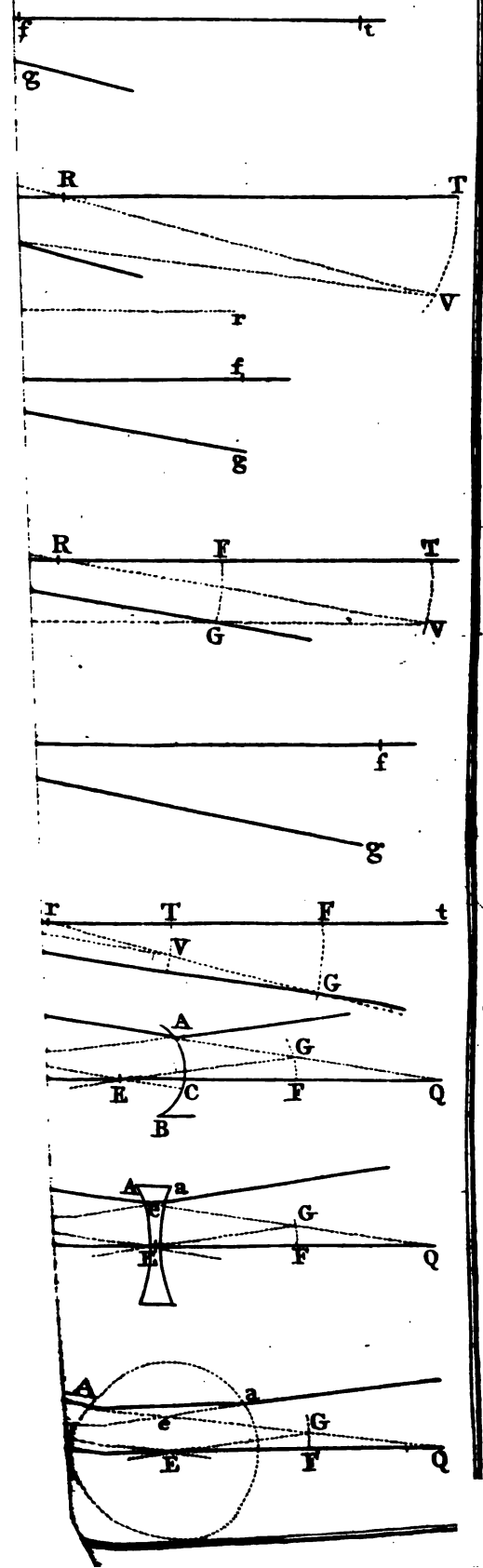
110. *Corol. 6.* The rule that was given for a sphere of an uniform density, will serve also for finding the focus of a pencil of rays refracted through any number of concentrick surfaces, which intercede uniform mediums of any different densities. For when rays come parallel to any line drawn through the common center of these mediums, and are refracted through them all, the distance of their focus from that center is invariable, as in an uniform sphere.

111. *Corol. 7.* When the focuses  $Q, q$  lye on the same side of the refracting surfaces, if the incident rays flow from  $Q$ , the refracted rays will also flow from  $q$ ; and if the incident rays flow towards  $Q$ , the refracted will also flow towards  $q$ ; and the contrary will happen when  $Q$  and  $q$  are on contrary sides of the refracting surfaces. Because the rays are continually going forwards.

The 75th and 76th articles are applicable to refractions as well as reflections.

“ After









“ After Des Cartes had discovered the true law of refraction<sup>a</sup>, it<sup>a</sup> Art. 13.  
 “ presently suggested to him the impossibility of collecting all the  
 “ rays of the same pencil into a single point by means of spherical  
 “ surfaces only. And not being aware that the aberration of rays  
 “ from the geometrical foci of the optical glasses then in use was  
 “ owing to any other cause than the *sphericity* of their figures, he  
 “ immediately applied himself to investigate what curves could  
 “ make all the rays belonging to the same pencil converge after re-  
 “ fraction to a single point. The seeming importance of this pro-  
 “ blem induced other mathematicians after him to prosecute the  
 “ solution of it: and among other surfaces, it was found that those,  
 “ which are generated by the revolution of *conic sections* round their  
 “ axes, would produce the effect. The limits of an abstract will  
 “ not permit us to relate all their discoveries upon this subject; in  
 “ the following problems and their corollaries are contained the  
 “ principal discoveries of Des Cartes and his demonstrations of the  
 “ more difficult cases. \*

## “ PROBLEM I.

“ *To collect every ray belonging to a pencil of parallel rays into one*  
 “ *point, when they are incident upon the refracting surface of a denser*  
 “ *medium.*

“ Take  $DK$  to  $HI$  as the sine of incidence to the sine of refraction Fig. 107.  
 “ tion; with the axis major  $DK$  and foci  $H, I$ , describe an ellipse  
 “  $DBK$ ; describe also a spheroid by the revolution of this ellipse  
 “ round its transverse axis: and all the rays, which are incident  
 “ upon the convex surface  $DBb$ , parallel to  $DK$ , will converge to  
 “ the farther focus  $I$ .

## “ PROBLEM II.

“ *To collect every ray belonging to a pencil of parallel rays into one*  
 “ *point, when they are incident upon the refracting surface of a rarer*  
 “ *medium.*

“ Take  $DK$  to  $HI$  as the sine of incidence to the sine of refraction Fig. 108.  
 “ tion; with the transverse axis  $DK$  and foci  $H, I$ , describe an hyper-  
 “ bola; describe also an hyperboloid by the revolution of this hy-  
 “ perbola round its transverse axis: and all the rays, which are inci-  
 “ dent upon the concave surface  $BDY$ , parallel to  $DK$ , will con-  
 “ verge to the farther focus  $I$ .

“ DE-

\* Des Cartes's Dioptr. Ch. 8.

Fig 107, 108. " DEMONSTRATION. Let  $AB, MD$  be two parallel rays, of which  $MD$  is perpendicular to the surface: join  $HB, IB$ ; through  $B$  draw two straight lines  $CBE, LBG$ , which are at right angles to each other, and one of which  $CBE$  is a tangent to the curve at the point of incidence  $B$ : through  $H$  draw  $HO$  parallel to  $LG$ , meeting  $IB$  in the hyperbola, or  $IB$  produced in the ellipse, in  $O$ , and the tangent in  $C$ . And the angles  $OBC, CBH$  being equal<sup>a</sup>, and  $HO$  cutting  $BC$  at right angles, the triangles  $OBC, CBH$  are equal; and therefore  $OI$  is equal to the sum or difference of  $HB$  and  $BI$ ; that is,  $OI$  equals  $DK$ <sup>b</sup>. Now  $LG$  being perpendicular to the surface, the angle  $ABN$  is either the angle of incidence or its complement to two right angles; and therefore the sine of the angle  $BNI$ , which is alternate to  $ABN$ , is equal to the sine of incidence<sup>c</sup>. And the sine of  $BNI$  is to the sine of  $NBI$  as  $BI$  to  $NI$ <sup>d</sup>; that is, as  $OI$  to  $HI$ <sup>e</sup>, or as  $DK$  to  $HI$ ; that is, as the sine of incidence to the sine of refraction<sup>f</sup>. Wherefore the sine of  $NBI$  must be equal to the sine of refraction, and  $BI$  is the course of the refracted ray.  $\mathcal{Q}. E. D.$

<sup>a</sup> Hamil.  
Con. Sect.  
B. II. p. 16, 17.  
<sup>b</sup> Hamil.  
Con. Sect.  
B. II. p. 14.

<sup>c</sup> Art. 83.  
<sup>d</sup> Art. 85.  
<sup>e</sup> Euc. VI. 2.  
<sup>f</sup> By construction.

Fig 109. " *Corol. 1.* If, therefore, with the center  $I$  and any radius  $IB$ , less than  $ID$ , a circular arch  $B\mathcal{Q}$  be described, and the figure  $B\mathcal{Q}D$  revolve round its axis  $D\mathcal{Q}$ , it will determine the shape of a piece of glass, which, placed in air, will collect into one point  $I$  all the rays that fall upon its convex surface parallel to  $DI$ . For the first surface  $BDB$  of this solid refracts them all to  $I$ ; and the radii of a circle being perpendicular to its circumference, the course of these rays will not be changed by the second surface  $B\mathcal{Q}B$ .

Fig 110. " Likewise, if from any point  $B$  in the hyperbola the right line  $B\mathcal{Q}$  be drawn perpendicular to the axis, the figure  $B\mathcal{Q}D$  by revolving round  $D\mathcal{Q}$  will generate a solid piece of glass, which, placed in air, will have the same effect upon all rays that are incident upon its plane surface parallel to the transverse axis.

" *Corol. 2.* Each of the glasses described in the preceding corollary will make all the rays, which are incident diverging from the same point, move parallel after refraction. For, if  $I$  be made the focus of incident rays, the refracted rays will be parallel to  $ID$  by article. 11.

Fig 111. " *Corol. 3.* Let an ellipse be constructed as before; and draw the circular arch  $RO$  with the center  $I$  and any radius  $RI$  greater than  $DI$ ;



“ *DI*; assume also any point *B* in the ellipse, which is not farther  
 “ from *D* than from *K*, draw the right line *OB* tending to *I*; and  
 “ the glass, which is generated by the revolution of the figure  
 “ *BDRO* about its axis *DR*, will make all rays, that are incident  
 “ upon the concave side parallel to its axis, diverge from the point  
 “ *I* after refraction. For the mediums continuing the same, it is  
 “ evident that the ray *PB* is as much bent out of its course by the  
 “ concave side of the refracter *DB*, as the ray *AB* is by the convex  
 “ side of the same refracter”. Since therefore *AB* and *PB* are in <sup>Art. 11. 17.</sup>  
 “ the same straight line, and *AB* moves in the line *BI* after refraction,  
 “ *BO*, which is in the same direction with *BI*, must be the  
 “ course of the ray *PB* after refraction.

“ Parallel rays may also be refracted diverging accurately from  
 “ the same point by means of a glass constructed in the following  
 “ manner. Describe an hyperbola *BDB* similar to that which is <sup>Fig. 112.</sup>  
 “ described in the second problem; assume any point *O*, from which  
 “ two straight lines can be drawn in such a manner that one of them  
 “ *OB* shall be parallel to the axis and cut the curve in some point *B*,  
 “ and that the other *OR* shall be perpendicular to the axis and not  
 “ cut the curve. Let the figure *OBDR* be turned round its axis,  
 “ and it will generate the solid required.

“ *Corol. 4.* Each of the glasses constructed in the preceding corollary  
 “ lary will make all rays, that are incident converging to the same  
 “ point, move parallel after refraction; for the reason given in  
 “ corol. 2d.

“ *Corol. 5.* If two glasses be constructed, each of which is similar <sup>Fig. 113.</sup>  
 “ to the elliptic refracter described in the first corollary, and be placed  
 “ in air, so that their axes are in the same straight line, and that their  
 “ convex surfaces are turned towards each other, all the rays, which  
 “ are incident diverging from the same point *I*, will after refraction  
 “ converge to the same point *i*.

“ Or two hyperbolic refracters, each of which is similar to the <sup>Fig. 114.</sup>  
 “ glass described in the first corollary, and placed in such a manner  
 “ that their axes are parallel, and that their plane surfaces are turned  
 “ towards each other, will produce the same effect. And since the <sup>Fig. 115.</sup>  
 “ refractions at the curved surfaces will be the same, when the plane  
 “ surfaces coincide as when they are asunder, a single convex glass,  
 “ which is formed by turning the hyperbolic arches *DB*, *db* round  
 “ their axes, will collect into one point *i* all the rays which diverge  
 “ from *I*.

“ *Corol.*

Fig. 116, 117. " *Corol. 6.* If it be required to make a pencil of rays, which are  
 " incident diverging from a single point *I*, diverge after refraction  
 " from a different point *i*, use two elliptic glasses, the figures and  
 " situations of which are represented in fig. 116, or the single hy-  
 " perbolic refracter, which is represented in fig. 117. The same  
 " glasses will also make converging rays converge to a different point  
 " after refraction.

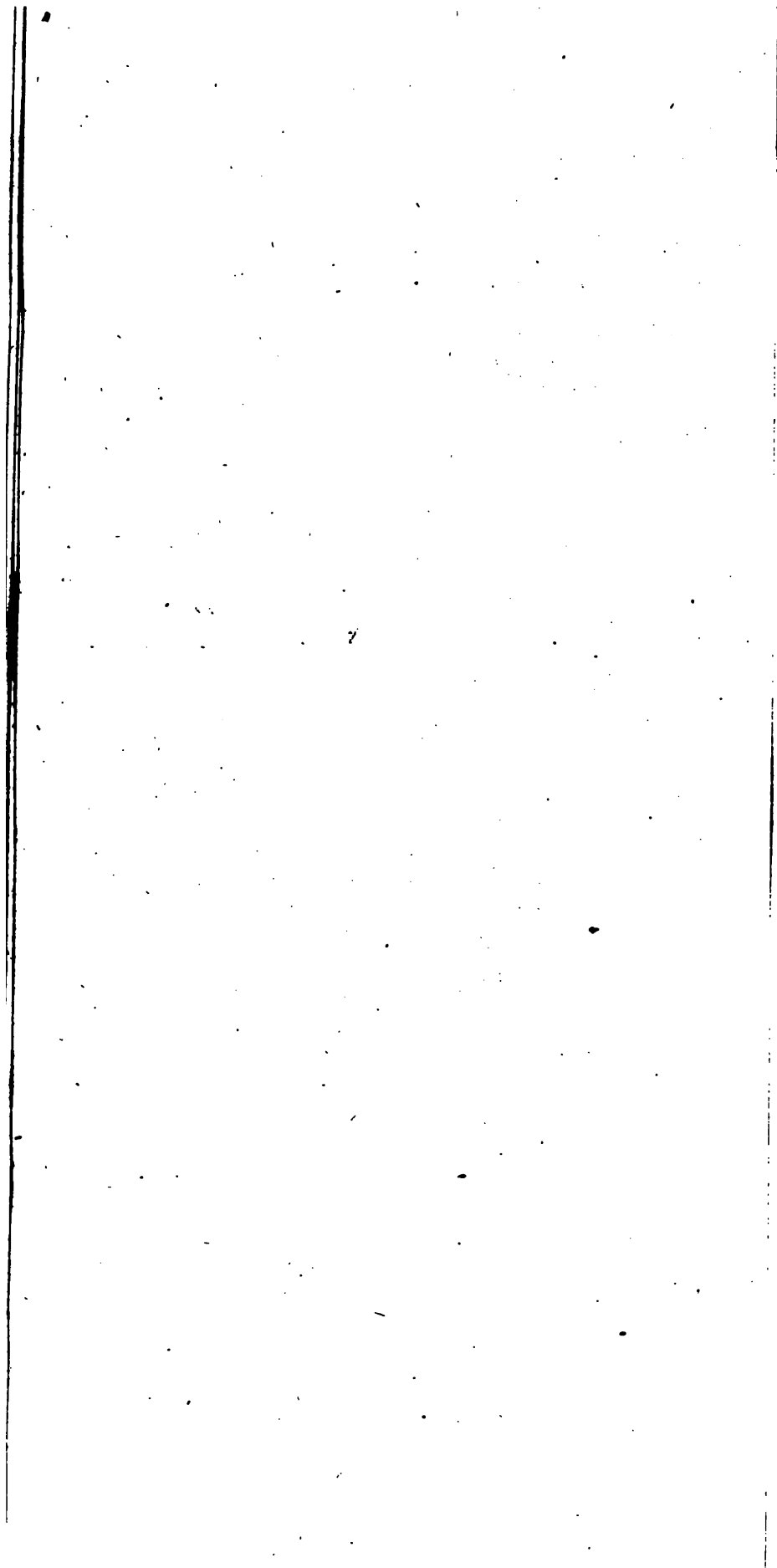
Fig. 118, 119. " *Corol. 7.* Lastly, a pencil of converging rays will be made to  
 " diverge accurately from a single point after refraction, either by  
 " means of two elliptic glasses, which are represented in fig. 118, or  
 " by the single hyperbolic glass represented in fig. 119.

" This apparent superiority of conical surfaces induced not only  
 " the mechanicks but the most eminent mathematicians to contrive  
 " engines for grinding and polishing these surfaces; till Sir Isaac  
 " Art. 28. " Newton discovered the different refrangibility of the rays of light <sup>a</sup>.  
 " He then saw that, however accurately rays of the same colour  
 " might be collected into a single point, rays of different colours,  
 " though they belong to the same incident beam, must move towards  
 " different focuses after refraction; and proved that the aberration  
 " arising from the unequal refrangibility of different kinds of light  
 " was 5449 times greater than the aberration, which is caused by the  
 " Art. 220. " spherical figure of a glass <sup>b</sup>. He observed besides, that, though a  
 " conical lens could refract such rays as are parallel to its axis more  
 " nearly to a single point than a spherical surface, yet the latter  
 " would be preferable to the former in respect of rays, which fall  
 " with some degree of obliquity; since the curvature of a sphere is  
 " uniform in every part, whereas different parts of an ellipse or hy-  
 " perbola have a different curvature. For these reasons and on ac-  
 " count of the insuperable difficulty that attended the construction  
 " of conical surfaces, all further attempts to correct the aberrations  
 " in refractions by means of such surfaces have long been discon-  
 " tinued."



Note

LB





## CHAP. VIII.

TO DETERMINE THE PLACE, MAGNITUDE AND SITUATION OF IMAGES FORMED BY REFRACTED RAYS.

## PROPOSITION I.

112. **I**MAGES formed by refractions at plane surfaces are similar to the objects, and are always erect, or in a similar situation to the objects, and on the same side of the planes\*.

Let  $PQR$  be an object radiating upon a refracting plane  $ACB$ ; Fig. 120, 121. to which draw the perpendiculars  $PA$ ,  $QC$ ,  $RB$ , &c. and in these perpendiculars take  $Ap$  to  $AP$ ,  $Cq$  to  $CQ$ ,  $Br$  to  $BR$  in the given ratio of the sine of incidence to the sine of refraction<sup>a</sup>; and the focuses  $p$ ,  $q$ ,  $r$ , &c. will constitute a similar image in a similar situation to the object; the parts  $pq$ ,  $qr$  being in the same ratio as  $PQ$ ,  $QR$ . This is self-evident when the object is parallel to the refracting plane; and when it is inclined, produce it till it cuts the plane in  $D$ ; and the produced image will also cut it in the same point  $D$ . For supposing the perpendicular  $BrR$  to move towards  $D$ , the lines  $BR$ ,  $Br$  being in a given ratio will vanish both together: and because the triangle  $pDP$  is cut by parallel lines  $qQ$ ,  $rR$ , it will be as  $pq$  to  $PQ$  (so  $qD$  to  $QD$ ) so  $qr$  to  $QR$ <sup>b</sup>; and alternately  $pq$  to  $qr$  as  $PQ$  to  $QR$ . In like manner if the rays that belong to the focuses  $p$ ,  $q$ ,  $r$  be refracted again by another plane, either parallel or inclined to  $AB$ , their second focuses will constitute a second image, similar to the first image, and consequently similar to the object; and so on.  $Q. E. D.$

## PROPOSITION II.

113. An image  $pqr$ , formed by a flat piece of glass  $ABba$  is upright, parallel and equal to the object  $PQR$ , and lies on the same side of the glass as the object; but nearer to it by a third part of the thickness of the glass.

Let  $PA$ ,  $QC$ ,  $RB$  be drawn perpendiculars to the surface  $ACB$ , Fig. 122, 123. and the focuses  $p$ ,  $q$ ,  $r$  of the several pencils that flow from  $P$ ,  $Q$ ,  $R$ ,

\* This proposition is only true, when the object is a straight line or a plane superficies.



- <sup>a</sup> Art. 87. lying in these perpendiculars<sup>a</sup>, the image must be upright. But we have also shewn, that each of the focuses  $p, q, r$ , lyes nearer to  $ACB$ , than the points  $P, Q, R$ , by a third part of  $Cc$ <sup>b</sup>; therefore the image must be so much nearer than the object, and parallel to it.  $Q. E. D.$
- <sup>b</sup> Art. 91.

## PROPOSITION III.

114. *An image formed by a prism is always upright, and equal to the object, and lies on the same side of the prism, and at the same distance from it as the object itself; provided the refracting angle of the prism, and the refractions made by it, be but small.*

- Fig. 124. Take two rays,  $PE, QE$ , which coming from the extremities of the object, pass through a point  $E$ , so near to the angular point of the refracting angle, that the distances between their points of incidence and emergence need not be mentioned. And since the total bendings of the rays  $PEN, QEO$  are equal<sup>c</sup>, they will cross each other, so as that the angle  $PEQ$  will be equal to the angle  $NEO$  or to  $pEq$ , made by the emergent rays produced backwards: and because the distance  $Ep$  of the focus  $p$ , of the pencil that flowed from  $P$ , is equal to  $EP$ <sup>d</sup>, and in like manner the focal distance  $Eq$  equal to  $EQ$ ; the image  $pq$  will be upright and equal to the object, and at an equal distance on the same side of the prism.  $Q. E. D.$
- <sup>c</sup> Art. 55.
- <sup>d</sup> Art. 90.

## PROPOSITION IV.

- Fig. 125. to 128. 115. *If an arch of a circle  $PQR$  described upon the center  $E$ , of a spherical surface, sphere or lens, be considered as an object, its image  $pqr$  will be a similar concentrick arch; whose length will be to the length of the object in the ratio of their distances from the common center  $E$ ; and the image will be erect or inverted, with respect to the object, according as they lye on the same side of the center or on contrary sides.*

- The proposition is evident by inspection of the 125th figure in all cases of refractions made by concentrick surfaces; because the parts of these surfaces are alike exposed to the parts of the concentrick object. And in a lens the focuses of all the pencils of parallel rays lye also in a concentrick arch  $GfH$ ; whence  $Pp$  and  $Qq$  being third proportionals to two pair of equal distances  $PG$  and  $PE, QF$  and  $QE$ <sup>e</sup>, are also equal; and so the image  $pqr$  is also a concentrick arch. Now since the axes of the pencils are reckoned straight lines passing through  $E$ <sup>f</sup>, the angles  $pEr, PER$  are equal; and therefore the ratio of the image to the object, is the same as of their distances from
- <sup>e</sup> Art. 107.
- <sup>f</sup> Art. 97.



from the center  $E$ . And according as their corresponding extremities  $P, p$  are on the same or on contrary sides of  $E$ , they lye on the same or contrary sides of their middle points  $Q, q$ ; that is, the image is accordingly erect or inverted.  $Q, E, D$ .

116. *Corol.* The smaller the circular object is with respect to its radius or distance from the center  $E$ , the nearer it approaches in shape to a straight line, and so does its similar image. Consequently a small straight object placed at a good distance from the center of the glass may be reckoned to have a straight image very nearly\*: Art. 68. though in strictness of geometry it is an arch of a conic section. And by these propositions the images of all objects may be determined, by finding the images of their out-lines.

## CHAP. IX.

## CONCERNING THE EYE AND MANNER OF VISION.

117. **C**ONSIDERING what has been said in the 92d and 115th articles, one might contrive a tolerable eye in this manner, by placing a pellucid hemisphere  $ABC$  to serve for the fore part, and another concentrick one  $DqE$ , opposite to the former, to serve for its bottom or back part; making the semidiameter,  $Oq$ , of the latter triple the semidiameter,  $OB$ , of the former; and then by filling the whole cavity of both with water. By this means rays of light flowing from the points  $P, Q, R$ , &c. of remote objects, after refraction at the surface  $ABC$ , will be collected to as many other points  $p, q, r$ , of the cavity  $DqE$ , and paint an image upon it. And because a spherical surface does not accurately refract all the rays of a large pencil to a single point<sup>a</sup>, but only those that go pretty near its axis; this imperfection might be remedied by covering the base  $AC$ , of the lesser hemisphere, all but a moderate hole about the center  $O$ ; which would answer the purpose much better than if the surface itself was covered, all but a hole in the middle about  $B$ . For in this latter case the surface  $ABC$  would not receive rays from the lateral points  $P, R$ , so directly as those from the middle of the object, to all which it is exposed alike when the hole is left open at the center  $O$ .

118. Though this construction of the eye appears not amiss at first sight, yet we shall see presently that the author of nature has

<sup>a</sup> A fictitious eye described by *Hugens* \*. Fig. 129.

<sup>a</sup> See the latter part of chap. 7.

\* *Opuscula Posthuma*. p. 112.

wisely varied some things for the better, and added others absolutely necessary; though in every thing we cannot perceive his designs. In the first place he would not make use of an entire hemisphere *ABC*, but retaining the middle part, has taken off pretty much from the sides, and yet without contracting the compass of objects taken in at one view. The reason of this was to bend inwards the edges of the larger hemisphere about *D* and *E*, thereby reducing the shape of the eye to a rounder figure, for the convenience of its motion every way in the cavity that contains it. He has therefore given it such a shape, as is expressed in this other figure, representing an human eye dissected through its axis, all the parts being twice as big as in the life to render them more conspicuous.

Fig. 130.

An human  
eye describ-  
ed.

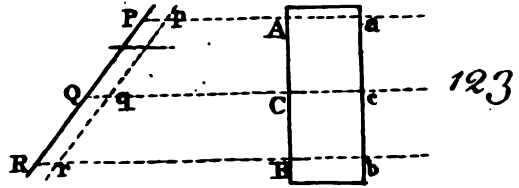
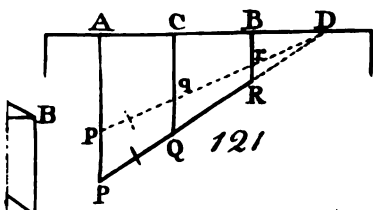
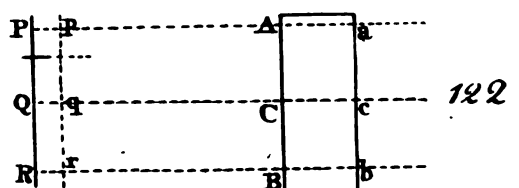
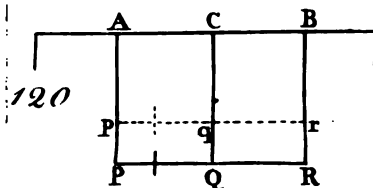
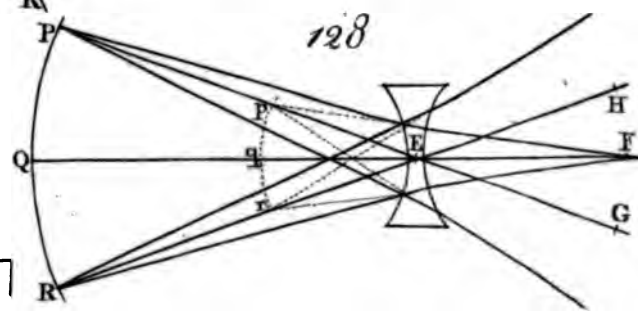
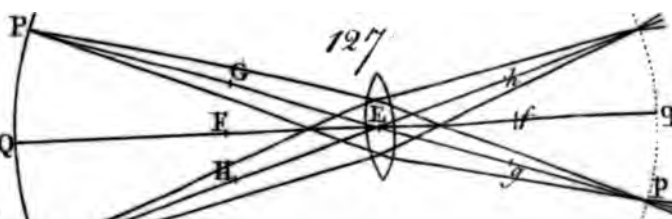
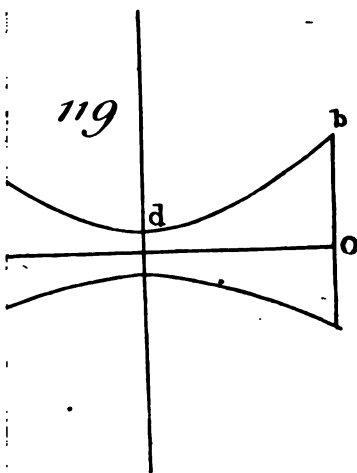
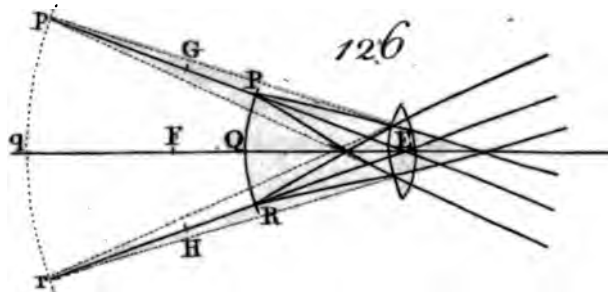
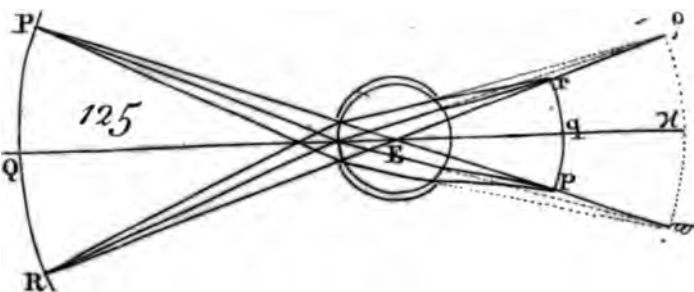
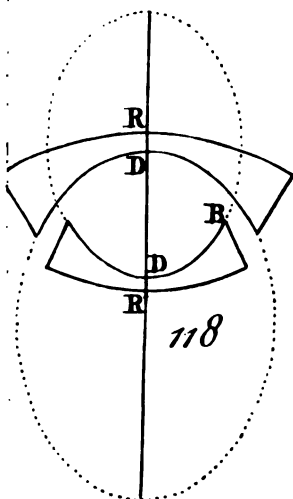
119. Here the transparent parts of the coat called the cornea is *ABC*; the remainder *ATTC* being opaque, and a portion of a larger sphere. Within this outward coat anatomists distinguish two others; the innermost of which is called the retina, being like a fine net composed of the fibres of the optick nerve *YVT* woven together, and is white about the parts *p, q, r*, at the bottom of the eye. The cavity of the eye is not filled with one liquor, but with three of different sorts. That contained in the outward space *ABCOEGFDO* is called the aqueous humor, being perfectly fluid like water; the other contained in the inward space *E p q r DFG* is a little thicker like the white of an egg, and is called the vitreous humor; the third humor *FG* is shaped like a lens of unequal convexities, lying between the two former, and fixed to the side coats by filaments or threads extended all round it, and is called the crystalline humor, being hard like the white of an egg boiled, but as clear as the other two, and differs from them in a greater degree of refractive power; whereby the rays that came from the points *P, Q, R*, having received a degree of convergence by the refraction of the cornea *ABC*, are made to converge a little more by other refractions at the surfaces of the crystalline *FG*; so that uniting in as many other points *p, q, r*, upon the retina, they represent the points of the object *P, Q, R*, from whence they came. And perhaps the rays are so directed by these secondary refractions at the crystalline, as to fit the cavity *p q r* intended to receive them; which otherwise must have been a portion of a larger sphere<sup>a</sup>, according to the fictitious design in the former figure.

<sup>a</sup> Art. 111.  
latter part.

The crystal-  
line makes all  
pictures di-  
stinct.

120. Besides this there was greater need of the lens *FG* upon another account; namely, to help the eye to conform itself for seeing objects distinctly at all distances, which was wanting in the fictitious eye. There are two ways of doing it by the help of this lens *FG*, in order









order to see things near at hand; either by moving it nearer to the outward cornea, or by increasing its convexity, or perhaps by doing both at once. If it is moved towards the cornea, this may be effected by the pressure of the muscles against the sides of the eye, and consequently against the vitreous humor; but if the crystalline alters its figure and becomes rounder for seeing near objects, the filaments *DF*, *EG*, whose greater tension helps to flatten it, may perhaps be slackened by the lateral pressure aforesaid; and possibly both these alterations are made at the same time. The hole or pupil *O* is not placed in the center of the cornea *ABC*, as in the fictitious eye, but somewhat nearer to its front. The reason is uncertain, unless this also may contribute to make the images coincide with the cavity of the retina, (in all their parts,) which otherwise must have been shaped according to a larger sphere<sup>a</sup>.

121. The diameter *AT* of the sphere of the eye is about an inch of the Rheinland foot, which is much the same as the old Roman foot; and the diameter of the outward cornea is about three fifths of an inch; the breadth of the pupil, *O*, has no fixt measure, being greater or smaller, as any one may try, according as less or more light shines upon the eye; it also contracts at the near approach of any small object, when we endeavour to view it distinctly. Its fabrick is admirable in this respect that while it alters its size it preserves its roundness<sup>1</sup>. So far Mr. *Hugens*<sup>2</sup>, to which I add something from Sir *Isaac Newton*<sup>3</sup>.

122. This account of the eye and of the cause of vision is farther confirmed by these arguments; that anatomists when they have taken off from the bottom of the eye that outward and thickest coat called the *dura mater*, can see through the thinner coats the pictures of objects lively painted thereon. And these pictures propagated by motion along the fibres of the optick nerves into the brain, are the cause of vision. For according as these pictures are perfect or imperfect, the object is seen perfectly or imperfectly. If the eye be tinged with any colour (as in the disease of the jaundice) so as to tinge the pictures in the bottom of the eye with that colour, then all objects appear tinged with the same colour.

123. If the humors of the eye decay by old age so as by shrinking to make the cornea and coat of the crystalline humor grow flatter than before, the light will not be refracted enough, and for want of a sufficient refraction will not converge to the bottom of the eye,

<sup>a</sup> Art. 111.  
latter part.

Some dimensions of an human eye.

Pictures on the retina the cause of vision.

Confused pictures in old men's eyes how caused, and mended by convex glasses.  
Fig. 131.

<sup>1</sup> See this accounted for in Mr. *Cheffelden's* Anatomy, p. 319. 3d Ed.

<sup>2</sup> Dioptrica prop. 31.

<sup>3</sup> Opticks pag. 12.

but



Fig. 132.

but to some place beyond it; and by consequence will paint in the bottom of the eye a confused picture; and according to the indistinctness of the picture the object will appear confused. This is the reason of the decay of sight in old men, and shews why their sight is mended by spectacles. For the convex glasses supply the defect of plumpness in the eye, and by increasing the refractions make the rays converge sooner, so as to convene distinctly at the bottom of the eye, if the glass has a due degree of convexity.

Confused pictures in short-sighted eyes, how caused and mended by concave glasses.

Fig. 133.  
\* Art. 108.

Fig. 134.

124. And the contrary happens in short-sighted men, whose eyes are too plump. For the refraction being now too great, the rays converge and convene in these eyes before they come at the bottom; and therefore the picture made in the bottom and the vision caused thereby will not be distinct, unless the object be brought so near the eye as that the place where the converging rays convene may be removed to the bottom<sup>a</sup>; or that the plumpness of the eye be taken off and the refraction diminished by a concave glass, of a due degree of concavity, as is represented in fig. 134; or lastly that by age the eye grows flatter till it comes to a due figure. For short-sighted men see remote objects best in old age, and therefore they are accounted to have the most lasting eyes. So far Sir *Isaac Newton*.

Glasses for defective eyes determined.

125. In order to determine the properest glasses for defective eyes, the limits of confused and distinct vision, or the distances of those places from the eye, where an object begins to appear confused, may be found by measuring the least distance from which a long-sighted person can see a pretty large print distinctly and read it readily: and likewise by measuring the greatest and the least distances, from which a short-sighted person can see a small print distinctly and read it readily: or still more exactly by placing the end of a long ruler close to the eye, or rather a little under it, and by observing the greatest and least distances at which the lines drawn lengthways along the ruler, begin to appear confused. I shall call those glasses the properest for defective eyes, which are the least concave or the least convex of any that will answer the purpose of distinct vision; for a reason to be mentioned hereafter.

Fig. 135.

126. Let  $Eq$  be the least distance from which any small objects appear distinct to the eye of a long-sighted person; and  $E\mathcal{Q}$  the least distance from which he wants to see them distinctly. Towards  $q$  take  $\mathcal{Q}F$  to  $\mathcal{Q}E$  as  $\mathcal{Q}E$  to  $\mathcal{Q}q$ , and  $EF$  will be the focal distance of a convex lens, which being put close to his eye, will make him see an object distinctly at any place between  $\mathcal{Q}$  and  $F$ , and possibly beyond  $F$ . For the rays that flow from  $\mathcal{Q}$  will emerge from the glass and will enter the eye as if they had come directly from  $q$ , to the naked eye;



eye<sup>a</sup>; and supposing  $\mathcal{Q}$  to recede from the eye,  $q$  will also recede<sup>a</sup> Art. 107. from it to infinity through places where the naked eye can see distinctly<sup>b</sup>. And therefore the refracted rays, diverging as from these<sup>b</sup> Art. 108. places, will also produce distinct vision of the object  $\mathcal{Q}$  as far as to  $F$ ; and still farther if the person can see distinctly by converging rays.

127. Therefore if he wants to see distinctly from no less a distance than half  $Eg$ , that is, only as near again as with his naked eye, a convex lens whose focal distance is  $Eg$  will be the properest; and will make him see an object distinctly at any distance not less than half  $Eg$ . For supposing  $\mathcal{Q}q$  and  $\mathcal{Q}E$  to be equal, the point  $F$  will fall upon  $q$  by the foregoing proportion.

128. Let  $EF$  be the greatest distance from which an object at  $F$  Glasses for short sighted persons. Fig. 136. appears distinct to the eye of a short-sighted person; and it will also be the focal distance of a concave lens, which being put close to his eye at  $E$ , will be the properest for seeing remote objects distinctly. Because the rays of a pencil, which come from any remote object, and consequently fall parallel upon the lens, will emerge from it to the eye, as if they had come directly to the naked eye from an object at  $F$ . And consequently the picture of the remote object formed upon the retina by rays refracted through the lens, will be as distinct as the picture of an object at  $F$  seen by unrefracted rays.

129. Let  $Eg$  be the least distance from which the same person can Fig. 137. see an object distinctly with his naked eye; then say as  $\mathcal{Q}F$  to  $\mathcal{Q}E$  so  $\mathcal{Q}E$  to  $\mathcal{Q}q$ , and placing  $\mathcal{Q}q$  towards  $F$ , the point  $q$  will be the nearest point, which he will be able to see distinctly through the lens abovementioned. For by art. 107, the rays of a pencil which fall upon the lens converging towards  $\mathcal{Q}$ , will after refraction converge to  $q$ ; and on the contrary, the rays which flow from  $q$  will emerge from the lens diverging from  $\mathcal{Q}$ ; and supposing the point  $q$  to recede from the eye, the point  $\mathcal{Q}$  will also recede from it to such places where the naked eye can see distinctly<sup>b</sup>. But if the point  $q$  approaches<sup>b</sup> Art. 108. towards the eye, the point  $\mathcal{Q}$  will also approach towards it, to such places where the naked eye cannot see distinctly, by the supposition.

130. Consequently if  $\mathcal{Q}F$ , the space between the limits of confused vision, be not less than  $\mathcal{Q}E$ , that one glass whose focal distance is  $EF$ , will make all objects appear to him distinct which are any where placed beyond  $F$ , the reach of his naked eye. For in this case  $\mathcal{Q}q$  cannot be greater than  $\mathcal{Q}F$ , as is manifest by the proportion above.

131. But if he wants a pair of concave spectacles to read or write Fig. 138. with, let the distance  $Eg$  be no greater than what is convenient for that purpose, and let  $\mathcal{Q}F$  be the limits of confused vision as before; and



and towards  $q$  take  $FG$  to  $FE$  as  $FE$  to  $Fq$ , and a concave lens whose focal distance is  $EG$  will be the properest for this purpose. For by art. 107, the rays of a pencil which fall on this glass converging towards  $F$  will converge to  $q$  after refraction, and on the contrary; and therefore he will see an object distinctly as far off as  $q$ ; and also nearer than  $F$ , if  $QF$  be but half of  $EF$ . For supposing rays to fall on the lens converging towards  $Q$ , say as  $QG$  to  $QE$  so  $QE$  to  $QH$ , and the refracted rays will converge to  $H$  and consequently  $H$  will be the nearest point that can be seen distinctly through this glass. But if  $Q$  bisects  $EF$  it is manifest that  $QH$  is less than  $QF$ ; because  $QG$ ,  $QF$ ,  $QH$  are now continual proportionals.

Directions for  
the choice of  
convex and  
concave  
spectacles.

132. Thus any person may be fitted with the properest glasses though he lives at a distance from the shops where they are sold, by sending the workman their focal distances computed by the foregoing rules. But if choice of glasses be at hand, he may be better fitted by trial; observing only to use those glasses which are the least concave or the least convex of any that will fit the eye. These are the glasses which I have computed and called the properest. For since they cannot be put quite close to the eye, the less any glass is concave, the less it diminishes the pictures of objects upon the retina. It will also accustom the eye to that conformation of its coats and humors, which is proper for seeing objects as far off as it can; and consequently may prevent the eye from growing more and more short-sighted. On the other hand, the less any glass is convex, the less it magnifies the pictures of objects upon the retina; and also obliges the eye to that conformation, which is requisite for seeing objects as near as it can. Both which may prevent the eye in some measure from growing more and more long-sighted. For when the picture upon the retina is very large, it need not be quite so distinct, as when it is smaller, to convey an idea of the same number of parts of an object; and consequently the eye will be more at liberty to recede from that conformation, which is proper for the glass; and to relapse into that to which it inclines, and which is only proper for seeing remote objects.

Equal parts  
of a small ob-  
ject subtend  
equal angles  
at the eye.  
Fig. 139.

133. When the perpendicular subtenſe  $BC$  of a small angle  $BAC$  is divided into any number of equal parts  $BH$ ,  $HI$ ,  $IC$ , the lines,  $HA$ ,  $IA$ , drawn from the points of division to  $A$ , will divide the angle  $BAC$  into the same number of parts, which will be nearly equal among themselves. For they would be so exactly if the line  $BC$  was an arch of a circle, drawn upon the center  $A$ ; from which it differs so much the less as the angle at  $A$  is smaller; and so the proposition is exactest in the smallest angles.

134. When



134. When the distance  $AB$  is double or treble of  $Ab$ , the subtense  $BC$  will be double or treble of the subtense  $bc$  of the same angle at  $A^1$ . Divide  $BC$  into its parts  $BH$ ,  $HI$ ,  $IC$ , each equal to  $bc$ , and the rays  $HA$ ,  $IA$ , will divide the angle  $BAC$  into as many equal parts<sup>b</sup>. Therefore when two angles  $bAc$ ,  $BAH$  are subtended by the same or by equal lines  $bc$ ,  $BH$ , the magnitude of the first angle  $bAc$ , will be to the magnitude of the second  $BAH$ , as the second distance  $BA$  to the first distance  $bA$ .

Small angles subtended by the same perpendicular are reciprocally as its distances from the angular points.  
<sup>a</sup> Art. 24.  
<sup>b</sup> Art. 133.

135. In determining the magnitude of pictures upon the retina, only one ray in each pencil need be considered; because when the picture is distinct, all the rays in any one pencil are collected to one and the same point of the retina. Or, which is much the same, we may suppose the pupil of the eye contracted to a point: and, for greater simplicity and ease of the imagination that this point  $O$  is a little hole at the center of a dark, hollow hemisphere  $DqE$ , admitting only single rays straight through it without any refraction at all. For then the lengths of these pictures  $pqr$  will increase and decrease as the angle  $pOr$  does, or as the angle  $POR$  does; which I am going to shew to be the property of the natural eye<sup>c</sup>: and if the semi-diameter  $Oq$ , of this hollow hemisphere, be about  $\frac{1}{8}$  of an inch, or of the axis of an human eye, the pictures of the same objects will always have the same bigness in both sorts of eyes<sup>c</sup>, very nearly.

The pupil of the eye may be considered as a point.

Fig. 129.

<sup>c</sup> Art. 136.

<sup>d</sup> Art. 142.

136. The diameters or lengths of the pictures of objects upon the retina are measured by, or proportionable to, the angles which the rays that come from the extremities of the object do make in falling on the eye; provided those angles be but small. For let two or more objects  $PQ$  and  $\pi\kappa$ , either parallel or oblique to each other, subtend the same angle  $POQ$  or  $\pi O\kappa$  at  $O$ ; and because the particles of light flowing from  $P$  and  $\pi$  describe the same line  $P\pi O$ , they will be refracted to the same point  $p$  upon the retina; and in like manner those that flow from  $Q$  and  $\kappa$  will be refracted to the same point  $q$ ; and so the pictures  $pq$  of the objects  $PQ$ ,  $\pi\kappa$ , which subtend the same angle at  $O$ , are the same in magnitude; which was the first thing to be proved.

Diameters of pictures on the retina are as the angles subtended by the object at the eye.  
 Fig. 140.

Now the pictures of objects painted upon the retina of a dead eye are found by experience to be perfectly well shaped and proportioned in their parts<sup>e</sup>; that is, the proportion of the parts  $pq$ ,  $qr$ , of the whole picture  $pqr$ , is the same as that of the parts  $PQ$ ,  $QR$ , of the whole object  $PQR$ , and this latter proportion is very nearly the same as that of the angles  $POQ$ ,  $QOR$  subtended by the parts  $PQ$ ,  $QR$ <sup>f</sup>; and so the proposition is proved when the objects  $PQ$ ,  $QR$  are both at the same distance from the eye. And since it was shewn just before,

<sup>e</sup> Art. 122.

<sup>f</sup> Art. 133.



fore, that the objects  $PQ$  and  $\omega\kappa$  have the same picture  $pq$ , it follows that the proportion of the pictures of the objects  $\omega\kappa$  and  $QR$  is the same as that of the angles  $\omega O\kappa$ ,  $QOR$ , subtended by them at the eye.

Are reciprocally as the distances of the object from the eye.

137. When an object approaches towards the eye, the diameter of its picture upon the retina increases in the same proportion as the distance between the eye and the object decreases; and on the contrary, it decreases in the same proportion as that distance increases. For the diameter of its picture increases in the same proportion as the angle increases, which the object subtends at the eye<sup>a</sup>; and this angle, when small, increases in the same proportion as the distance between the eye and object decreases<sup>b</sup>.

<sup>a</sup> Art. 136.

<sup>b</sup> Art. 134.

Brightness of pictures not altered by the distance of the object from the eye.

138. The degree of brightness of the picture of an object painted upon the retina continues the same, at all distances between the eye and the object; provided none of the rays be stopt by the way, and that the pupil does not alter its aperture. For instance, when the eye approaches as near again to the object, the picture upon the retina becomes double in length and double in breadth, and consequently quadruple in surface; for the surface would be double, if its length alone or breadth alone was double. The quantity of rays received through the same aperture of the pupil, at half the distance from the object, is also quadruple<sup>c</sup>: and being equally spread over four times the quantity of surface of the retina, they are just as dense as before when the object was at twice the distance.

<sup>c</sup> Art. 25.

Faintness of pictures of remote objects how caused.

139. It follows then that the faint appearance of remote objects is owing to the opacity of the atmosphere, which hinders part of their light from coming to the eye. Accordingly we find that the sun, moon, and stars appear very faint when near the horizon, and brighter continually as they rise higher; because the tract of vapours, which lies in the way of the rays, is longest and thickest near the horizon; and becomes thinner and shorter as the objects rise higher, and consequently does less obstruct the passage of the rays.

Their degrees of brightness by day-light and moon-light compared.

140. The sensibility of the eye, or its power to discern objects, without inconvenience, by different quantities of light, is vastly extensive. For instance, the disproportion in the quantities of light, cast upon the horizon by the sun and moon, at any equal altitudes, I find is no less than 90 thousand to 1, when the moon is full; or no less than 180 thousand to 1, when the moon is in the quarters. And the proportion between those parts of the lights of the sun and moon, whatever they be, which are reflected to our eyes from the same object by day and by night, can hardly be different from the proportion of the whole lights. Allowing then that the aperture of the pupil



pupil may possibly be 8 or 9 times less by day than by night, (that is about 3 times less in diameter,) yet the proportion in the quantities of day-light and moon-light, received by the eye from the same object, to illuminate a picture of the same bigness, will be no less than 20 thousand to 1, when the nights have a middle degree of moon-light; I say no less, because the numbers here given are deduced from a rule, which is built upon this principle; that the moon reflects all the light received from the sun; which cannot be true, by reason of the appearance of very large obscure places in her body; and in all probability a great part of the incident light is buried and lost even in the brightest places.

The rule I mentioned is this, day-light is to moon-light as the surface of an hemisphere, whose center is at the eye, to the part of that surface which appears to be possessed by the enlightened part of the moon: so that the whole heavens covered with moons would only make day-light. This will be evident enough from the following considerations, though I invented it another way. Day-light is made by innumerable reflections of the sun's rays from all sorts of bodies till at last they come to our eyes: for if this were not so, we could see nothing in the world, even in the day time, but the sun and stars and self-shining substances<sup>a</sup>. Accordingly we find that day-light is<sup>a</sup> Art. 2. much the same, whether the sun shines out or not, in the place we are in; because his light is reflected to us from a vast quantity of earth, air and clouds extended all round us, perhaps to a hundred miles or more. So that the absence of the sun's rays from a particular place scarce alters day-light. Another thing is that the moon by day appears like a cloud in the air of a middle degree of brightness; some appearing duller and some brighter than the moon itself. The rays of the sun being therefore intercepted in the night from all the visible clouds, and being reflected to us by the moon only, it follows that day-light is to moon-light, as the apparent surfaces of all the visible clouds, to the apparent surface of the visible part of the moon, considered as the only cloud which remains enlightened. And these two lights, whatever be the distances of the moon and clouds, are just the same as if those bodies were all placed at any equal distances from us, and composed the surface of an hemisphere<sup>b</sup>; whose parts are the true measures of the parts of the<sup>b</sup> Art. 138. light which comes to us.

141. A vast disproportion between the lights of the sun and moon appears also by experiments made with burning-glasses; either by refraction of the rays through very broad lenses, or by reflection from very broad concave-glasses or metals: which by collecting the

And confirmed by experiments with burning-glasses.



rays of the sun into a small round image at the focus, do excite a more violent heat and burn quicker than the hottest wind-furnaces: as appears by their melting and calcining the hardest metals, and by vitrifying bricks and stones, in much less portions of time than a minute<sup>1</sup>. Yet the rays of the moon being collected by the same glasses, do not excite the least sensible heat; nor do they sensibly affect the nicest thermometer, when cast upon the ball of it<sup>2</sup>, though the brightness of the light be very sensibly increased. By measuring the breadth of the round image at the focus, and by comparing it with the breadth of the glass itself, it appears that some of these burning-glasses collect the incident rays into a space about 2 thousand times less than they possessed at their incidence. But by the preceding calculation, the light of the full moon must be condensed about 90 thousand times<sup>3</sup>, to make it as dense and as warm as the direct rays of the sun. It is no wonder then that the heat of the moon's rays is not sensible in the focus of the glass, being then even 40 or 50 times thinner than the direct rays of the sun. For it is found by experiments made with these glasses that the degrees of heat are proportionable to the densities of the rays: which being compared with a scale of the degrees of heat and warmth of several natural bodies, determined by Sir *Isaac Newton*, in the philosophical transactions<sup>4</sup>, it appears there is a vast disproportion between the degrees of light which the eye can bear and be sensible of, and the degrees of its heat which the touch can bear and be sensible of.

<sup>1</sup> Art. 140.

Vision limited both by magnitude and distance.

142. Dr. *Hook* assures us that the sharpest eye cannot well distinguish any distance in the heavens, suppose a spot of the moon's body, or the distance of two stars, which subtends a less angle at the eye than half a minute; and that hardly one of a hundred can distinguish it when it subtends a minute<sup>4</sup>. If the angle be not bigger than this, the two stars appear to the naked eye, as if they were but one. I have been present at making the experiment, when a friend of mine, who had the best eyes of all the company, could scarce perceive a white circle upon a black ground, or a black circle upon a white ground, or against the sky-light, when it subtended a less angle at the eye than two thirds of a minute; or which is the same thing, when its distance from the eye exceeded 5156 times its own diameter: which agrees well enough with Dr. *Hook's* observation.

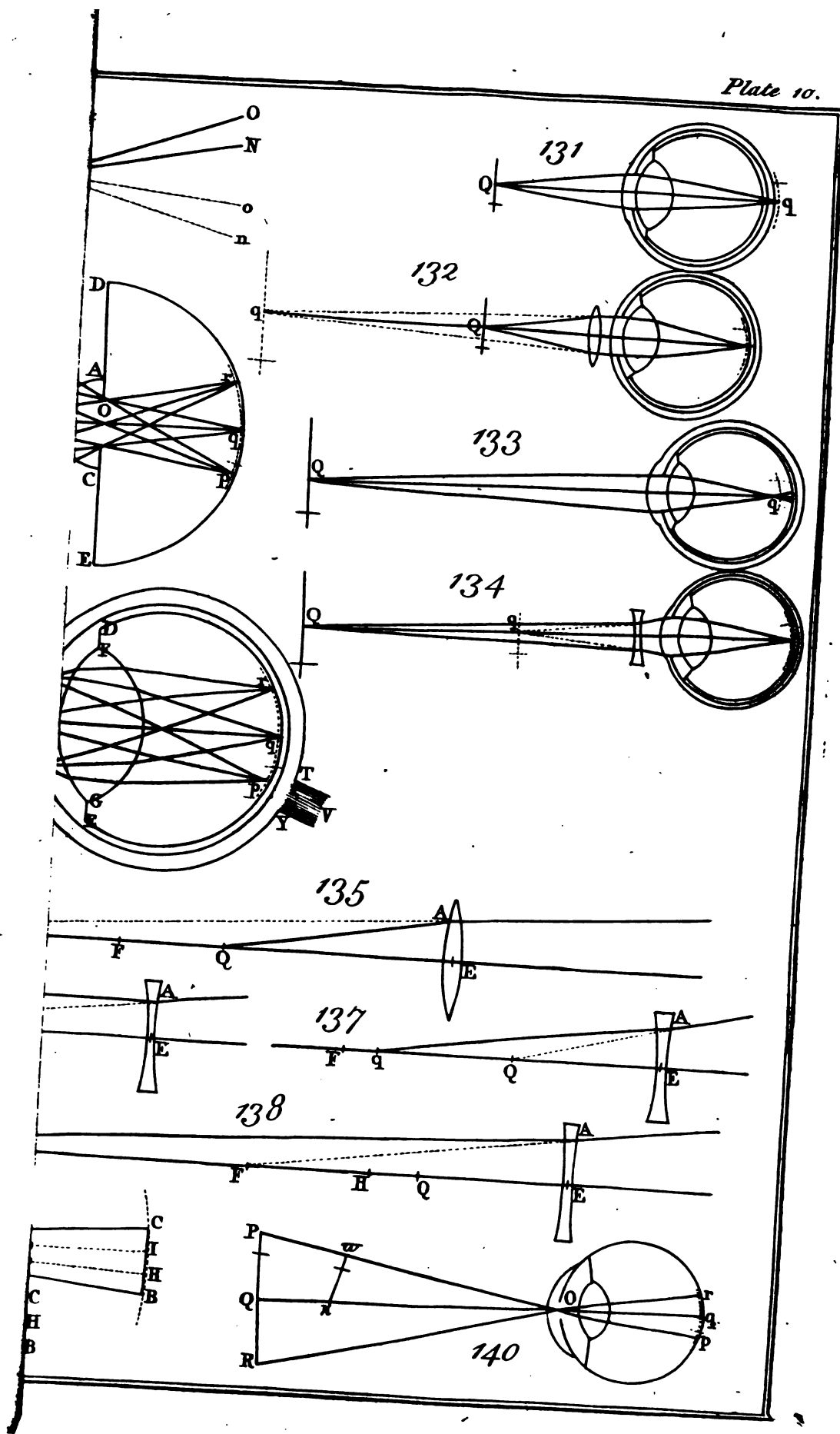
<sup>1</sup> Phil. Trans. abr. by *Lowth.* Vol. 1. p. 211. and by *Jones* Vol. 4. p. 190.

<sup>2</sup> Ibid. *Lowth.* p. 213. and Mem. de l'Acad. Roy. des Scien. ann. 1705. p. 455. 8°.

<sup>3</sup> N°. 270. or *Jones's* abr. Vol. 4. part. 2. p. 1. It. Mem. de l'Acad. an. 1703. p. 233.

<sup>4</sup> Animadversions on Hevelii machina cœlestis p. 8.





## CHAP. X.

## CONCERNING VISION WITH GLASSES.

Apparent directions of visible points defined.

\* Art. 18.  
Fig. 4.

Fig. 141. to 152.

Fig. 153.

\* Art. 54.

146. **A**NY small object or point of an object, seen by refracted or reflected rays, appears somewhere in the direction of that line, which the visual ray describes after its last refraction or reflection in falling upon the eye.

In the experiments to prove the laws of reflection and refraction\*, the pin at *B*, seen by a ray reflected from the water, appeared somewhere in the line *AC* produced, which the visual ray *BCA* described after reflection at *C*, when it advanced to the eye. And as the whole line *CE* appeared lifted up by the refraction at the water, as if it had been a continuation of the line *AC* straight on, so if a straight oar be in part immersed obliquely in water, it appears crooked, as if the part immersed was broken at the surface, and lifted higher. For this part of the oar is seen in the direction of rays which are bent downwards by refraction at their emergence from the water, and consequently advance to the eye as if they came from a place in the water which is higher than the real place of the oar. In like manner any point *P* of an object seen by the ray *PAO* twice refracted, either by passing through the edge of a prism, or of a concave or convex lens, or through the sides of a globe or decanter, or of a drinking glass filled with any transparent liquor; or seen by a ray *PAO* reflected from a plane or spherical looking-glass, appears to the eye at *O*, somewhere in the direction of the last refracted or reflected ray *AO*. Lastly, an object *P* viewed by the eye at *O*, through a multiplying glass, appears at one view in as many different places, *p*, *p* 1, *p* 2, situated in as many different directions *OA*, *OB*, *OC* of the last refracted rays produced, as the glass has different surfaces *DE*, *EF*, *FG* differently inclined to the opposite surface *DH*. For these surfaces, like so many different prisms, give the visual rays *PIAO*, *PKBO*, *PLEO*, so many different bendings at *I* and *A*, *K* and *B*, *L* and *C*, and make them fall upon the eye at *O* in as many different directions *AO*, *BO*, *CO*<sup>b</sup>. And in all these instances when the reflecting or refracting surfaces of the water or glasses are shaken by the wind, or otherwise, the objects seen by reflection or refraction appear to shake and tremble; because the last directions of the visual rays are shaken and varied by those motions.

Now



Now the reason why an object or point of an object appears always in the direction of the last refracted or reflected ray, is, because the place of its picture upon the retina is the same as it would be if the object was really removed from its proper place into the direction of that ray, and was seen by direct rays. And having no sensation of the previous reflections or refractions of the rays at the glasses, but only of their action upon a certain place of the retina, we form the same judgement of the place of the object as we used to do in the more common cases of direct vision.

It is manifest then, that any point  $P$  of an object  $PQ$  seen by refractions or reflections, appears somewhere in the line  $pO$ , drawn from the corresponding point  $p$  of its last image to the eye at any place  $O$ . Because all the rays which flowed from  $P$  do after the last refraction or reflection flow from or towards the corresponding point  $p$  of the last image. The reason why I say the last image will be mentioned in the 155th article.

147. It is also manifest why an object seen by refracted or reflected rays appears sometimes upright and sometimes inverted. For when the refracted or reflected rays  $AO$ ,  $CO$ , have the same situation with respect to each other, as two rays that come directly from the same points of the object to the eye, these points will appear in the same situation with respect to each other in both cases<sup>a</sup>. But if the rays that come from these points shall have crossed each other before they arrive at the eye, they will then have a contrary situation to that of two rays coming directly from the same points to the eye; and consequently these two points will appear in the glass in a contrary situation<sup>b</sup>. And one may add that in the former case, the picture upon the retina of the eye will have the same position, though not the same magnitude, as if the glass was removed, and will have a contrary position in the latter case.

148. The apparent magnitude of an object,  $PQ$ , seen by refracted or reflected rays either upright or inverted, is a quantity of visible extension, measured by the angle,  $AOC$ , which the two rays,  $AO$ ,  $CO$ , that came from its extremities,  $P$ ,  $Q$ , do make, after their last refraction or reflection, in falling on the eye. Or in other words, the object appears greater or smaller in proportion as that angle  $AOC$  is greater or smaller. Because its extremities appear in the directions of the last refracted or reflected rays  $OA$ ,  $OC$ <sup>d</sup>; and also because its picture upon the retina is greater or smaller in proportion as these rays constitute a greater or smaller angle at the eye<sup>e</sup>.

149. Therefore the apparent magnitude of an object,  $PQ$ , is also measured by the angle  $pOq$  which its last image  $pq$  subtends at the

And determined.

Fig. 141. to

153.

Their apparent situation determined.

Art. 146.

Art. 146.

Apparent magnitude in glasses defined.

Art. 135.

Art. 146.

Art. 136.

And determined.



the eye. For the lines  $AO$ ,  $PO$  are but one line continued, and so are  $CO$ ,  $QO$ , and therefore the angles  $AOC$ ,  $POQ$  are the same when the image lies before the eye, and are equal when it lies behind it.

How it varies.

<sup>a</sup> Art. 144.

<sup>b</sup> Art. 134.

When invariable.

Fig. 154. to 157.

<sup>c</sup> Art. 104. 70.

150. Hence the apparent magnitude of an object increases and decreases in proportion as the eye approaches to or recedes from its last image, (just as if it was a real object <sup>a</sup>.) placed either before or behind the eye. For when the image is fixed, the angle  $POQ$ , when small, increases in the same proportion as  $OQ$  decreases, and on the contrary <sup>b</sup>.

151. Hence if the last image be removed to an infinite distance, that is, if the object be placed in the principal focus of a lens, sphere, or concave looking-glass <sup>c</sup>, its apparent magnitude to the eye at any place whatever will be invariably the same; and equal to its apparent magnitude seen by the naked eye, supposing it put into the place of the center of the sphere, lens, or reflecting concave. For since all the rays of any one pencil, are parallel to its axis  $PE$ , the angle  $COA$ , which measures the apparent magnitude at any point  $O$ , is every where equal to the angle  $QEP$  at the center  $E$ .

Fig. 158, 159.

The apparent magnitude of the object will also be invariable where-ever it be placed, when the eye is fixed at the principal focus of any glass which makes parallel rays converge to the eye. For conceiving them to flow back again from the eye to the object, they will fall upon the same points of the object from whence they came while it is moved in any place along the axis of the glass: and no other rays but these can return from the same points of the object to the eye in that place: therefore the several parts of the object will always be seen under the same angles, and consequently will appear of the same magnitudes <sup>d</sup>.

<sup>e</sup> Art. 148.

Compared to the true magnitude.

<sup>f</sup> Art. 149.

<sup>g</sup> Art. 143.

152. The apparent magnitude of an object seen by reflected or refracted rays being measured by the angle which its last image subtends at the eye <sup>e</sup>, and its apparent magnitude to the naked eye in any place being measured by the angle which the object itself subtends at the eye in that place <sup>f</sup>, it follows that the former apparent magnitude is to the latter, as the former angle to the latter angle. For the measures of things and the things measured by them are proportionable.

When equal to the true.

<sup>h</sup> Art. 108.

153. Consequently the apparent magnitude of an object seen in a glass, will be equal to its apparent magnitude to the naked eye in the same place, if the glass was removed; First, when the object touches any thin lens, or any single surface. For the image is then equal to the object and coincides with it <sup>g</sup>. Secondly, when the eye touches any thin lens or any reflecting surface. For then the ray

$PAO$



$PAO$  will pass from the object to the eye through the middle of the lens very nearly, and therefore being almost straight<sup>a</sup> will make<sup>a</sup> Art. 60. nearly the same angle with the axis as an unrefracted ray would do: and when the point of incidence,  $A$ , coincides with  $C$  at any reflecting surface, the incident and reflected rays  $PA$ ,  $AO$ , produced, will also make equal angles with the axis or perpendicular  $QC$ <sup>b</sup>; and so<sup>b</sup> Art. 8. the object will appear under the same angle as it would do to the naked eye turned about. Thirdly, when the eye is at the center of a reflecting concave. For then the incident and reflected rays  $PA$ ,  $AO$  will coincide with the direct ray  $PE$ <sup>c</sup>, and consequently will<sup>c</sup> Art. 10. make the same angles with the axis. Fourthly, when the object is at the center of the reflecting concave. For then the reflected image is also at the center and is equal to the object<sup>d</sup>. Fifthly, when a ray<sup>d</sup> Art. 72. 79. coming directly from  $P$  to  $O$ , would make an angle with the axis equal Fig. 146, 148. to the angle  $AOC$ , which the refracted or reflected ray  $PAO$  makes<sup>e</sup> 152. with it on the other side.

154. These cases being excepted the apparent magnitude of an object seen through a concave lens is always less than the true; and when it is seen upright through a convex lens, or a globe, it is greater than the true. For the ray  $PAO$ , coming from the extremity of the object to the eye, is bent by the concave lens from its axis, and therefore makes a less angle with it at the eye than a ray coming directly from that extremity to the eye. But the same ray is bent by the convex lens towards its axis, and therefore makes a greater angle at the eye than the direct ray: and the apparent magnitudes are measured by these angles. Less than the true through a concave and greater through a convex glass.

155. What has hitherto been demonstrated concerning the apparent magnitude of an object  $PQ$ , will continue in force if you suppose the object  $PQ$  to be an image formed by another glass or other glasses. For the rays diverge from either of them in the same manner, and for this reason I have always called  $pq$  the last image of the object. The whole applied to vision through any number of glasses.

156. The place of the eye at  $O$  being given, to determine what part of an object is visible in a given portion or aperture  $AC$  of any refracting or reflecting glass, draw  $OA$  to the edge of the aperture and produce it till it cuts the image in  $p$ , and through the center of the glass draw  $pE$  cutting the object in  $P$ ; and  $PQ$  will be the part in view in the aperture  $AC$ . For the whole pencil of rays flowing from  $P$  will belong to  $p$  after refraction or reflection<sup>e</sup>, and consequently some one of those rays will advance to the eye in the line  $AO$  drawn through  $p$ . If the image be at an infinite distance all the rays that belonged to  $p$  will be parallel to the axis of the pencil; Fig. 154. to 157. therefore



therefore  $PQ$  is now determined by drawing  $EP$  parallel to  $OA$ .  
 In a plane looking-glass,  $pP$  must be drawn from  $p$  parallel to  $qQ$ ,  
 or perpendicular<sup>a</sup> to the glass, to cut off the part  $PQ$  visible in the  
 aperture  $AC$ . For this glass may be considered as having a center  
 at an infinite distance from it.

How it va-  
 ries.

157. Hence if the glass and object be fixed, the part in view in a  
 given aperture will decrease perpetually while the eye recedes from  
 the glass; unless the image be behind the eye. For then it will de-  
 crease only till the eye arrives at the image, and after the eye has  
 passed by the image it will increase perpetually. The reason is be-  
 cause the object and image, being fixed in their places, do both  
 increase or both decrease together, being both terminated by two  
 lines  $Pp, Qq$  that meet or cross in  $E$  the center of the glass.

When great-  
 est and least.

158. Therefore the part in view is greatest when the eye is close  
 to the glass, and least when close to the image; and, in this latter  
 case, it appears infinitely magnified. For conceiving the distance  
 $Oq$  infinitely diminished, the parts  $pq, PQ$  cut off by the lines  $AOp$ ,  
 $pEP$  will both be infinitely diminished; but the magnitude of the  
 angle at  $O$ , subtended by  $pq$  or by  $AC$ , continues finite while the  
 angle subtended by  $PQ$  at  $O$  is infinitely diminished: and so the dis-  
 proportion between these angles, that is, between the apparent and  
 true magnitudes of the particle  $PQ$ <sup>b</sup> is infinitely great. It appears  
 also infinitely confused, when the pupil is open, for the reason given  
 in the following articles.

The size of a  
 looking-glass  
 sufficient to  
 see all ones  
 own body.  
 Fig. 149.  
<sup>c</sup> Art. 66.  
<sup>d</sup> Art. 24.

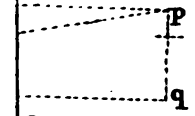
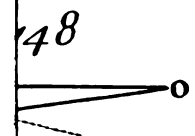
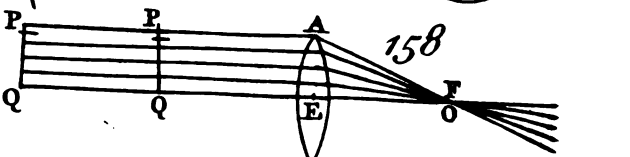
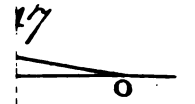
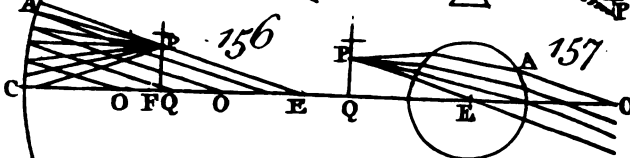
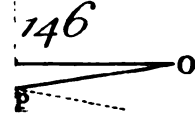
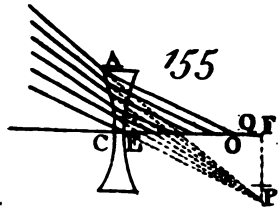
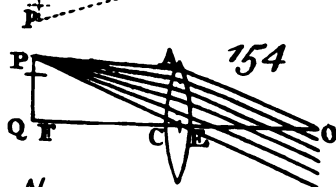
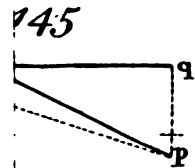
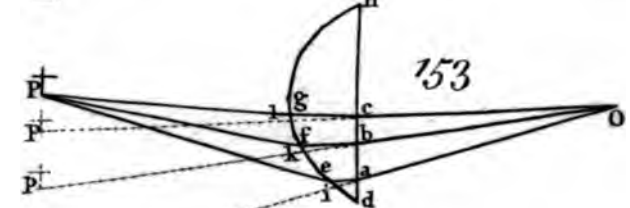
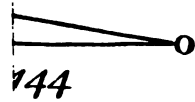
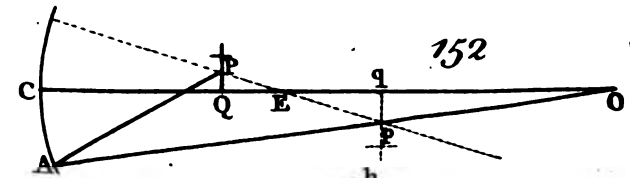
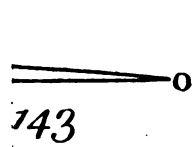
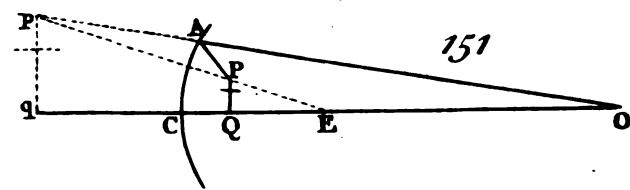
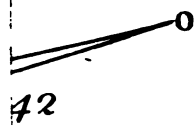
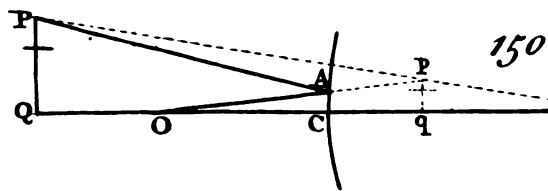
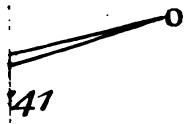
Vision when  
 confused by  
 glasses.

<sup>e</sup> Art. 135.

159. When a person views himself in a plane looking-glass he  
 appears to himself to fill the same part of the glass wherever he  
 stands: and the length and breadth of this part is always half the  
 length and breadth of the corresponding part of his own body. For  
 when  $O$  and  $Q$  coincide,  $OC$  is half of  $Oq$  or  $Qq$ <sup>c</sup>, and therefore  
 $AC$  is half of  $pq$ <sup>d</sup> or  $PQ$ .

160. Hitherto I have considered the pupil of the eye as no bigger  
 than a point, admitting but a single ray from every point of the  
 object<sup>e</sup>; by which means the picture upon the retina would be dis-  
 tinct in all cases. But when the pupil is open, if the image formed  
 by the glass be nearer to the eye than the least distance at which we  
 can see objects distinctly with the naked eye, the appearance through  
 the glass will be confused. Because the rays diverge too much  
 from so near an image to be reduced by the eye to a distinct pic-  
 ture upon the retina. On the other hand, when the rays converge  
 to an image behind the eye, they will be collected to a distinct  
 picture before they arrive at the retina, because the eye is not na-  
 turally used to conform itself to converging rays; and so the vision  
 will





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object from the glass, is then the focal distance of the glass. Now if the glass be a small round globule whose diameter is  $\frac{1}{3}$  of an inch, such as are easily made, its focal distance  $Eg$  being three quarters of its diameter<sup>a</sup>, is  $\frac{1}{9}$  of an inch; and if  $qL$  be 8 inches, the usual distance at which we view minute objects, this globule will magnify at the rate of 8 to  $\frac{1}{9}$  or of 160 to 1.

<sup>a</sup> Art 95:13.  
Astronomical telescope, how much it magnifies, and why.  
Fig. 162.

164. An astronomical telescope is composed of two convex glasses in the following manner.  $PQ$  represents the semidiameter of a remote object, and  $pg$  its picture formed by the convex lens  $L$ , which being next to the object is called the object-glass. In the axis of this glass,  $QLg$  produced,  $EA$  represents another glass more convex than  $L$ , so placed, that as  $qL$  is the focal distance of the glass  $L$ , so  $qE$  is the focal distance of the glass  $E$ ; and  $EL$  the sum of their focal distances. In this situation of the glasses, I say the object will appear to the eye at any point  $O$ , distinct, inverted and magnified at the rate of  $qL$  to  $qE$ , that is of the focal distance of the object glass to the focal distance of the eye-glass.

For the rays which diverge from the point  $q$  of the picture  $pg$ , being refracted by the eye-glass, will emerge upon the eye at  $O$  in lines parallel to the axis  $qEO$ ; because  $qE$  is supposed to be the focal distance of the eye-glass; and for the same reason, the rays which diverge from any collateral point  $p$ , of that picture  $pg$ , will emerge from the eye-glass, after refractions at  $A$ , in lines parallel to the line or ray  $pE$ ; this line being the axis of an oblique pencil of rays, part of which diverge from  $p$  upon the glass. An eye therefore which can see distinctly by pencils of parallel rays being placed anywhere at  $O^*$ , among the intersections of these pencils, will see the points of the object distinctly.

Now to the eye at  $O$  the apparent magnitude of the picture  $pg$ , or object  $PQ$ , is measured by the angle  $EOA^b$ , or by the equal an-

Fig. 162.

<sup>b</sup> Art. 148.  
\* PROPOSITION. The point  $O$ , where the axes  $LB$ ,  $LA$  of the extreme pencils cross the axis  $LE$ , is a little beyond the principal focus of the eye-glass.

For we may suppose all the axes of the several pencils to proceed from the point  $L$  as from a focus of incidence; and their focus after refraction may be found by article 104. But the first term  $Lq$ , in the rule there laid down, is considerably greater than the second  $qE$ , by construction: wherefore the fourth term, that is, the distance of  $O$  from the other principal focus is but small, when compared with the focal length of the eye-glass.  $Q. E. D.$

N. B. The point  $O$ , thus determined, is the place from which an eye can see the most possible of an object, through a given telescope.

The situation of common eyes in a double microscope is determined after the same manner.

Hugens's Dioptr.

gle



gle  $qEp$ ; but to the naked eye at  $L$ , if the glass was removed, the apparent magnitude of the object is measured by the angle  $QLP$ , or by the equal angle  $qLp$ ; the oblique axis  $PLp$  being straight<sup>a</sup>.<sup>a</sup> Art. 61. Therefore the former apparent magnitude is to the latter, as the angle  $qEp$ , to the angle  $qLp$ ; and consequently as the latter distance  $qL$ , to the former  $qE$ <sup>b</sup>.

<sup>b</sup> Art. 134.

165. The object which appeared inverted in the former telescope<sup>c</sup>, will appear upright and distinct through two more convex eye-glasses subjoined to it; at a distance from each other, equal to the sum of their focal distances; and when their focal distances are equal to each other, the object will be magnified just as much as it was before. For the pencils of parallel rays  $EOF$ ,  $AOB$ , &c, which are continued to the glass  $FB$ , will be formed by it into a second image  $\pi\kappa$ ; and the focus  $\pi$ , of any oblique pencil  $OB$ , will be determined by the intersection of the line  $\pi\kappa$ , perpendicular to the common axis of the glasses, and of the oblique axis  $F\pi$ , drawn parallel to the incident rays  $OB$ <sup>d</sup>. This point  $\pi$  being the focus of incident rays on<sup>d</sup> the last glass  $GC$ , the emergent rays  $CD$  will be parallel to their oblique axis  $\pi G$ ; because the rays that flow from  $\kappa$  are supposed to emerge parallel to the direct axis. Therefore to the eye at  $D$ , where these emergent pencils cross, the object will appear distinct, and upright<sup>e</sup>. And when the glasses  $F$  and  $G$  are exactly equal, the image  $\pi\kappa$  will be exactly in the middle between them; and so the triangles  $\pi F\kappa$ ,  $\pi G\kappa$  will be exactly equal. Consequently the angle  $CDG$ , which now measures the apparent magnitude to the eye at  $D$ , will be equal to the angle  $\pi G\kappa$  or  $\pi F\kappa$  or  $BOF$  or  $AOE$ , which measured it before to the eye at  $O$ .

<sup>c</sup> A telescope made of four convex glasses considered. Fig. 163.

<sup>e</sup> Art. 147.

<sup>d</sup> Art. 116.

<sup>e</sup> Art. 147.

166. In a telescope of a given length the quantity of objects taken in at one view, depends upon the breadth of the eye-glass. For as  $AE$  is greater or smaller, the angle  $ALE$  or its equal  $PLQ$  is also greater or smaller; and this angle takes in all the objects that can be seen at one view on one side of the axis of the telescope.

How much they take in at one view. Fig. 162, 163.

167. The difference between the astronomical telescope and Galileo's telescope or a common perspective-glass is this; instead of the convex eye-glass placed behind the image to make the rays of each pencil go parallel to the eye, there is placed a concave eye-glass  $AE$  as much before it; which opens the rays of each pencil that converged to  $q$  and  $p$ , and makes them emerge parallel upon the eye; as is evident by conceiving the rays to go back again through the eye-glass, whose focal distance we supposed was  $Eq$ . The eye must be put close to the glass to receive as many pencils as possible; and then, supposing an emerging ray of an oblique pencil produced backward along

Galileo's telescope considered. Fig. 164.



along  $AO$ , the apparent magnitude of the object is measured by the angle  $AOE^a$  or its equal  $qEp$ , which is to the angle  $qLp$  (or  $QLP$ , the measure of the true magnitude,) as  $qL$  to  $qE$ , as before in the other telescope. It is manifest, by the 147th article, that objects in this telescope appear upright.

This takes in less than the former.

168. The quantity of objects taken in at one view in this telescope does not depend upon the breadth of the eye-glass, as in the astronomical telescope, but upon the breadth of the pupil of the eye. Because the pupil is less than the eye-glass, and the lateral pencils do not now converge to, but diverge from the axis of the glasses. Upon this account the view being narrower is not so pleasant as in the former telescope.

See *If. Newton's reflecting telescope.* Fig. 165.

169. Sir *Isaac Newton's* reflecting telescope magnifies the diameter of a remote object in the proportion of the focal distance of the reflecting concave to the focal distance of the convex eye-glass, and shews it inverted. Let  $ST$  be an image of a remote object  $PQ$  formed by reflections from a large concave surface  $AC$ , and terminated by the lines  $PESA$ ,  $QETC$  drawn through its center  $E$ . Now because this image cannot be viewed through an eye-glass placed directly before it (for then the spectator would intercept the rays that are coming to the concave) therefore let the several pencils of rays which converge towards it in coming from the broad concave  $AC$ , be reflected sideways from a small polished plane, represented by  $ac$ ; and then the second image  $st$ , formed by this plane, will be equal to the first image  $ST^b$ . Let  $tl$  be the focal distance of a small convex eye-glass  $kl$  and the rays which flow from any point  $s$  will be refracted through this glass, to the eye at  $o$ , in the lines  $ko$  drawn parallel to the oblique axis  $sl$ ; and so the apparent magnitude of the object,  $PQ$ , to the eye at  $o$ , will be measured by the angle  $kol$  or  $s/t^c$ : but to the naked eye at  $E$ , it is measured by the angle  $PEQ$  or  $SET$ . Therefore the former apparent magnitude is to the latter, as the angle  $s/t$  to the angle  $SET$  or, (because their subtenses  $st$ ,  $ST$  are equal,) as  $ET$  to  $lt^d$  or as  $CT$  to  $lt$ , when the object is remote<sup>e</sup>. Note that the plane  $acb$  is much too broad in comparison to the concave  $ACB$ , which could not be helped in so small a draught. That the appearance of the object is inverted or turned from right to left, is evident by the 147th article.

Why so much shorter than others.

170. Dioptrick telescopes which magnify much being very long and troublesome to be managed, Sir *Isaac Newton* proposed this method to shorten telescopes<sup>f</sup>; which answers to admiration; as appears by a table in the 12th chapter, of the lengths of both sorts of



telescopes which magnify equally with equal distinctness. The reason why dioptrick telescopes cannot be shortened as much as these, and still magnify as much, by diminishing the focal distances of the eye-glasses<sup>a</sup>, in short is this. The images made by refractions<sup>a</sup> Art. 164. through the convex object-glasses, being much more imperfect than those which are made by reflections from concave surfaces, will not bear to be magnified so much by so small eye-glasses<sup>b</sup>, without appearing confused; and the chief cause of those imperfections in the pictures is the unequal refrangibility of rays of different colours<sup>c</sup>.<sup>c</sup> Art. 220.

171. The following description of Mr. *Gregorie's* reflecting telescope differs from the author's chiefly in this; that he directs his larger reflecting concave to be made of a parabolick figure, and his lesser of an elliptical one, instead of the spherical surfaces now used; which are the only figures that can be polished without insuperable difficulties. Mr. *Gregorie's* reflecting telescope described. Fig. 166.

It is proposed to make a reflecting telescope with two concave metals and a convex eye-glass and to shew its effects. Let the given focal distances of the lesser and the larger concave and of the convex eye-glass, be equal respectively to the lines  $t, T, q$ ; and in a given line  $ctqCl$ , designed for their common axis, take in one and the same direction,  $ct = t$ ,  $tq = T$ ,  $qC = \frac{t \times t}{T}$  and  $ql = q$ ; and place the eye-glass at  $l$ , the lesser concave at  $c$ , and the larger at  $C$ ; so that their concavities may respect each other; and let the incident rays, as  $QA, QB$ , be reflected from the larger to the lesser concave, and from thence to the larger again, where let them pass through a moderate hole made in the middle of it at  $C$ , and then be refracted through the eye-glass  $kl$  to the eye at  $o$ ; I say a remote object will appear distinct and upright and magnified in the ratio of  $T \times T$  to  $t \times q$ ; that is, of the square of the focal distance of the larger concave, to the rectangle under the focal distances of the lesser concave and of the eye-glass.

172. For a pencil of rays  $QA, QB$  coming parallel to the common axis, will be reflected from the larger concave  $ACB$  to its principal focus  $T$ ; where crossing one another, and falling upon the lesser concave  $acb$ , they will be reflected from it to the point  $q$ . For since the focal distance  $TC = T = tq$  by construction; by taking away the common part  $Tq$ , we have  $tT = qC = \frac{t \times t}{T}$  by construction; that is, we have  $tT, tc, tq$  continual proportionals, as they should be<sup>d</sup>; and since  $pl$  is the focal distance of the eye-glass  $kl$ , the<sup>d</sup> Art. 71. rays



rays that flow from  $q$  will emerge from it in parallel lines, and therefore will produce a distinct appearance of the remote point  $\mathcal{Q}$  from which they came.

Fig. 167.

<sup>a</sup> Art. 79.

<sup>b</sup> Art. 79.

<sup>c</sup> Art. 98.

173. Let  $ST$  be the image of the object  $P\mathcal{Q}$  formed by reflection from the large concave; and it will be terminated by the line  $PES$ , drawn through  $E$ , the center of this concave, parallel to the rays  $PA$ ,  $PA$  that flow from  $P^a$ . Again, the rays that flow from this image  $ST$ , will be reflected from the lesser concave and form a second image  $pq$ ; which will be terminated by the line  $Se\phi$  drawn through the center  $e$  of this concave<sup>b</sup>; and the rays that diverge from  $\phi$  will emerge from the eye-glass  $kl$  in the lines  $ko$  parallel to the line  $\phi l$ , drawn through the center of the eye-glass<sup>c</sup>. Therefore the object  $P\mathcal{Q}$  will appear upright, because the rays  $ko$  lye on the same side of the common axis  $\mathcal{Q}l^o$  as the point  $P$  from which they came.

<sup>d</sup> Art. 151.

155.

<sup>e</sup> Art. 134.

<sup>f</sup> Art. 172.

174. In the second image  $pq$  take a line  $qs$  equal to the first image  $TS$ ; and if the image  $pq$  was equal to  $qs$ , the object would appear through the eye-glass under an angle equal to  $qls^d$ ; which is to the angle  $PE\mathcal{Q}$  or  $SET$ , under which it appears to the naked eye at  $E$ , as  $TE$  or  $TC$  to  $ql^e$ ; and so the object would be magnified in the same ratio as in Sir *Isaac Newton's* telescope. But since the triangles  $epq$ ,  $eST$  are similar; and since we had  $tq$  to  $te$  (as  $te$  to  $tT^f$ , and disjointly as  $eq$  to  $eT$ , that is,) as  $pq$  to  $ST$  or  $qs$ ; it appears that  $pq$  is bigger than  $qs$ , and also the visual angle  $kol$  or  $\phi lq$  bigger than  $qls$ , in the said ratio of  $tq$  to  $te$ . And so the object being farther magnified in this ratio of  $tq$  to  $te$  or, by construction, of  $TC$  to  $tc$ , is magnified in the whole in the compound ratio of  $TC$  to  $tc$ , and of  $TC$  to  $ql$ , that is in the ratio of  $TC$  square, to the rectangle under  $tc$  and  $ql$ .

To adapt it to a near object.

<sup>g</sup> Art. 72.

<sup>h</sup> Art. 172.

And to short-sighted eyes.

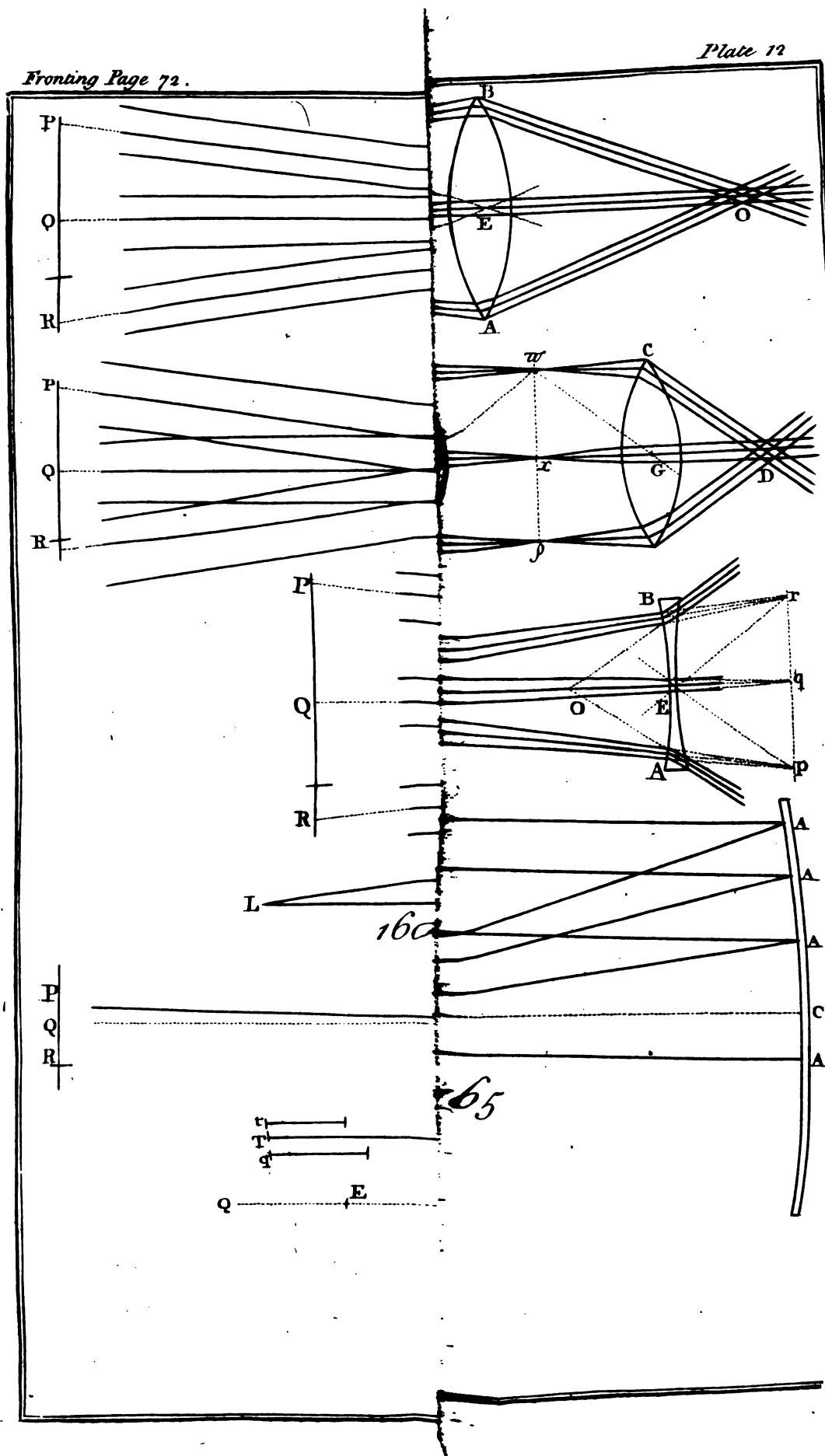
175. For viewing near objects the little concave must be removed a little from the large one. Because while a remote object approaches, its image  $TS$  will also approach towards  $t^g$ ; and while  $tT$  is diminished, its reciprocal  $tq$  will be increased<sup>h</sup>.

176. Therefore to fit this telescope for a short-sighted person, since the eye-glass is usually fixt, the little speculum must be moved somewhat nearer to the large one. For then the interval  $tT$  will also be diminished and its reciprocal  $tq$  will be increased; and so the rays will fall upon the eye-glass diverging from a nearer point than its focal distance; and consequently will emerge from it diverging upon the eye.

A more general rule for the magnifying power.

177. By a farther contraction of the interval between the concaves, the image  $pq$  may be projected through the hole in the large concave,









concave, to any given place behind it; and by removing the eye-glass to the same distance from the images as before, the vision would become distinct again; and the object would be more magnified than before, as much as the ratio of  $tq$  to  $te$  or  $tc$  is made bigger than the ratio of  $TC$  to  $tc$ ; as appears by the demonstration<sup>a</sup>. But by<sup>a</sup> Art. 174. enlarging the image  $pq$ , it becomes more obscure and imperfect, and consequently the appearance of the object less bright and distinct. Besides, as the image becomes larger, the less of it, and of the object, can be seen at one view through a given eye-glass.

178. All things being fixt in their places, the diameter of an object taken in at one view, is proportionable to the breadth of the eye-glass, if the hole in the large concave does not limit it. For the angle of reflection  $pce$ , at the middle point of the lesser concave, being equal to the angle of incidence  $ecS$ ; it appears, that while  $pq$  and  $kl$  are increased or diminished in any ratio, the image  $ST$  and the object  $PQ$  will also be increased or diminished in the same ratio.

179. Now if an eye-glass of a given focal distance and convexity, be made very broad, it will become too thick; and so the rays will fall too obliquely upon one or both its surfaces near the margin of it; and this obliquity will cause too many of them to be reflected, and the rest that are transmitted, to be too much refracted, in comparison to those pencils that pass through the middle of the said lens. Therefore to increase the visible area of the object, it is necessary to project the image  $pq$  two or three inches beyond the hole in the large concave, and to intercept the rays that are tending towards it, with a thinner and broader convex glass  $fg$  put close to the backside of this concave; which glass will cause the rays to converge quicker than before, and to form an image  $vx$  nearer to it, and smaller, than  $pq$ ; both being terminated by a line  $pvg$  drawn through the center of this glass. And then the rays of each pencil diverging from this new image  $vx$ , must be received by another convex eye-glass  $bi$ , that shall make them emerge towards the eye in parallel lines. A meniscus glass, whose convex side is placed towards the converging pencils  $fvb$ , is fittest for this purpose; because the rays will pass through its edges less obliquely, than through a glass of any other shape.

180. To prevent collateral rays, that pass by the sides of the smaller concave, through the hole in the larger, and those also which are reflected from the imperfect margins of them both, from entering into the eye; it is necessary to place a thin plate with a proper hole in it to circumscribe the image at  $x$ , and also another very small

K

hole

The visible area is as the breadth of the eye-glass.

Is enlarged by two eye-glasses.

Fig. 168, 169.

The eye-stops.



hole at  $o$ , where all the pencils cross one another immediately before they enter the eye. The breadth of this latter hole must be no bigger than that of the principal pencil at  $o$ , and the places of them both must be exactly adjusted; otherwise the telescope can have no good effect.

The little concave speculum may be changed for a convex one. Fig. 170.

181. Telescopes of this kind are sometimes made with a little convex speculum instead of the concave one. If their focal distances be equal, and the vertex of the convex  $de$ , be placed at  $e$ , where the center of the concave was, the telescope will magnify in the same ratio as before; but will shew the object inverted; unless it be set upright by three convex eye-glasses, as in a dioptrick telescope. For a pencil of rays converging from the large concave towards its focus  $T$ , being intercepted by the little convex  $de$ , will be reflected by it to the same point  $q$  as before by the little concave  $bc$ . For the point  $t$  being the principal focus of both these little speculums, we have  $tT, te$  (or  $tc$ ) and  $tq$  continual proportionals as before<sup>a</sup>. Through any point  $S$  of the first image  $ST$  and through the center  $e$  of the little concave, draw  $Sep$  terminating the image  $pq$  formed by this concave<sup>b</sup>; in like manner through  $c$  the center of the little convex  $de$ , and through the same point  $S$ , draw  $cSr$  terminating the image  $qr$  formed by this convex. These images  $qp, qr$  lye on contrary sides of the axis, and therefore the object appears in contrary positions. But these images are equal, and of consequence the object appears equally magnified. For we have  $tq : te :: te : tT :: tq = te : te = tT$ , that is  $:: eq : eT :: cq : cT$ . And the triangles  $peq, Tes$  being similar, and also  $qcr, Tcs$ , we have  $pq : ST (:: eq : eT :: cq : cT ::) qr : ST$ ; therefore  $pq$  is equal to  $qr$ .

Double microscope considered. Fig. 171.

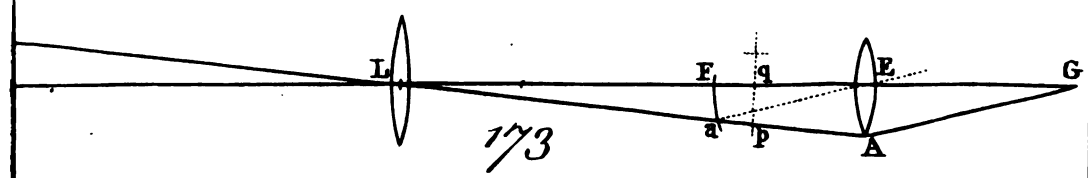
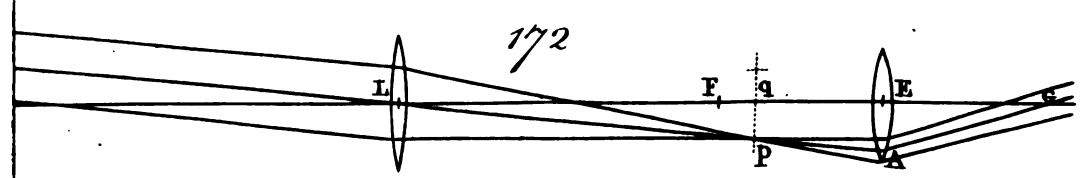
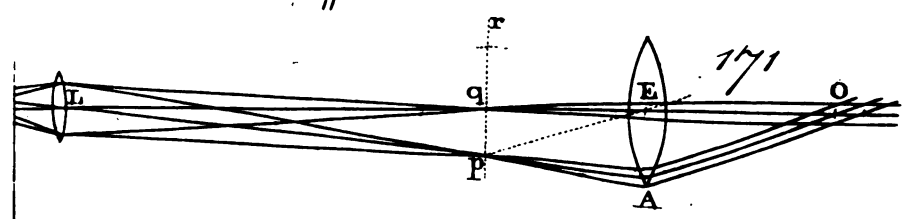
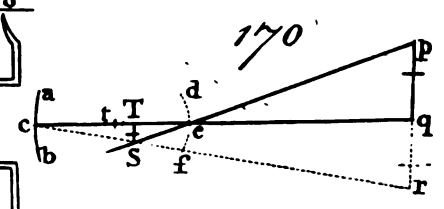
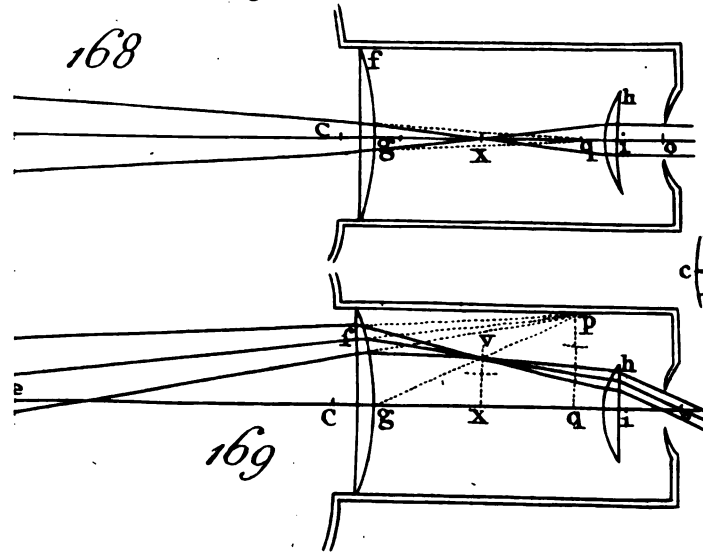
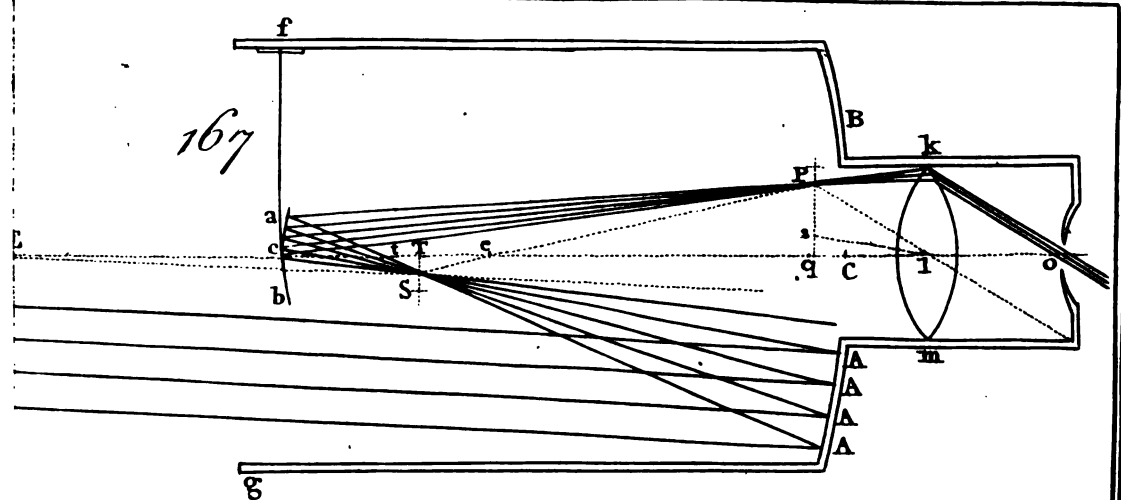
182. A double microscope is composed of two convex glasses placed at  $E$  and  $L$ . The glass  $L$  next the object  $PQ$  is very small and very much convex, and consequently its focal distance  $LF$  is very short; the distance  $LQ$  of the small object  $PQ$  is but a little greater than  $LF$ ; so that the image  $pq$  may be formed at a great distance from the glass<sup>c</sup>, and consequently may be much greater than the object itself<sup>d</sup>. This picture  $pq$  being viewed through a convex eye-glass  $AE$ , whose focal distance is  $qE$ , appears distinct as in a telescope. Now the object appears magnified upon two accounts; first because if we viewed its picture  $pq$ , with the naked eye, it would appear as much greater than the object, at the same distance, as it really is greater than the object, or as much as  $Lq$  is greater than  $LQ$ <sup>e</sup>; and secondly because this picture appears magnified through the eye-glass as much as the least distance at which it can be seen distinctly with the naked eye, is greater than  $qE$ , the focal

<sup>a</sup> Art. 108.

<sup>b</sup> Art. 115.

<sup>c</sup> Art. 115.









focal distance of the eye-glass<sup>a</sup>. For example, if this latter ratio be<sup>a</sup> Art. 162.  
5 to 1, and the former ratio of  $Lq$  to  $LQ$  be 20 to 1, then upon both  
accounts the object will appear 5 times 20, or 100 times greater  
than to the naked eye.

183. To fit these telescopes and microscopes to short-sighted eyes, To fit tele-  
the glasses  $E$  and  $L$  must be placed a little nearer together; so that scopes and  
the rays of each pencil may not emerge parallel but may fall di- microscopes  
verging upon the eye<sup>b</sup>; and then the apparent magnitude will be to defective  
eyes.  
altered a little but scarce sensibly; as is demonstrated in the next<sup>b</sup> Art. 104.  
article. 108. 111.

184. Suppose the interval  $LE$  between the two convex glasses to A more gene-  
be greater or less than the sum of their focal lengths, and let  $EF$  be ral demon-  
the focal distance of the eye-glass, and  $Lq$  that of the object-glass; stration of  
I say the apparent magnitude will be to the true as  $LF$  to  $FE$ , that telescopes.  
is, as the interval of the glasses diminished by the focal distance of Fig. 172, 173,  
the eye-glass, to the focal distance of the eye-glass. For the axes of 174, 175.  
all the pencils which pass through  $L$ , as  $PLA$ , will be refracted by  
the eye-glass to a focus  $G$ , where the eye being placed will see the  
whole object  $PQ$ , though the aperture of the pupil and of the object-  
glass be never so small<sup>c</sup>; and the object  $PQ$  will appear under the<sup>c</sup> See note to  
angle  $AGE$ . But  $L$  being a focus of incident rays upon the eye- Art. 164.  
glass, we have  $LF : LE :: LE : LG^d$ , and disjointly  $LF : FE ::^d$  Art. 107.  
( $LE : EG ::$  the angle  $EGA$ , to the angle  $ELA^e$  or  $PLQ^f$ ) as the<sup>e</sup> Art. 134.  
apparent magnitude to the true. <sup>f</sup> Art. 61.

185. Hence according as the interval of the glasses is greater or  
less than the sum of their focal distances, the apparent magnitude is  
to the true, in a greater or less proportion than that of the focal  
distances.

186. The brightness of the appearance through a given telescope The apparent  
or microscope is more or less in proportion to the aperture of the brightness  
object-glass. For supposing it covered with paper, all but a small through  
hole in the middle, the magnitudes of the pictures  $pq$  in the focus them.  
of the glasses, and of that upon the retina would not be altered; but  
the hole at  $L$  being smaller than before, there are fewer rays in every  
pencil, and consequently in every point of those pictures, and so they  
appear more obscure. If the aperture and object-glass remain the  
same, things appear brighter or fainter according as the focal di-  
stance of the eye-glass is longer or shorter; that is, according as the  
telescope or microscope magnifies less or more<sup>g</sup>. For the same<sup>g</sup> Art. 164.  
quantity of light spread over a smaller or larger picture or part of 182.  
the retina will make it brighter or duller.



All appearances the same when the eye is out of the axis of the glass or glasses.  
Fig. 176.

187. Hitherto I have supposed the eye to be always placed at some point  $O$  in the common axis of the refracting or reflecting surfaces. Now let it be placed at any point  $o$  in a line  $Oo$  perpendicular to the axis  $Qq$ ; I say that all the appearances will be the same or at least not sensibly different from what they were before. For let  $pq$  be the last image of an object, and  $PQ$  the last but one, or the object itself; draw the lines  $po, qo$  meeting the next surface in  $a$  and  $c$ ; and the points  $P$  and  $Q$  will appear to the eye at  $o$  in the directions of those lines  $oa, oc$ . Whence drawing  $pO$  meeting the surface in  $A$ ; since the directions  $OA, oa$ , in which  $P$  is seen, lye the same way from the directions  $OC, oc$ , in which  $Q$  is seen, it is evident that the apparent situation of the extremities  $P, Q$  is the same at both places of the eye; and also the apparent magnitude, which is measured by the angle  $aoc^a$  or  $poq$  or  $pOq$  or  $AOC$ . For the small angles  $poq, pOq$ , being subtended by the same image  $pq$ , very nearly at equal distances  $po, qO$  from  $o$  and  $O$ , are very nearly equal. The apparent brightness of the object is also the same; because the density of the rays, that enter the pupil, at any part of the perpendicular plane represented by  $Oo$ , is nearly the same<sup>b</sup>. For the rays flow from or towards the last image  $pq$  just as if it was a luminous body. And lastly the degree of apparent distinctness or confusion is the same also, because the angles which the pupil, placed at  $O$  or at  $o$ , subtends at  $p$  and  $q$ , or the mutual inclinations of the rays in each pencil are very nearly equal.

A general observation upon vision.

<sup>c</sup> Art. 160.

<sup>d</sup> Art. 148.

<sup>e</sup> Art. 147.

The portable camera obscura.  
Fig 177.

188. This general observation upon vision is worth remembering. That the apparent distinctness and confusion of an object depends upon the mutual inclination of the rays to each other in any one pencil when they fall upon the eye<sup>c</sup>; the apparent magnitude, upon the inclination of the rays of different pencils to each other when they fall upon the eye<sup>d</sup>; the apparent situation, upon the real situation of the extream pencils when they fall upon the eye<sup>e</sup>; and the apparent brightness and obscurity, upon the quantity of rays in every pencil.

189. The portable *camera obscura*, or dark chamber, commonly sold in the shops, will require but a short description. The theory is this; the rays that come from the object  $PQR$  after passing the lens  $E$  are tending to form an image  $pqr$ ; but being reflected upwards by the looking-glass  $ABC$ , they form an horizontal image  $wxq$  upon a glass plane, whose unpolished side lyes uppermost; upon which a copy of the picture may be sketched out with a black lead pencil; and to the spectator facing the object, the picture appears upright.



upright. This figure represents a section of the machine through the axis of the tube that holds the lens, and through the middle of the square box and the looking-glass within it. The section of the side opposite to the tube is not here represented, it being a door that opens sideways; the edges of the rough glass at the top are placed in two grooves upon the sides of the box; and being taken off, it is placed in a drawer *ef* at the bottom of the box; the looking-glass *ABC* may also be drawn out of the grooves in the sides of the box and lodged in the same drawer. The square wooden tube consists of 3 parts; the innermost that carries the lens, draws outwards or inwards to make the pictures distinct. The parts *gb* and *ik*, being fixt together and to the box with small bolts, may be taken asunder and put into the box; then the lid *at* at the top, and the door at the end, being both shut and fixt, the machine becomes more commodious for carriage. The inside of the lid whose section is *at*, has two wings, that open to right angles on each side of it, and rest upon the sides of the box, to shade the image upon the rough glass.

190. The construction of the magick lantern is briefly this; *ABCD* is a tin lantern, from whose side there proceeds a square or round arm or tube *bnkclm*, consisting of two parts; the outermost whereof *nklm* slides over the other, so as that the whole tube may be lengthened or shortened thereby. In the end of the arm *nklm* is fixt a convex glass *kl*: about *de* there is a contrivance for admitting and placing an object *de* painted in dilute and transparent colours on a plane thin glass; which object is there to be placed inverted. This is usually some ludicrous or frightful representation, the more to divert the spectators: *bhc* is a deep convex glass, so placed in the other end of the prominent tube, that it may strongly cast the light of the flame *a* on the picture *de* painted on the plane thin glass. And here it is to be noted, that the glass *bhc* is only designed for the strong illumination of the picture *de*, and has nothing to do in the representation; and therefore in some of these lanterns, instead of the glass *bhc*, we shall find a concave speculum so placed, that it may strongly cast the light of the flame *a* on the picture at *de*; and sometimes both are used.

191. Wherefore let us now consider the picture *de* as a very light-some object of distinct colour and parts. And let us conceive *de* more remote from the glass *kl* than its focus. It is then manifest that the distinct image of the object *de*, shall be projected by the glass *kl* on the opposite white wall *FH* at *fg*; and here it shall be represented erect. For now the whole chamber *EFGH* is dark, the lantern

Description  
of the magick  
lantern by  
Mr. Molyneux.  
Fig. 178.

<sup>a</sup> Art. 104.  
208. 115.

lantern  $ABCD$  inclosing all the light; so that in effect this appearance of the magick lantern is no more than what we are told concerning the representation of outward objects in a dark room by a convex glass; and here we may observe, that if the tube be contracted, and thereby the glass  $kl$  brought nigher the object  $de$ , the representation  $fg$  shall be projected so much the larger; and so much the more distant from the glass  $kl$ ; according to the rules before laid down<sup>a</sup>. So that the smallest picture at  $de$  may be projected at  $fg$  in any greater proportion required, within due limits. From whence the name of *Lanterna Megalographica*. And consequently, protracting the tube and drawing the glass  $kl$  more distant from the object  $de$ , will diminish the representation  $fg$ , and project it nigher the glass  $kl$ .

A comparison  
of different  
ways of illu-  
minating mi-  
croscopical  
objects, pi-  
ctures in a  
magick lan-  
tern, &c.  
Fig. 179, 180,  
181.

192. Of a luminous object  $QR$  let  $qr$  be the image formed by reflection from a concave surface, or by refraction through a convex lens, or sphere  $AC$ ; whose center is  $E$ , principal focus  $F$ , axis  $QEF$ , and semiaperture  $AC$ ; and let a perpendicular  $FG$ , to the axis, cut the outermost ray  $QA$  in  $G$ ; I say the brightness of the several pictures  $qr$ , will be very nearly as  $\overline{FG}^2$  directly and  $\overline{FE}^2$  inversely.

For, not regarding the small losses of light by the several reflections and refractions, the quantity collected to the point  $q$  is very

<sup>b</sup> Art. 25. nearly as  $\frac{\overline{AC}^2}{CQ^2}$ <sup>b</sup>, and consequently the quantity in the area of the

whole picture  $qr$ , as  $\frac{\overline{AC}^2}{CQ^2} \times \overline{QR}^2$  or  $\frac{\overline{FG}^2}{FQ^2} \times \overline{QR}^2$ . But the area of

the picture is as  $qr^2 = \frac{\overline{Eq}^2}{EQ^2} \times \overline{QR}^2 = \frac{\overline{FE}^2}{FQ^2} \times \overline{QR}^2$ . Because, in the

<sup>c</sup> Art. 71. reflecting concave, we have  $Fq : FE :: FE : FQ$ <sup>c</sup>; and consequently  $Eq : EQ :: FE : FQ$ ; and in the lens and sphere we have  $Qq : QE$

<sup>d</sup> Art. 107.  $:: QE : QF$ <sup>d</sup>, and consequently  $Eq : EQ :: FE : FQ$ . Therefore the brightness of the picture, or the density of the rays in its area,

being as their quantity directly and the area inversely, is as  $\frac{\overline{FG}^2}{\overline{FE}^2}$

very nearly; and the more exactly as the aperture is smaller and the object farther off.

193. *Corol. 1.* In a given speculum, lens or sphere, the brightness of the picture of a given object, is as  $\overline{FG}^2$ ; and therefore increases continually with the distance of the luminous object from the focus  $F$ .

194. *Corol.*



194. *Corol.* 2. If the luminous object be very remote, and the apertures of several specula, lenses and spheres be equal to one another, the degrees of brightness of the several pictures formed by them are reciprocally as the squares of their respective focal distances very nearly.

195. *Corol.* 3. Consequently if the several apertures be equal portions of equal spheres, the degrees of brightness of the several pictures formed by a concave speculum, a double-convex glass, a glass sphere and a plano-convex glass, are respectively as the squares of the decreasing musical progression 12, 6, 4, 3. Because the respective focal distances are  $\frac{1}{2}$ ,  $\frac{1}{3}$ ,  $\frac{1}{4}$ ,  $\frac{1}{6}$  of the diameter of the given sphere, by Art. 69, 103, 95; and the reciprocals of an arithmetical progression are called a musical progression.

196. *Corol.* 4. Therefore the concave speculum has greatly the advantage of the sphere and lenses, for illuminating microscopical objects, and also for burning things in the sun-shine, though not in so great a proportion, as will appear by the next corollary.

A comparison of the burning powers of glasses and speculums,

197. *Corol.* 5. Though the rays in two pictures of the sun formed by similar specula, be equally dense<sup>a</sup>, yet the picture formed by the larger speculum, being proportionably larger, will burn things more vehemently than the smaller; because the burning particles of matter communicate and propagate their heat to one another. And when the specula are similar, the aberrations of the rays from the peripheries of the pictures, are also similar<sup>b</sup>.

<sup>a</sup> Art. 192.

<sup>b</sup> Art. 213.

## CHAP. XI.

TO DETERMINE THE ABERRATIONS OF RAYS, FROM THE GEOMETRICAL FOCUS, CAUSED BY THEIR UNEQUAL REFRACTIBILITY, AND ALSO BY THE SPHERICALNESS OF THE FIGURE OF REFLECTING AND REFRACTING SURFACES.

## PROPOSITION I.

198. **L**ET the common sine of incidence be to the sine of refraction of the least refrangible rays as  $I$  to  $R$ , and to the sine of refraction of the most refrangible rays as  $I$  to  $S$ ; and the diameter of the least circular space into which heterogeneous parallel rays can be collected, by a spherical surface or by a plano-convex lens, will be to the diameter of its aperture, in the constant ratio of  $S - R$  to  $S + R - 2I$ .

Fig. 182.

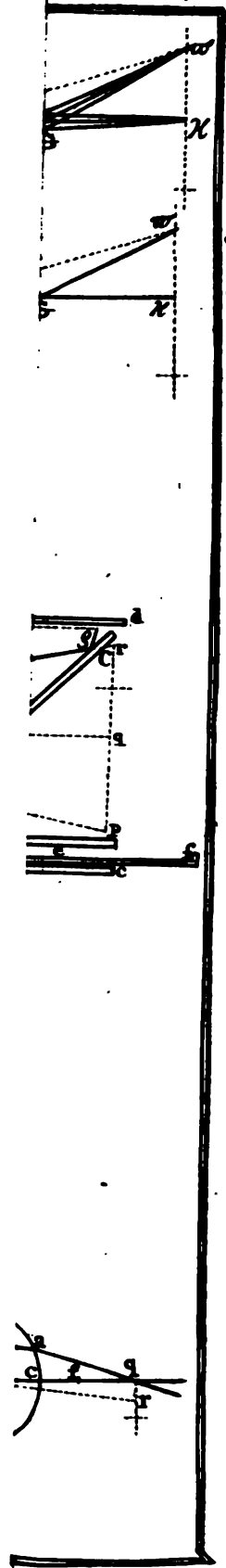
For let an heterogeneous ray  $PA$  fall upon a spherical surface  $ACB$ , and let it be separated by refraction into the rays  $AF$ ,  $Af$ , cutting the axis  $EC$ , drawn parallel to  $PA$ , in  $F$  and  $f$ . Take the arch  $CB$  equal to  $CA$ , and let another heterogeneous ray  $PB$ , coming parallel to  $PA$  be refracted into the lines  $BF$ ,  $Bf$ , cutting the two former rays in  $R$  and  $S$ . Join  $RS$  and produce it till it meets the incident rays produced in  $I$  and  $K$ , and the perpendiculars  $EA$ ,  $EB$  to the refracting surface at the points  $A$ ,  $B$ , in  $H$  and  $L$ . And when  $AB$ , the breadth of the aperture or of the pencil, is but moderate, and consequently the refractions at  $A$ ,  $B$  but small, the angles of incidence and refraction,  $HAI$ ,  $HAR$ ,  $HAS$ , or the arches that measure them, or their perpendicular subtenses  $HI$ ,  $HR$ ,  $HS$ , will be to each other very nearly in the same given ratios as those of the sines  $I$ ,  $R$ ,  $S$  of those angles<sup>a</sup>. And disjointly the differences of those subtenses will be proportionable to the differences of these sines: that is, the line  $RS : RI :: S - R : R - I$ , and doubling the consequents,  $RS : 2RI$  or  $IK - RS :: S - R : 2R - 2I$ ; and conjointly  $RS : IK$  or  $AB :: S - R : S + R - 2I$ . From this given ratio of  $RS$  to  $AB$  in which they increase or decrease together, it appears that all the intermediate rays which fall upon  $AB$  will pass through  $RS$ . And when parallel rays fall perpendicularly upon the plane side of a plano-convex lens, they are refracted only at their emergence from its convex surface; and so the aberrations are the same in both cases. Q. E. D.

<sup>a</sup> Art. 68. 84.

199. Corol.



Plate 14.







199. *Corol. 1.* Hence the diameter  $RS$ , of the circle of aberrations that contains all the incident rays, is a 55th part of the diameter  $AB$  of the aperture of a plano-convex glass, whatever be its focal distance. For supposing  $AR$  and  $AS$  to be the outmost red and indigo rays, their sines of incidence and refractions  $I, R, S$  are to each other as 50, 77, 78<sup>a</sup>. Whence  $S - R$  is to  $S + R - 2I$  as <sup>a</sup> Art. 31. 1 to 55.

200. *Corol. 2.* The diameter of the least circle that can receive the rays of any single colour or of several contiguous colours is also determinable from the proportions of their sines. Thus all the orange and yellow is contained in a circle whose breadth is the 260th part of the breadth of the aperture of the plano-convex glass; the sines of the outmost orange  $AR$  and yellow  $AS$  being to the common sine of incidence as  $77\frac{1}{2}$  and  $77\frac{1}{3}$  to 50<sup>b</sup>. <sup>b</sup> Art. 31.

201. *Corol. 3.* In different surfaces, or plano-convex glasses, the angles of aberration  $RAS$  are as the breadths of the apertures  $AB$  directly and as the focal distances  $CF$  inversely; because any angle, as  $RAS$ , is as its subtense  $RS$  directly and as its radius  $AR$  or  $CF$  inversely.

202. *Corol. 4.* If several glasses of several sorts or shapes have the same focal distance, and the same aperture; the diameter of the circle of aberrations of heterogeneous parallel rays from their principal focus, will be the same in them all; being the same as in a plano-convex glass when its plane side is turned to the incident rays<sup>c</sup>; and <sup>c</sup> Art. 57. 58. may therefore be determined by art. 198. And when the rays in the Fig. 182. incident pencil are either parallel or inclined to the axis of the lens, the diameter of the circle of aberrations is as its distance from the lens; because the angle  $RAS$  is invariable.

203. *Corol. 5.* Therefore with respect to these aberrations by colours separately considered, it is indifferent which side of the lens is turned to the incident rays; because its focal distance is the same in both positions<sup>d</sup>. <sup>d</sup> Art. 101.

#### LEMMA.

204. *The versed sines  $AB, AC$  of very small arches  $BD, CD$ , of Fig. 183. 184. unequal circles  $BDG, CDH$ , that have the same right sine  $AD$ , are reciprocally proportionable to their diameters  $BG, CH$  very nearly; that is  $AB : AC :: CH : BG$ .*

For since the rectangles under  $BAG$  and  $CAH$  are each equal to the square of  $AD$ <sup>e</sup>, and consequently to each other; their sides are <sup>e</sup> Euc. VI. reciprocally proportionable<sup>f</sup>, that is  $AB$  is to  $AC$  as  $AH$  to  $AG$  or <sup>f</sup> Euc. VI. <sup>13.</sup> as <sup>14.</sup>

L



as  $CH$  to  $BG$  very nearly, when the versed sines are incomparably less than the diameters themselves<sup>a</sup>.  $\mathcal{Q}$ .  $E. D.$

## PROPOSITION II.

Fig. 185.

205. When homogeneous parallel rays  $NA, EC$  fall upon a spherical surface  $AC$  whose center is  $E$ , the longitudinal aberration  $FT$  of any refracted ray  $AT$  from  $F$  the focus of the pencil, is to the versed sine of the arch  $AC$  intercepted between the point of incidence and the axis  $ECF$ , in the given ratio of the square of the sine of refraction, to the rectangle under the sine of incidence and the difference of the sines very nearly; and the aberration is the same when the rays fall perpendicularly upon the plane side of a plano-convex lens.

<sup>b</sup> Art. 93.<sup>c</sup> Art. 92.<sup>d</sup> Art. 204.<sup>e</sup> Art. 68.

For when the refraction is made in the passage of a ray  $NA$  from a denser to a rarer medium, the intersection  $T$  of the refracted ray  $AT$  with the axis  $ECF$ , lies between the refracting surface and its focus  $F$ . With the center  $T$  and semidiameter  $TA$  having described the arch  $AD$  cutting the axis in  $D$ , draw the sine  $AP$  of the arches  $AC, AD$ , and also  $EN$  and  $EM$  the sines of incidence and refraction, for which put  $n$  and  $m$ ; then because the triangles  $ETM, ATP$  are similar, it will be as  $ET : TA$  or  $TD :: (EM : AP$  or  $EN ::)  $EF : FC$ <sup>b</sup>; and disjointly  $TF : EF :: (FC - TD$  or)  $TF - CD : FC$ ; and alternately  $TF : TF - CD :: EF : FC$ ; and disjointly  $TF : CD :: (EF : EC ::)  $m : m - n$ <sup>c</sup>. Again since  $(PD : PC ::)  $CE : DT$ <sup>d</sup> or  $FC$ <sup>e</sup>, and conjointly)  $CD : CP :: (EF : FC ::)  $m : n$ ; by compounding this and the foregoing proportion, it will be as  $TF : CP :: mm : m - n, n$ .  $\mathcal{Q}$ .  $E. D.$$$$$

206. Corol. 1. The segment  $ACBPA$  may be considered as a plano-convex lens; and when rays fall parallel upon its plane side, the longitudinal aberration of the extreme ray falling upon  $A$  is equal to  $\frac{2}{3}$  of its thickness  $PC$ , as appears by putting 3 and 2 for  $m$  and  $n$  respectively.

207. Corol. 2. Also this aberration  $FT = \frac{mm}{m - n, n} \times \frac{AP^2}{2EC} =$

<sup>f</sup> Euc. VI. 13.  $\frac{mm}{m - n} \times \frac{AP^2}{2CF}$ . For  $PC = \frac{AP^2}{2EC}$  very nearly<sup>f</sup>, and  $EC = \frac{m - n}{n}$   
<sup>g</sup> Art. 92.  $\times CF$ <sup>g</sup>.

208. Corol. 3. Let the refracted ray  $ATG$  produced, cut the line  $FG$ , perpendicular to the axis, in  $G$ , and the lateral aberration  $FG$



$$FG = \frac{mm}{nn} \times \frac{AP^3}{2EC^2} = \frac{mm}{m-n} \times \frac{AP^3}{2CF^2}. \text{ For } FG : TF :: AP :$$

$$TP \text{ or } CF \text{ or } \frac{n}{m-n} \times CE.$$

209. *Corol. 4.* When the semidiameter of the convexity or the focal distance is given, the longitudinal aberrations are as the squares, and the lateral aberrations as the cubes, of the linear apertures of a plano-convex lens.

## PROPOSITION III.

210. *When parallel rays QA, EC are reflected from a spherical concave ACB whose center is E and whose aperture ACB is but small, the longitudinal aberration TF of the extremity ray AT from the geometrical focus F, is equal to half the versed sine CP of the semiaperture AC very nearly.* Fig. 186.

In fig. 185. imagine  $EM$ , the sine of refraction to be diminished to nothing, and then to become negative and equal to  $EN$  the sine of incidence, and the refraction of the ray to be changed to reflection as in fig. 186; and by the former proposition it will be as  $TF : CP :: mm : -m-n, n :: nn : -2nn :: 1 : -2$ .

But the particular proof is this. By the last lemma the versed sine  $CP$  nearly equals  $\frac{1}{2}$  the versed sine  $PD$  of the arch  $AD$  whose center is  $T$  and semidiameter  $TA$  or  $TE$  or  $\frac{1}{2}$  the semidiameter of the arch  $AC$  very nearly. But  $2TF = 2TE - 2EF = ED - EC$  Art. 69,  $= CD$  exactly or  $CP$  nearly. Therefore  $TF = \frac{1}{2}CP$  nearly.

211. *Corol. 1.* We had  $2TF = CD$  exactly; which is the excess of the secant  $ED$  of the arch  $AC$  above its radius  $EA$ . For joining  $AD$  the angle  $DAE$  in the semicircle  $DAE$  is a right one.

212. *Corol. 2.* The longitudinal aberration  $TF = \frac{AP^2}{4CE}$ . For  $CP = \frac{AP^2}{2CE}$  nearly<sup>b</sup>.

<sup>b</sup> Euc. VI. 13.

213. *Corol. 3.* The lateral aberration  $FG = \frac{AP^3}{2CE^2}$ . For  $FG : FT :: AP : PT$  or  $\frac{1}{2}CE$  nearly.

214. *Corol. 4.* When the diameter of the concave, or its focal distance, is given, the longitudinal aberrations are as the squares, and the lateral ones as the cubes of the diameters of the apertures.

## PROPOSITION IV.

215. *When parallel rays of any one sort are refracted by a plano-convex object-glass, or when rays of all sorts are reflected by a spherical concave, the diameter of each circle of aberrations caused by the sphericalness of the figures, is equal to  $\frac{1}{2}$  the lateral aberration of the extream ray in each; and therefore is given by the former propositions.*

Fig. 187, 188. Let  $\alpha Y\tau$  be any refracted or reflected ray cutting the axis  $ECT$  in  $\tau$ , and the extream ray  $ATG$ , that comes from the contrary side of the axis, in  $Y$ . Draw  $YX$  perpendicular to the axis, and supposing the line  $ATG$  immoveable, as the point of incidence  $\alpha$  moves from the vertex  $C$ , the perpendicular  $XY$  will first increase; because the angle  $C\tau\alpha$  continually increases, and afterwards will decrease, because the line  $T\tau$  continually decreases; and when  $XY$  is the greatest, it is evident that all the rays, incident upon the same side of the axis as itself, will pass through it. To find its greatest quantity, let the incident ray  $q\alpha$  cut the chord  $APB$  in  $\beta$ , and supposing the variable aperture  $P\beta = v$ , the variable  $TX = x$  and the given lines  $PA = a$ ,  $PT = f$ ,  $TF = b$ ; by cor. 4. prop. 2 and 3, the aberration  $F\tau$  is to the aberration  $FT (b)$  as  $\pi\alpha^2$  or  $P\beta^2 (vv)$  to  $PA^2 (aa)$ . Wherefore  $F\tau = \frac{vv}{aa}b$  and thence  $TF - F\tau = T\tau = \frac{b}{a} \times \frac{aa - vv}{a}$ . Again  $PT (f) : PA (a) :: TX (x) : XY = \frac{ax}{f}$ ; also  $\pi\alpha (v) : \pi\tau$  or  $PT (f) :: XY \left( \frac{ax}{f} \right) : X\tau = \frac{ax}{v}$ . Hence again  $T\tau$ , or  $X\tau + XT = \frac{ax}{v} + x = \frac{b}{aa} \times \frac{aa - vv}{a}$  found before; or  $\frac{x}{v} \times \frac{aa - vv}{a} = \frac{b}{aa} \times \frac{aa - vv}{a}$ . Whence  $x = \frac{b}{aa}v \times \frac{aa - vv}{a}$ , and therefore  $x$  or  $TX$  is the greatest possible when the rectangle  $v \times a - v$ , or  $P\beta \times \beta B$  is greatest, that is when its sides <sup>a</sup>  $P\beta$ ,  $\beta B$  are equal<sup>a</sup>, or when  $v = \frac{1}{2}a$ . Substitute this value for  $v$  in the last equation and it gives the greatest value of  $x = \frac{1}{4}b$ , or the greatest  $TX = \frac{1}{4}TF$ , and therefore the greatest  $XY = \frac{1}{4}FG$ , because  $TX : XY :: TF : FG$ , and this  $XY$  turned about the axis.  $PX$  describes the circle of aberrations through which all the rays, falling up  $AB$  will just pass. Q. E. D.

<sup>a</sup> Euc. VI. 13.  $P\beta$ ,  $\beta B$  are equal<sup>a</sup>, or when  $v = \frac{1}{2}a$ . Substitute this value for  $v$  in the last equation and it gives the greatest value of  $x = \frac{1}{4}b$ , or the greatest  $TX = \frac{1}{4}TF$ , and therefore the greatest  $XY = \frac{1}{4}FG$ , because  $TX : XY :: TF : FG$ , and this  $XY$  turned about the axis.  $PX$  describes the circle of aberrations through which all the rays, falling up  $AB$  will just pass. Q. E. D.



## PROPOSITION V.

216. *If the density of the reflected rays in the circle of aberrations be uniform, it is to the density of the incident rays falling perpendicularly upon a plane AP, as the whole surface of the sphere of which the speculum is a portion, to the area of a circle whose diameter is the versed sine PC of the small arch AC very nearly, and the more exactly as this arch is smaller; supposing also that all the incident rays are reflected.* Fig. 183.

For since the very same rays pass through two circles described by the lines AP and XY turned about EC; their densities in these circles are reciprocally as the circles themselves; that is, the density of the reflected rays, is to the density of the incident rays, as  $AP^2$

to  $XY^2$ , or  $\frac{1}{16} FG^2$ , or  $\frac{AP^6}{16 \times 4 CE^4}$ ; that is, putting  $D$  for  $2CE$ , as  $4D^4$  to  $AP^4$ ; that is, as  $4D^2$  to  $PC^2$ , (because  $D, AP, PC$  are very nearly continual proportionals<sup>d</sup>;) that is, as the area of 4 great circles of the sphere, or the whole surface of the sphere<sup>e</sup>, to the area of a circle whose diameter is  $PC$  very nearly. <sup>a</sup> Euc. XII. 2. <sup>b</sup> Art. 215. <sup>c</sup> Art. 213. <sup>d</sup> Euc. VI. 8. <sup>e</sup> Cor. Archim. de Sph. & Cyl.

217. *Cor. 1.* Therefore the greatest density of the reflected rays is at the focus  $F$ , considered as a physical point; and is immensely greater than the density of the incident rays. For the proposition above becomes geometrically exact when  $AP$  is infinitely diminished, and  $XY$  comes to its limit at  $F$ ; and the density at  $F$  is always the same whether a slender pencil falls upon the speculum or a large one, because the outward rays are reflected wide of the focus  $F$ .

218. *Corol. 2.* In like manner when rays fall parallel upon the plane side of a plano-convex lens, (putting  $m$  to  $n$  for the ratio of majority of the sines of incidence and refraction) their greatest density at their focus  $F$ , is to the density of the incident rays, as the whole surface of a sphere whereof the lens is a portion, to the area

of a circle whose diameter is  $\frac{m}{n} PC$ , or in glass  $\frac{2}{3}$  of the versed sine of the smallest aperture of the lens; that is immensely great. It follows from art. 208.

219. *Cor. 3.* Therefore the density of reflected or refracted rays in the several points of an image of a very remote object, is also immensely greater than the density of the incident rays of any one pencil. For it would be immensely great, if all the rays of every pencil were rejected, except a few that go near to their axes, and those outward rays being scattered upon points collateral to each point.

point of the image, help to increase the density of the rays in the whole image.

### PROPOSITION VI.

220. *The circle of aberrations caused by the sphericalness of the figure of the object-glass of a telescope, compared with the circle of aberrations caused by the unequal refrangibility of rays, is altogether inconsiderable.*

Newt. Opt.  
p. 83.

For if the object-glass be plano-convex and the plane side be turned towards the object, and the diameter of a sphere whereof this glass is a segment be called  $D$ , and the semidiameter of the aperture of the glass be called  $S$ , and the sine of incidence out of glass into air be to the sine of refraction as  $n$  to  $m$ ; the rays which come parallel to the axis of the glass shall in the place where the image of the object is most distinctly made, be scattered all over a little circle

whose diameter is  $\frac{mm}{nn} \times \frac{S^3}{DD}$  very nearly, if they were all equally refrangible by article 215 and 208. As for instance, if the sine of incidence  $n$  be to the sine of refraction  $m$  as 20 to 31, and if  $D$ , the diameter of the sphere to which the convex side of the glass is ground, be 100 foot or 1200 inches, and consequently the telescope about 100 foot long<sup>a</sup>, and  $S$  the semidiameter of the aperture be

<sup>a</sup> Art. 92.

2 inches; the diameter of this circle of aberrations, that is  $\frac{mm}{nn} \times \frac{S^3}{DD}$ , will be  $\frac{31 \times 31 \times 8}{20 \times 20 \times 1200 \times 1200}$  or  $\frac{961}{72000000}$  parts of an inch.

<sup>b</sup> Art. 199.

But the diameter of the little circle through which these rays are scattered by unequal refrangibility, will be about the 55th part of the breadth of the aperture of the object-glass<sup>a</sup>, which is here 4 inches. And therefore the aberration arising from the spherical figure of the glass, is to the aberration arising from the different re-

frangibility, as  $\frac{961}{72000000}$  to  $\frac{4}{55}$ , that is as 1 to 5449; and therefore being in comparison so very little, deserves not to be considered in the theory of telescopes. If we suppose the little circle of aberrations arising from unequal refrangibility, to be 250 times narrower than the circular aperture of the object-glass, it would contain all the orange and yellow, and would permit the other fainter and darker colours to pass by it<sup>c</sup>, which perhaps may scarce affect the sense; yet even in this case the aberration caused by the spherical figure, would be to the aberration caused by the unequal refrangibility,

<sup>c</sup> Art. 200.  
Newt. Opt.  
p. 88.



lity, in a 100 foot telescope, but as  $\frac{961}{72000000}$  to  $\frac{4}{250}$ , or only as 1 to 1200, which sufficiently proves the proposition. *Q. E. D.*

221. *Corol. 1.* If the focal distances and apertures of a reflecting concave and a plano-convex glass be both the same, the diameter of the circle of aberrations, caused by their figures, will be above 30 times less in the reflecter than in the refracter. For these diameters

are  $\frac{AP^3}{16CF^2}$  and  $\frac{mm}{m-n} \times \frac{AP^3}{4CF^2}$  by art. 215, 213, and 208; which are as  $\frac{1}{4}$  to  $\frac{mm}{m-n}$  or  $\frac{31 \times 31}{11 \times 11}$ . Hence if the length of each tele-

scope be 100 foot, the lateral aberrations in the reflecter would be  $30 \times 5449$  or 163470 times less than the lateral aberrations caused by unequal refrangibility in the refracter.

222. *Corol. 2.* The number of pencils, some of whose rays are mixed together in every point of a confused picture, is as the area of the circle of aberrations of the rays in any one pencil; and consequently the mixture of the rays of different pencils, caused by the sphericity of the figure of an object-glass, if they were all alike refrangible, would be to their mixture caused by their unequal refrangibility, as 1 to  $5449 \times 5449$  or 29691601 in the present instance. For conceiving any point in the confused picture to be a center of a circle of aberrations, it is manifest that all other equal circles of aberrations, whose centers fall upon the first mentioned circle will cover its center; that is some rays of as many pencils will be mixed in this center as there are points in the circle itself; or, which is the same thing, the number of pencils mixed in this center is as the area of the circle of aberrations.

## CHAP. XII.

A REFRACTING OR REFLECTING TELESCOPE BEING GIVEN, WHOSE APERTURE AND EYE-GLASS ARE ADJUSTED BY EXPERIENCE, TO DETERMINE THE LENGTH, APERTURE AND EYE-GLASS OF ANOTHER TELESCOPE, THROUGH WHICH AN OBJECT SHALL APPEAR AS BRIGHT AND DISTINCT AS IN THE GIVEN ONE, AND MAGNIFIED AS MUCH AS SHALL BE REQUIRED.

## PROPOSITION I.

223. **I**N all sorts of telescopes and double microscopes, the apparent indistinctness of a given object, is as the area of a circle of aberrations in the focus of the object-glass directly, and as the square of the focal distance of the eye-glass inversely.

For in vision with the naked eye or with glasses, the apparent indistinctness of a given object, is as the area of a circle of aberrations in its picture painted upon the retina. Because any one sensible point of the retina, being the center of a circle of aberrations, will at once be affected by a mixture of the rays of as many distinct pencils, as there are sensible points in the area of that circle<sup>a</sup>; and so will at once convey to the mind a mixt or confused sensation of the same number of visible points in the object, from whence those pencils flowed; and this number of points is as the magnitude of the area of a circle of aberrations, whatever be the magnitude of a sensible point of the retina. Now in vision with telescopes, the diameter of a circle of aberrations in the picture upon the retina, is as the apparent magnitude of the diameter of the corresponding circle of aberrations in the common focus of the glasses<sup>b</sup>, that is as the angle subtended by this diameter at the center of the eye-glass<sup>c</sup>; that is as the diameter itself directly, and the focal distance of the eye-glass inversely<sup>d</sup>. And so the area of that circle of aberrations upon the retina, is as the area of the corresponding circle of aberrations in the focus of the object-glass directly, and as the square of the focal distance of the eye-glass inversely<sup>e</sup>.

224. *Corol.* In all sorts of telescopes and double microscopes a given object appears equally distinct, when the focal distances of the eye-

<sup>a</sup> Art. 222.

<sup>b</sup> Art. 136.

<sup>c</sup> Art. 164.

<sup>d</sup> Art. 86.

<sup>e</sup> Enc. XII. 2.



eye-glasses are as the diameters of the circles of aberrations in the focus of the object-glasses.

225. The alteration in the confusion which may arise from aberrations caused by the eye-glasses, is not here regarded, as being inconsiderable. We only consider the confusion of those points in the image which lye very near the axis of the telescope, as of the point *q* in fig. 162. Now if this point was perfectly distinct the rays going from it would emerge from the eye-glass in parallel lines without sensible error; because the breadth of this cylinder of rays is exceeding small compared to the breadth of the eye-glass, being in proportion to the breadth of the aperture of the object-glass as their focal distances; and the refractions at so small a distance from the axis are sufficiently true and regular. It is the largeness of the aperture of the object-glass and of its focal distance, which causes the irregularity in its refractions. Add to this that the differently refrangible rays cannot be separated sensibly in going so short a distance as between the eye-glass and the eye. Besides this we find by experience that objects and images distinct in themselves, appear sufficiently distinct through very small eye-glasses when their apertures are small.

### PROPOSITION II.

226. *In refracting telescopes the apparent indistinctness of a given object, is directly as the area of the aperture of the object-glass, and inversely as the square of the focal distance of the eye-glass.*

This appears from prop. 1, because the area of the circle of aberrations at the focus of the object-glass is as the area of its aperture<sup>a</sup>; and because the aberrations arising from the eye-glass<sup>b</sup>, and from the sphericity of the figure of them both are inconsiderable<sup>c</sup>.

<sup>a</sup> Art. 198.

<sup>b</sup> Art. 225.

<sup>c</sup> Art. 220.

227. *Corol.* In refracting telescopes a given object appears equally distinct, when the diameters of the apertures of their object-glasses, are as the focal distances of their eye-glasses.

### PROPOSITION III.

228. *In all sorts of telescopes and double microscopes the apparent brightness of a given object is as the square of their linear apertures directly and as the square of their linear amplifications inversely.*

For if the squares of the linear amplifications, that is if the areas of the pictures upon the retina were the same, their brightness

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would

would be as the quantities of light coming through the areas of the apertures, that is as the squares of the linear apertures; and if the apertures or quantities of light were the same, the brightness of the pictures would be as their areas inversely or as the squares of the linear amplifications inversely. Therefore when neither the apertures nor the amplifications are the same, the brightness is as the square of the linear apertures directly, and as the square of the linear amplifications inversely. *Q. E. D.*

229. *Corol. 1.* Hence in refracting and reflecting telescopes a given object appears equally bright, when their linear apertures are as their linear amplifications, that is as the focal distances of the object-glasses directly and as the focal distances of the eye-glasses inversely.

230. *Corol. 2.* If the breadth of the aperture of a given object-glass and the focal distance of the eye-glass be each increased in any given ratio, the distinctness will remain the same as before<sup>a</sup>; and the linear amplification will be diminished in the same ratio<sup>b</sup>; but the apparent brightness will be increased in a ratio quadruplicate of the former ratio by this proposition; and on the contrary.

<sup>a</sup> Art. 227.

<sup>b</sup> Art. 164.

Dioptr. p.  
215.

231. *Hugens* observes that the same degrees of distinctness here demonstrated do not exactly agree with experience, as he found by looking at the same object through different telescopes, or through the same telescope with different apertures; and that through the larger aperture the object appeared not quite so distinct as through the smaller. He found also that in viewing objects of different brightness through the same aperture, the apparent indistinctness of the brighter object was a little greater than that of the duller: and therefore the aperture adjusted for the duller planets may be somewhat larger than for the brighter.

#### PROPOSITION IV.

232. *In reflecting telescopes the apparent indistinctness of a given object is as the sixth power of the diameter of the aperture of the object-metal directly, and as the fourth power of its focal distance inversely, and also as the square of the focal distance of the eye-glass inversely.*

For the area of a circle of aberrations in the focus of the object-metal is as the sixth power of its linear aperture directly and as the fourth power of its focal distance inversely<sup>c</sup>; and therefore the apparent indistinctness of the object, is as the sixth power of the linear aperture

<sup>c</sup> Art. 215.  
213.



aperture directly, as the fourth power of the focal distance of the object-metal inversely, and as the square of the focal distance of the eye-glass inversely<sup>a</sup>. *Q. E. D.*

<sup>a</sup> Art. 223.

233. *Corol.* In reflecting telescopes a given object appears equally distinct when the cubes of the linear apertures of the object-metals, are as the solids whose bases are the squares of the focal distances of the object-metals, and heights are the focal distances of the eye-glasses: or when the focal distances of the eye-glasses are as the cubes of the linear apertures of the object-metals, applied to the squares of their focal distances.

### PROPOSITION V.

234. *In refracting telescopes of various lengths a given object will appear equally bright and equally distinct, when their linear apertures and focal distances of their eye-glasses are severally in a subduplicate ratio of their lengths or focal distances of their object-glasses: and then also their linear amplifications will be in a subduplicate ratio of their lengths.*

For to shew the object equally bright, the rectangle under the linear aperture and the focal distance of the eye-glass must be as the length of the telescope<sup>b</sup>, and to shew it equally distinct the linear<sup>b</sup> Art. 229. aperture must be as the focal distance of the eye-glass<sup>c</sup>; and there-<sup>c</sup> Art. 227. fore to perform both things together, the square of the linear aperture, and also the square of the focal distance of the eye-glass, must be severally (as the rectangle under each, or) as the length of the telescope; and consequently the linear aperture, and also the focal distance of the eye-glass, as the square root of that length. Now the linear amplification was as the linear aperture<sup>d</sup>, or by this demon-<sup>d</sup> Art. 227. stration, as the square root of the length of the telescope. *Q. E. D.*

235. *Hugens's* standard telescope 30 foot long, or 360 inches, *Diop.* p. 210. bears an aperture whose breadth is 3 inches, and an eye-glass whose focal distance is 3 inches and 3 tenths. From whence he has given us the following table of apertures and eye-glasses for other telescopes<sup>e</sup>, computed by the following rule.

<sup>e</sup> Art. 244.

Multiply the number of feet in the focal distance of any proposed object-glass by 3000, and the square root of the product will give the breadth of its aperture in hundredth parts of an inch. And the same breadth of the aperture, increased by a tenth part of itself, gives the focal distance of the eye-glass in hundredth parts of an inch. And the magnifying powers are as the breadths of the apertures.

For since the standard telescope has 30 foot focal distance of its

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object-

object-glass, put  $F$  for the number of feet in any other focal distance, and say by the proposition as  $\sqrt{30}$  to  $\sqrt{F}$ , so is the standard aperture 3 inches or 300 centesimals or  $\sqrt{300 \times 300}$ , to the aperture sought; which therefore is  $\sqrt{3000F}$  in centesimals of an inch. The focal distance of the eye-glass of the standard telescope is  $3\frac{1}{10}$  inches, that is a tenth part more than the breadth of the aperture of the object-glass; consequently the focal distance of the new eye-glass must be a tenth part more than the linear aperture of the new object-glass, by the last proposition.

Diop. p. 215. 236. He also adds the following directions how to suit these telescopes to all sorts of objects seen either by day or by night. They are proportioned in the following table for astronomical observations, and therefore will require more light when used in the day time. For when the eye is dazzled with the brightness of the day, objects will appear through them but obscure, which in the night are sufficiently bright. Therefore (says *Hugens*) when I used these telescopes to observe objects by day-light, by experience I found it requisite to change the eye-glasses for others whose focal distances were double the former. By this means the apparent brightness became quadruple, because the surfaces of the images in the bottom of the eye were diminished in the same proportion <sup>a</sup>. For as the aperture remains unaltered, so does the quantity of light, and therefore it illuminates a lesser space so much the more. Now if the aperture was increased without changing the eye-glass, the brightness would be increased too, but then the mist arising from greater aberrations would also be greater; and therefore this remedy must not be used.

<sup>a</sup> Art. 164.

237. But one may ask this question, since by substituting an eye-glass of a longer focal distance, the apparent indistinctness hitherto examined is diminished, why may not the aperture of the object-glass be so far increased, till the same degree of indistinctness returns again as belongs to a telescope regulated by the table? For from hence more light is gained and the distinctness is not altered <sup>b</sup>. The answer is this, which I hinted before <sup>c</sup>, that the mist arising from *Newton's* aberration, though the same in quantity, becomes more sensible in proportion to the brightness of the image. For the brightness of the mist increases at the same time. And we find by experience, that as soon as the apertures of those day-light telescopes are increased, the mist arising from the aberrations of a brighter object begins to be troublesome. The apertures therefore must not be altered.

<sup>b</sup> Art. 230.

<sup>c</sup> Art. 231.

238. Again one may ask, if a telescope fitted for Saturn be applied to the Moon, which is 100 times brighter (I mean in each equal parts



parts, though not in the whole, as being 10 times nearer to the Sun;) one may ask I say whether the breadth of the aperture and the focal distance of the eye-glass may not both be lessened in the same proportion to make the regions of the moon no brighter than those of saturn, but much greater in appearance than before. For instance, in a 30 foot telescope, if 3 inches, the breadth of the aperture be reduced to  $\sqrt{\frac{9}{100}}$  of an inch, which is somewhat less than ( $\sqrt{\frac{9}{9}}$  or) a third part of the former, and also the focal distance of the eye-glass be shortened in the same proportion; the proportion of the apparent brightness in these two telescopes, the object being the same, would be quadruplicate of 3 to  $\sqrt{\frac{9}{100}}$ <sup>a</sup> that is as 100 to 1; <sup>a</sup> Art. 230. and since the regions in the moon are 100 times brighter in themselves than those in saturn, the moon would appear in the darker telescope just as bright as saturn did in the lighter. But the apparent indistinctness hitherto considered would also be the same in both<sup>b</sup>, and the amplification of the moon would be greater than<sup>b</sup> Art. 230. that of saturn in the ratio of 3 to  $\sqrt{\frac{9}{100}}$ <sup>c</sup>, which is more than triple.<sup>c</sup> Art. 164. So that this reduction of the aperture and eye-glass seems very advantageous; but in reality it is quite otherwise; and that for two reasons. First because the minute parts of the moon may be better discerned when all the light remains in the telescope, than when it is reduced to an 100th part, though not in the same proportion. The other reason is that when the aperture is too much contracted, the out-lines that circumscribe the pictures in the eye become confused; which is carefully to be minded, and also what are the limits of this confusion. This is certain that as the aperture is contracted, the slender pencils or cylinders of rays, that emerge from the eye-glass into the eye, are also contracted in the same proportion. Now if the breadth of one of these pencils be less than  $\frac{1}{3}$  or  $\frac{1}{6}$  of a line, that is less than  $\frac{1}{60}$  or  $\frac{1}{120}$  part of an inch, the out-lines of the pictures are spoiled, for some unknown reason in the make of the eye, whether in the choroid, or in the retina, or in the humors it is uncertain. For by looking through an hole, in a thin plate, narrower than  $\frac{1}{3}$  or  $\frac{1}{6}$  of a line, the edges of objects begin to appear confused and so much the more as the hole is made narrower. Now it is easy to shew in the last mentioned telescope that the cylinder of rays is too slender. For by adding  $\frac{1}{10}$  of the aperture to itself<sup>d</sup>, the focal distance of the eye-glass becomes  $\sqrt{\frac{9}{100}} + \frac{1}{10} \sqrt{\frac{9}{100}}$ , that is  $\frac{11}{10} \sqrt{\frac{9}{100}}$  of an inch; and by similar triangles subtended at the common focus *q* by the aperture and cylinder sought<sup>e</sup>, it is as the focal distance of the object-glass, to the focal distance of the eye-glass, so the breadth of the aperture, to the breadth of the cylinder; that is as  
 as

<sup>a</sup> Art. 235.

<sup>e</sup> Fig. 162.

as 30 feet or 360 inches to  $\frac{11}{10} \sqrt{\frac{9}{10}}$  inches, so is  $\sqrt{\frac{9}{10}}$  of an inch to  $\frac{11}{4000}$  inch or almost  $\frac{1}{10}$  of a line; which is much less than  $\frac{1}{2}$ . But in the telescope regulated in the table, it is as 360 to  $3\frac{1}{10}$  so 3 to  $\frac{11}{400}$  of an inch or almost  $\frac{1}{10}$  of a line for the breadth of that cylinder; which can possibly do no harm. Hence we learn that the breadth of the aperture and focal distance of the eye-glass cannot be contracted much more than  $\frac{1}{10}$  of themselves; for even then the breadth of the cylinder at the eye will not much exceed  $\frac{1}{10}$  of a line. The same is to be understood of telescopes of all lengths regulated as in the table, the breadth of the cylinder being the same in all. For by the proportion just mentioned it equals the breadth of the aperture multiplied into the focal distance of the eye-glass and divided by the focal distance of the object-glass, and consequently it is proportionable to the linear aperture directly and the linear amplification inversely; which two ratios must compound a ratio of equality to preserve the same apparent brightness, by art. 229.

239. Hence though we transferred one of these telescopes from Saturn to Venus which is 225 times brighter, being 15 times nearer to the Sun, yet the breadth of the aperture must not be contracted above  $\frac{1}{10}$  part of the whole; and if too much light still remains, it must be diminished by darkening the eye-glass with the smoak of a candle. For a greater contraction of the aperture is hurtful for another reason, that all the little bubbles and veins in the eye-glass become more conspicuous by intercepting the whole or a greater part of those little cylinders above mentioned, and consequently the particles of the object they came from.

240. Upon the whole I conclude we may lengthen our telescopes at pleasure, according to the laws of the table, with good success; since not only the brightness and distinctness remain unaltered, but also the breadth of the pencils that enter the eye. Lastly to observe exceeding small stars and especially the Satellites of Jupiter and Saturn, the best way is to increase very much both the aperture and focal distance of the eye-glass. For since they appear like points even through the telescopes, there is nothing gained by endeavouring to increase their diameters; but their brightness must be increased as much as possible; and this is chiefly done by increasing the aperture. By doubling its breadth, the light received into it becomes quadruple, and then by doubling also the focal distance of the eye-glass, the distinctness returns to the same as at first<sup>a</sup>. But still the brightness will not become 16 times greater, according to cor. 2. prop. 3, but only 4 times; because as I said the picture of the star upon the retina is but a sensible point, whose brightness cannot

<sup>a</sup> Art. 227.



cannot therefore be increased by a diminution of its breadth, but only by an addition of new light. The case is different when we view the moon and primary planets through the same telescope, whose several parts receive 16 times more light than before. Thus by widening the apertures we very much increase the power of the telescope for finding out small stars and the satellites of Saturn, so that perhaps with a 30 foot glass, whose aperture is 6 inches or double the usual one, as much may be done as with another of 120 foot whose aperture by the table is also 6 inches. So far from *Hugens*.

## PROPOSITION VI.

241. *In reflecting telescopes of various lengths a given object will appear equally bright and equally distinct, when their linear apertures and also their linear amplifications are as the square-square roots of the cubes of their lengths: and consequently when the focal distances of their eye-glasses are also as the square-square roots of their lengths.*

Put  $A$  for the linear aperture of the reflecting concave,  $L$  for its focal distance or the length of the telescope,  $F$  for the focal distance of the eye-glass; and when the distinctness is given  $A^3$  is as  $FLL^2$ ; <sup>a</sup> Art. 233.

and when the brightness is given the amplification or  $\frac{L}{F}$  is as  $A^b$ , <sup>b</sup> Art. 229.

that is  $F$  is as  $\frac{L}{A}$ . Therefore when the distinctness and brightness

are both given,  $A^3$  is as  $\frac{L^3}{A}$ ; or  $A^4$  as  $L^3$ ; or  $A$  as  $\sqrt[4]{L^3}$ . The am-

plification  $\frac{L}{F}$  was as  $A$ , that is as  $\sqrt[4]{L^3}$ ; and therefore  $F$  is as

$\frac{\sqrt[4]{L^4}}{\sqrt[4]{L^3}}$  or  $\sqrt[4]{L}$ , Q. E. D.

242. In the reflecting telescope made and described by *John Hadley*, Esq; F. R. S. in the Philosophical Transactions N<sup>o</sup>. 376 and 378,  $L = 62\frac{1}{2}$  inches,  $F = \frac{1}{3}$  or  $\frac{2}{15}$  or  $\frac{1}{7.5}$  of an inch. For he uses 3 eye-glasses and as many apertures for the reflecter whose breadths are

$4\frac{1}{2}$ , 5,  $5\frac{1}{2}$  inches. Hence the linear amplifications or  $\frac{L}{F}$  are  $187\frac{1}{2}$ ,  $208\frac{1}{3}$ ,  $227\frac{2}{3}$  respectively. Taking the middle eye-glass and aperture for a standard I computed the following table for telescopes of other lengths by this Rule. Call the number of inches in the length of

of any telescope  $L$ , and the focal distance of its eye-glass will be equal to  $60 \sqrt[4]{10} L$  in thousandth parts of an inch. The quotient of  $L$

\* Art. 169. divided by  $60 \sqrt[4]{10} L$  or  $F$  gives the amplification<sup>a</sup>, which multiplied by 24 will always give the linear aperture in thousandth parts of an inch. For by the proposition  $\sqrt[4]{L}$  is as  $F$ ; that is  $\sqrt[4]{62\frac{1}{2}}$  or  $\sqrt[4]{\frac{125}{2}}$  or  $\sqrt[4]{\frac{625}{10}}$  or  $5 \sqrt[4]{\frac{1}{10}}$  is to  $\sqrt[4]{L}$  as  $\frac{3}{10}$  or 300 millesimals in the given eye-glass, to the millesimals in the correspondent eye-glass or in  $F = 60 \sqrt[4]{10} L$ . And the aperture being as the amplification by the proposition, say, as the amplification given or  $208\frac{1}{2}$  is to  $\frac{L}{F}$ , the amplification found, so is 5 inches, the aperture given, to the aperture sought  $= \frac{5}{208\frac{1}{2}} \times \frac{L}{F} = \frac{24}{1000} \times \frac{L}{F}$  inches.

243. Were it not for the unequal refrangibility of rays, refracting telescopes, though not so short as these<sup>b</sup>, would also be proportioned by this rule<sup>c</sup>: which not agreeing with experience, shews again that the aberrations arising from the spherical figure are inconsiderable in comparison to the other aberrations arising from the unequal refrangibility of the rays.

<sup>a</sup> Art. 221.

<sup>c</sup> Art. 209.

214.



Length of the telescope or focal dist. of the object-glass. Feet.	Linear aperture of the object-glass. Inch & Dec.	Focal dist. of the eye-glass. Inch & Dec.	Linear amplification, or magnifying power.	Length of the telescope or focal dist. of the concave Feet. $\frac{1}{2}$	Focal dist. of the eye-glass. Milles. Inch. 0. 167	Linear amplification, or magnifying power. 36	Linear aperture of the concave-metal. Milles. Inch. 0. 864
1	0. 55	0. 61	20	1	0. 199	60	1. 440
2	0. 77	0. 85	28	2	0. 236	102	2. 448
3	0. 95	1. 05	34	3	0. 261	138	3. 312
4	1. 09	1. 20	40	4	0. 281	171	4. 104
5	1. 23	1. 35	44	5	0. 297	202	4. 848
6	1. 34	1. 47	49	6	0. 311	232	5. 568
7	1. 45	1. 60	53	7	0. 323	260	6. 240
8	1. 55	1. 71	56	8	0. 334	287	6. 888
9	1. 64	1. 80	60	9	0. 344	314	7. 536
10	1. 73	1. 90	63	10	0. 353	340	8. 160
13	1. 97	2. 17	72	11	0. 362	365	8. 760
15	2. 12	2. 32	77	12	0. 367	390	9. 360
20	2. 45	2. 70	89	13	0. 377	414	9. 936
25	2. 74	3. 01	100	14	0. 384	437	10. 488
30	3. 00	3. 30	109	15	0. 391	460	11. 040
35	3. 24	3. 56	118	16	0. 397	483	11. 592
40	3. 46	3. 81	126	17	0. 403	506	12. 143
45	3. 67	4. 04	133				
50	3. 87	4. 26	141				
55	4. 06	4. 47	148				
60	4. 24	4. 66	154				
70	4. 58	5. 04	166				
80	4. 90	5. 39	178				
90	5. 20	5. 72	189				
100	5. 48	6. 03	199				
120	6. 00	6. 60	218				
140	6. 48	7. 13	235				
160	6. 93	7. 62	252				
180	7. 35	8. 09	267				
200	7. 75	8. 53	281				
220	8. 12	8. 93	295				
240	8. 48	8. 83	308				
260	8. 83	9. 71	321				
280	9. 16	10. 08	333				
300	9. 49	10. 44	345				
400	10. 95	12. 05	398				
500	12. 25	13. 47	445				
600	13. 42	14. 76	488				

245. These propositions, in *Hugens's* table for refracting telescopes, are measured by the Rheinland foot which is to the English foot as 139 to 135; so that taking their lengths of as many English feet, their apertures and eye-glasses and linear amplifications should be severally diminished in the subduplicate ratio of 139 to 135 by art. 234. that is nearly in the ratio of 139 to 137 or about  $\frac{1}{60}$  or  $\frac{1}{70}$  part of the whole.



## CHAP. XIII.

## CONCERNING THE RAIN-BOW.

## LEMMA I.

Fig. 189. 246. **T**HE ratio of the tangents,  $CT$ ,  $CV$ , of any two angles,  $CBD$ ,  $CBE$ , is compounded of the ratio of their sines,  $CD$ ,  $CE$ , taken directly, and of their cosines,  $BD$ ,  $BE$ , taken inversely.

For the right angled triangles  $BCT$ ,  $BDC$  are equiangular, and so are the right angled triangles  $BCV$ ,  $BEC$ . Therefore the ratio of  $CT$  to  $CV$ , which is compounded of  $CT$  to  $CB$  and of  $CB$  to  $CV$ , or of  $CD$  to  $DB$  and of  $EB$  to  $EC$ , is the same as the ratio of the rectangle under  $CD$ ,  $EB$  to the rectangle under  $DB$ ,  $EC$ , which  
 \* Euc. VI. 23. is compounded of the ratio of  $CD$  to  $CE$  and of  $EB$  to  $DB$ <sup>a</sup>, that is of the sines directly and cosines inversely. Q. E. D.

## LEMMA II.

Fig. 190, 191. 247. The least increment of an angle of incidence, is to the contemporary increment of the angle of refraction, as the tangent of the angle of incidence, to the tangent of the angle of refraction.

Let two rays  $AB$ ,  $aB$ , containing a very small angle  $ABa$ , be refracted at  $B$  along the lines  $BE$ ,  $Be$  by a plane or by any curve-surface. From any point  $C$ , of the line  $BC$  perpendicular to that surface, draw  $CDd$  cutting the incident rays (produced) at right angles<sup>b</sup> in  $D$  and  $d$ ; and likewise  $CEe$  cutting the refracted rays (produced) at right angles in  $E$  and  $e$ . Then because  $CD$  is to  $CE$  and  $Cd$  to  $Ce$  in the same ratio of the sines, disjointly we have  $Dd$  to  $Ee$  as  $CD$  to  $CE$ . Now the ratio of the small angles ( $ABa$  or)  $DBd$  and  $EBe$ , which are the contemporary increments or decrements of the angles of incidence and refraction, being compounded of the ratio of  $Dd$  to  $Ee$  and of  $BE$  to  $BD$ <sup>c</sup>, that is of  $CD$  to  $CE$  and of  $BE$  to  $BD$ , is the same as the ratio of the tangents of incidence and refraction<sup>d</sup>. Q. E. D.

\* Art. 68.  
 \* Art. 86.  
 \* Art. 246.

P R O-



## PROPOSITION I.

248. *When a ray of light is refracted into a circle, and successively reflected within it any given number of times before it emerges out of the circle by a second refraction; let the angle of refraction be multiplied by the number of successive reflections increased by an unite; and the excess of the resulting angle above the angle of incidence will be equal to half the angle contained under the incident and the emergent ray produced till they meet: that is, the excess abovementioned is equal to half that angle, under the incident and the emergent ray, in which the refracting circle lyes, when the number of reflections is odd; and is equal to half the other angle, under the same rays, which is the complement of the former to two right angles, when the number of reflections is even.*

For let  $ABCDE$  be a great circle of a sphere whose center is  $O$ , Fig. 192. to 195.  
and let an incident ray  $SA$  be refracted at  $A$  to  $B$ , and be reflected from  $B$  to  $C$ ; and at  $C$  let it either go out by refraction to  $G$ , or be reflected to  $D$ ; where let it either go out by refraction to  $H$ , or be reflected to  $E$ ; and so on. Art. 35, &c. And when the number of reflections is odd, a line  $OR$  drawn through the center  $O$  and the middlemost point of reflection, will bisect the angle at  $R$  under the incident and the emergent ray produced: because the reflections and refractions on each side of the line  $OR$  are equal in number and magnitude; the chords  $AB, BC, CD, DE$  described by the reflected ray being equal to one another. And for the same reason when the number of reflections is even, a line  $OT$ , drawn through the center  $O$  perpendicular to the chord that joins the two middlemost points of reflection, will bisect one of the angles at  $T$  under the incident and the emergent ray produced; and a line  $TV$ , perpendicular to  $TO$ , will bisect the other angle under them, which is the complement of the former to two right ones. Hence the line  $TV$  is parallel to the middlemost chord, because  $TO$  is perpendicular to them both. Draw a diameter  $POQ$  parallel to the incident ray  $SAM$ , and let it cut the reflected rays  $BC, CD, DE$  produced, in  $\beta, \gamma, \delta$ , respectively. Join  $OA, OB$  and in fig. 192. the sums of the three angles in each of the triangles  $OAB, OAR$ , are equal to one another; take away the common angle  $AOB$ , and the sum of the equal angles  $OAB, OBA$  in the first triangle, will be equal to the sum of the angles  $OAR, ORA$  in the second triangle. And by subtracting the angle of incidence  $OAR$  or  $OAM$  from both sums, we have  $2OAB - OAM = ORA = BOQ$ . Hence in fig. 193. the angle  $STV$  or  $P\beta C$ , being an external angle of the triangle  $OB\beta$ , equals  $OBC +$   

$N 2$ 
 $BOQ$



$BOQ = OAB + 2OAB - OAM = 3OAB - OAM$ . Hence again in fig. 194. the angle  $SRO$  or  $POC$ , being an external angle of the triangle  $OC\beta$ , equals  $OCB + P\beta C = OAB + 3OAB - OAM = 4OAB - OAM$ . Hence again in fig. 195. the angle  $STV$  or  $P\gamma D$ , being an internal angle of the triangle  $CO\gamma$ , equals  $OCD - CO\gamma = 5OAB - OAM$ , throwing away two right angles. For  $CO\gamma = 2$  right angles  $- POC = 2$  right angles  $- 4OAB + OAM$ . And so forward continually. Therefore if the number of successive reflections increased by an unite be called  $m$ , it appears that  $mOAB - OAM$  equals half the angle under the incident and emergent rays. *Q. E. D.*

## PROPOSITION II.

249. *Things remaining as they were, let the angle of incidence increase from nothing till it becomes a right angle; and the angle under the incident and the emergent ray, after any given number of reflections called  $n$ , will first increase and then decrease again; and will be the greatest of all when the tangent of the angle of incidence, is to the tangent of the angle of refraction, as  $n + 1$  to 1.*

Fig. 192. to  
195.  
\* Art. 248.

For putting  $m = n + 1$ , we had half the angle under an incident and the emergent ray equal to the excess of  $mOAB$  above  $OAM$ ; which excess, when the angles  $OAB$ ,  $OAM$  are very small, will also be but small; and will increase so long as the successive increments of  $mOAB$  shall exceed the contemporary increments of  $OAM$ ; and will decrease again when the successive increments of  $mOAB$  are exceeded by the increments of  $OAM$ ; and consequently will be the greatest of all when  $m$  times the least increment of  $OAB$  is equal to once the contemporary increment of  $OAM$ ; that is when the least increment of the angle of incidence  $OAM$  is to the contemporary increment of the angle of refraction  $OAB$ , and consequently the tangent of incidence is to the tangent of refraction<sup>b</sup>, as  $m$  to 1. *Q. E. D.*

\* Art. 247.

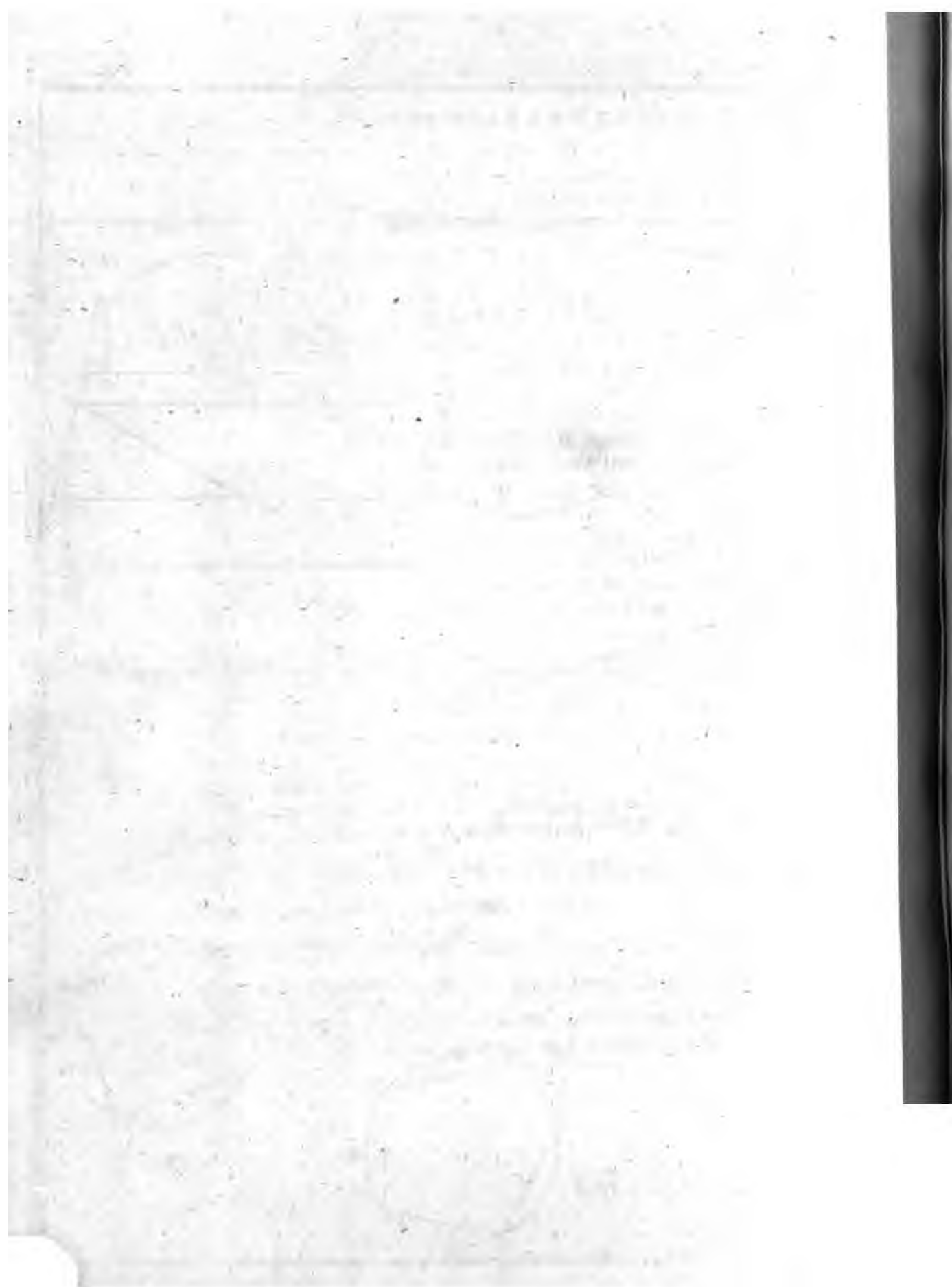
## PROPOSITION III.

250. *It is proposed to find two angles, whose sines shall be in a given ratio of  $I$  to  $R$ , and whose tangents shall be in another given ratio of  $m$  to 1.*

Fig. 196.

In any given line  $CEDA$ , let  $CA$  be to  $CD$  as  $I$  to  $R$ , and  $CA$  to  $CE$  as  $m$  to 1; with the center  $C$  and semidiameter  $CD$  describe an arch  $DB$ , cutting a circle  $ABE$  whose diameter is  $AE$ , in  $B$ ; draw  $ABF$ , and joining  $BC$ , the sine of the angle  $CBF$  will be to the sine





fine of  $CAF$  as  $I$  to  $R$ ; and the tangent of  $CBF$  to the tangent of  $CAF$  as  $m$  to  $1$ ; and consequently  $CBF$ ,  $CAF$  are the angles required. For in the triangle  $CAB$  the fine of the angle  $CBA$  or  $CBF$ , is to the fine of  $CAF$ , as  $CA$  to  $CB^a$  or  $CD$ , as  $I$  to  $R$  by <sup>a</sup> Art. 85. construction. Join  $BE$  and compleat the parallelogram  $CEBG$ ; and  $CG$  produced will cut  $ABF$  at right angles in  $F$ , because  $ABE$  is a right angle in the semicircle  $ABE$ . Therefore the lines  $FC$ ,  $FG$  are tangents of the angles  $CBF$ ,  $GBF$  or  $CAF$  to the radius  $BF$ ; and the tangent  $FC$  is to the tangent  $FG$  as  $FA$  to  $FB^b$  or <sup>b</sup> Euc. VI. 2. as  $CA$  to  $CE^b$  or as  $m$  to  $1$  by construction. Q. E. D.

251. *Corol. 1.* When parallel rays of the sun fall upon a spherical drop of rain, let the given ratio of  $I$  to  $R$  stand for the ratio of the fine of incidence to the fine of refraction; and let  $n$  be any given number of successive reflections made by every ray before it emerges out of the drop, and let  $m = n + 1$ ; and by these propositions it appears, that half the greatest angle which any of the emergent rays can make with the incident rays, is equal to  $m \times \text{ang. } CBF - CAF$ . For  $CBF$  and  $CAF$  or  $GBF$  are angles whose fines are as  $I$  to  $R$ , and whose tangents are as  $m$  to  $1$ ; and consequently are the angles of incidence and refraction of that ray, whose incident and emergent parts produced contain the greatest angle.

252. *Corol. 2.* The foregoing construction for determining the angle  $CBF$  is Dr. Halley's <sup>c</sup>, and Sir Isaac Newton's rule for calculating it, is this that follows. As  $\sqrt{mm - 1} \times RR$  is to  $\sqrt{II - RR}$ , so is the tabular radius to the cosine of the angle of incidence  $CBF$ . Whence this angle and its fine are given by the tables, and from thence by the ratio of  $I$  to  $R$  the tabular fine of the angle of refraction and the angle itself are also given. The rule may thus be demonstrated. We had  $CA : CB :: I : R$  and  $FA : FB :: m : 1$ .

Hence  $CA^2 = \frac{II}{RR} CB^2$ , and  $AF^2 = mm BF^2$ ; and so  $\frac{II}{RR} CB^2 - mm BF^2 = (CA^2 - AF^2 = FC^2 =) CB^2 - BF^2$ . Hence <sup>d</sup> Euc. I. 47.

$\frac{II}{RR} CB^2 - CB^2 = mm BF^2 - BF^2$ , and  $\sqrt{II - RR} \times CB^2 = \sqrt{mm - 1} \times RR \times BF^2$ ; and by resolving this equality into a proportion, and by extracting the roots, we have  $\sqrt{mm - 1} \times RR : \sqrt{II - RR} :: CB : BF :: \text{radius} : \text{cosine ang. } CBF$ .



## PROPOSITION IV.

*To explain the Phænomena of the Rain-bow.*

Design.

253. Having premised such mathematical principles as are necessary for an exact computation of the apparent diameters and breadths of the Rain-bows, I will here subjoin Sir *Isaac Newton's* entire explication of the colours of the bows, and of the manner in which they are formed; taking the liberty here and there of making a few additions to it; for the sake of such readers as may not be so skilful as those that he generally writes to.

Newt. Opt.  
P. 147.

254. This bow never appears but where it rains in the sun shine, and may be made artificially by spouting up water which may break aloft, and scatter into drops, and fall down like rain. For the sun shining upon these drops, certainly causes the bow to appear to a spectator standing in a due position to the rain and sun. And hence it is now agreed upon that this bow is made by refraction of the sun's light in drops of falling rain. This was understood by some of the ancients, and of late more fully discovered and explained by the famous *Antonius de Dominis* archbishop of *Spalato* in his book *de Radiis Visus & Lucis* published by his friend *Bartolus* at *Venice* in the year 1611, and written above 20 years before. For he teaches there how the interior bow is made in round drops of rain by two refractions of the sun's light, and one reflection between them; and the exterior bow by two refractions and two sorts of reflections between them in each drop of water; and proves his explications by experiments made with a phial full of water, and with globes of glass filled with water, and placed in the sun to make the colours of the two bows appear in them. The same explication *Des-Cartes* has pursued in his *Meteors* and mended that of the exterior bow. But while they understood not the true origin of colours, it is necessary to pursue it a little farther.

Fig. 197.

255. For understanding therefore how the bow is made, let a drop of rain or any other spherical transparent body be represented by the sphere *BNFG* described with the center *C* and semidiameter *CN*. And let *AN* be one of the sun's rays incident upon it at *N*, and be thence refracted to *F*, where let it either go out of the sphere by refraction toward *V*, or be reflected to *G*; and at *G* let it either go out by refraction to *R*, or be reflected to *H*; and at *H* let it go out by refraction towards *S*, cutting the incident ray in *Y*. Produce *AN* and *RG* till they meet in *X*, and upon *AX* and *NF* let fall the perpendiculars *CD* and *CE*, and produce *CD* till it falls upon



upon the circumference at  $L$ . Parallel to the incident ray  $AN$  draw the diameter  $BQ$ ; and let the sine of incidence out of air into water, be to the sine of refraction as  $I$  to  $R$ . Now if you suppose the point of incidence  $N$  to move from the point  $B$  continually till it comes to  $L$ , the arch  $QF$  will first increase and then decrease, and so will the angle  $AXR$  which the rays  $AN$  and  $GR$  contain; and the arch  $QF$  and angle  $AXR$  will be biggest when  $ND$  is to  $NC$  as  $\sqrt{II-RR}$  to  $\sqrt{3RR}$ , in which case  $NE$  will be to  $ND$  as  $2R$  to  $I^a$ . Also the angle  $ATS$ , which the rays  $AN$  and  $HS$  contain will first decrease and then increase; and grow least when  $ND$  is to  $NC$  as  $\sqrt{II-RR}$  to  $\sqrt{8RR}$ , in which case  $NE$  will be to  $ND$  as  $3R$  to  $I$ ; and so the angle which the next emergent ray (that is the emergent ray after three reflections) contains with the incident ray  $AN$  will come to its limit, when  $ND$  is to  $NC$  as  $\sqrt{II-RR}$  to  $\sqrt{15RR}$ ; in which case  $NE$  will be to  $ND$  as  $4R$  to  $I$ . And the angle which the ray next after that emergent, (that is the ray emergent after four reflections) contains with the incident, will come to its limit, when  $ND$  is to  $NC$  as  $\sqrt{II-RR}$  to  $\sqrt{24RR}$ ; in which case  $NE$  will be to  $ND$  as  $5R$  to  $I$ ; and so on infinitely, the numbers 3, 8, 15, 24, &c. being gathered by continual addition of the terms of the arithmetical progression 3, 5, 7, 9, &c.

256. Now it is to be observed, that as when the sun comes to his tropicks, days increase and decrease but a very little for a great while together; so when by increasing the distance  $CD$ , these angles come to their limits, they vary their quantity but very little for some time together; and therefore a far greater number of rays which fall upon all the points  $N$  in the quadrant  $BL$  shall emerge in the limits of these angles than in any other inclinations. Add to this, that of all the rays which fall upon the quadrant  $BL$ , those contiguous ones can only emerge parallel to one another, which emerge in the limits of these angles; and that all other contiguous rays emerge diverging from points either behind or before the drop; and consequently will fall much thinner upon the eye, at a great distance from the drop, than the parallel rays. For those rays which converge to points behind the eye, placed at a great distance from a small drop, are not sensibly different from parallel rays. This will appear by observing that while the arch  $BN$  is continually increasing from nothing, and the angle  $AXR$ , for example, is also increasing; the successively emergent rays, being continually less and less inclined to the incident ones or to the fixt line  $BQ$ , are also successively inclined in small angles to one another; and the same

Art. 252.

Art. 246.

249.

Fig. 198.

Fig. 197.



same thing is manifest while the angle  $AXR$  is decreasing; the successive rays being more and more inclined to  $BQ$ ; consequently in the limit between the increase and decrease of this angle the contiguous incident rays must emerge parallel to one another.

Fig. 199.

257. And farther it is to be observed, that the rays which differ in refrangibility will have different limits of their angles of emergence; and by consequence according to their different degrees of refrangibility, emerge most copiously in different angles; and being separated from one another appear each in their proper colours. Add to this that although the heterogeneous rays of any slender pencil whatever, as  $AN$ , will be separated by refractions at the drop into rays  $NFGR$  of one colour, and  $Nfgr$  of another, as by refractions through a prism; yet these emergent rays  $GR$ ,  $gr$  will not affect the eye with their distinct colours, unless they be in the limits of the angles  $AXR$   $Axr$ ; because every where within these greatest angles, an infinite number of such coloured pencils being variously inclined to one another are mixt together, and consequently appear white or without distinct colours. And the same may be said of the rays emerging any where within the greatest angle  $NYS$ . Fig. 197.

<sup>a</sup> Art. 31.

258. Now what these angles are may be gathered first by computing the angles of incidence and refraction by art. 252, and then the angles  $AXG$ ,  $ARS$  themselves by the 248th article. For in the least refrangible rays the sines  $I$  and  $R$  are 108 and 81<sup>a</sup>, and thence by computation the greatest angle  $AXR$  will be found 42 degrees and 2 minutes; and the least angle  $ARS$ , 50 degrees and 57 minutes. And in the most refrangible rays the sines  $I$  and  $R$  are 109 and 81, and thence by computation the greatest angle  $AXR$  will be found 40 degrees and 17 minutes, and the least angle  $ARS$ , 54 degrees and 7 minutes.

Fig 200.

259. Suppose now that  $O$  is the spectator's eye, and  $OP$  a line drawn parallel to the sun's rays. Let  $POE$ ,  $POF$ ,  $POG$ ,  $POH$  be angles of 40 degr. 17 min.; 42 deg. 2 min.; 50 deg. 57 min.; 54 deg. 7 min. respectively; and these angles turned about their common side  $OP$ , shall with their other sides  $OE$ ,  $OF$ , and  $OG$ ,  $OH$  describe the verges of two rain-bows  $AFBE$  and  $CHDG$ . For if  $E$ ,  $F$ ,  $G$ ,  $H$  be drops placed any where in the conical superficies described by  $OE$ ,  $OF$ ,  $OG$ ,  $OH$  and be illuminated by the sun's rays  $SE$ ,  $SF$ ,  $SG$ ,  $SH$ ; the angle  $SEO$  being equal to  $POE$  or 40 deg. 17 min. shall be the greatest angle in which the most refrangible rays can after one reflection be refracted to the eye; and therefore all the drops in the line  $OE$  shall send the most refrangible rays most copiously



copiously to the eye; and thereby strike the senses with the deepest violet colour in that region. In like manner the angle  $SFO$  being equal to the angle  $POF$  or 42 deg. 2 min. shall be the greatest in which the least refrangible rays after one reflection can emerge out of the drops; and therefore these rays shall come most copiously to the eye from the drops in the line  $OF$ , and strike the senses with the deepest red colour in that region. And by the same argument the rays which have intermediate degrees of refrangibility shall come most copiously from drops between  $E$  and  $F$  and strike the senses with the intermediate colours in the order which their degrees of refrangibility require; that is in the progress from  $E$  to  $F$  or from the inside of the bow to the outside in this order, violet, indigo, blue, green, yellow, orange, red. But the violet by the mixture of the white light of the clouds will appear faint and incline to purple. It may be farther observed, that all the rays but the violet in the line  $SE$  will emerge from  $E$  in a greater angle than  $SEO$  made by the violet, and consequently will pass below the eye; and that all the rays but the red in the line  $SF$  will emerge from  $F$  in a lesser angle than  $SFO$  made by the red, and consequently will pass above the eye; by which means only red will appear in the line  $SF$  and only violet in the line  $SE$ .

260. Again the angle  $SGO$  being equal to the angle  $POG$  or 50 deg. 57 min. shall be the least angle in which the least refrangible rays can after two reflections emerge out of the drops; and therefore the least refrangible rays shall come most copiously to the eye from the drops in the line  $OG$ , and strike the sense with the deepest red in that region. And the angle  $SHO$  being equal to the angle  $POH$  or 54 deg. 7 min. shall be the least angle in which the most refrangible rays after two reflections can emerge out of the drops; and therefore these rays shall come most copiously to the eye from the drops in the line  $OH$ , and strike the sense with the deepest violet in that region. And by the same argument the drops in the regions between  $G$  and  $H$  shall strike the sense with intermediate colours in the order which their degrees of refrangibility require; that is in the progress from  $G$  to  $H$ , or from the inside of the bow to the outside in this order; red, orange, yellow, green, blue, indigo, violet. And since these four lines  $OE$ ,  $OF$ ,  $OG$ ,  $OH$  may be situated any where in the abovementioned conical surfaces, what is said of the drops and colours in these lines is to be understood of the drops and colours every where in those surfaces.

261. Thus shall there be made two bows of colours, an interior and stronger by one reflection in the drops, and an exterior and



fainter by two; for the light becomes fainter by every reflection. And their colours shall lye in a contrary order to one another; the red of both bows bordering upon the space  $GF$ , which is between the bows. The apparent breadth of the interior bow  $EOF$ , measured cross the colours, shall be 1 deg. 45 min. and the breadth of the exterior,  $GOH$ , shall be 3 deg. 10 min. and the apparent distance between them,  $GOF$ , shall be 8 deg. 55 min. the greatest semidiameter of the innermost, that is, the angle  $POF$  being 42 deg. 2 min. and the least semidiameter of the outermost,  $POG$ , being 50 deg. 57 min.

Fig. 201.

262. These are the measures of the bows as they would be were the sun but a point, for by the breadth of his body the breadths of the bows will be increased and their distance decreased by half a degree. And so the breadth of the interior Iris will be 2 deg. 15 min. that of the exterior 3 deg. 40 min., their distance 8 deg. 25 min., the greatest semidiameter of the interior bow 42 deg. 17 min., and the least of the exterior 50 deg. 42 min. For let  $SEO$  be the limit of all the angles under the rays of any one colour, which coming from the center of the sun are reflected from the drop at  $E$  to the eye at  $O$ . In the ray  $SE$  take any point  $S$  at pleasure, and let the angles  $ESM$ ,  $ESN$  and also  $EOM$ ,  $EON$  be severally equal to a quarter of a degree, that is to half the apparent breadth of the sun. And joining  $OS$ , since the sums of the angles at the base  $OS$ , of the several triangles  $OSM$ ,  $OSE$ ,  $OSN$ , are equal among themselves, their vertical angles at  $M$ ,  $E$ ,  $N$  are also equal among themselves. Consequently the angle  $SMO$  will be the limit of all the angles contained under the incident and emergent rays of the same colour as before, which came from  $m$  the highest point of the sun; and  $SNO$  the limit of all the angles contained under the incident and emergent rays of the same colour as before, which came from  $n$  the lowest point of the sun. Therefore if all the rays of the sun were of the same colour, or alike refrangible, the apparent breadth of the bow, measured by the angle  $MON$ , would be but half a degree or equal to the apparent breadth of the sun measured by the angle  $MSN$  or  $mSn$ . But since his rays are differently refrangible, conceive the drop  $E$  to be placed any where in the inward or outward verges of the bows above described, upon supposition that the sun was but a point; and then it is manifest that the angle  $EOM$  must be added to the inside, and  $EON$  to the outside of the angles which the breadths of those bows subtend at  $O$ , to obtain their apparent breadths. A rain-bow is therefore a circular image of the sun reflected to the eye from the farther surfaces of innumerable drops of

of falling rain, and dilated in breadth by the unequal refrangibility of rays of different colours.

263. And such are the dimensions of the bows in the heavens found to be very nearly, when their colours appear strong and perfect. For once by such means as I then had I measured the greatest semidiameter of the interior iris about 42 degrees, the breadth of the red, yellow, green in that iris 63 or 64 minutes, besides the outmost faint red obscured by the brightness of the clouds, for which we may allow 3 or 4 minutes more. The breadth of the blue was about 40 minutes more besides the violet, which was so much obscured by the brightness of the clouds that I could not measure its breadth. But supposing the breadth of the blue and violet together to equal that of the red, yellow and green together; the whole breadth of this iris will be about  $2\frac{1}{4}$  degrees, as above. The least distance between this iris and the exterior iris was about 8 degrees and 30 minutes. The exterior iris was broader than the interior, but so faint, especially on the blue side, that I could not measure its breadth distinctly. At another time when both bows appeared more distinct I measured the breadth of the interior iris 2 deg. 10 min. and the breadth of the yellow and green in the exterior iris was to the breadth of the same colours in the interior as 3 to 2.

264. Whoever has a mind to repeat these observations after Sir *Isaac Newton* may observe, that the apparent semidiameter of the bow, (or of any ring of colours in either of the bows) is equal to the apparent altitude of its highest point added to the sun's altitude, and consequently may be measured by a common quadrant. For let  $SOP$  be the axis of the bows passing through the sun at  $S$  and the eye at  $O$ ,  $GOH$  an horizontal line,  $E$  the highest point of any ring of either of the bows, whose apparent semidiameter  $EOP$  is required. It is manifest that the angle  $EOP = EOH + HOP = EOH + SOG$ . Fig. 202.

265. This explication of the rain-bow is yet farther confirmed by the known experiment (made by *Antonius de Dominis* and *Des Cartes*) of hanging up any where in the sun-shine a glass-globe filled with water, and viewing it in such a posture that the rays which come from the globe to the eye may contain with the sun's rays an angle of either 42 or 50 degrees. For if the angle be about 42 or 43 degrees the spectator supposed at  $O$ , shall see a full red colour in that side of the globe which is opposed to the sun; as it is represented at  $F$ : and if that angle be made less, suppose by depressing the globe to  $E$ , there will appear other colours yellow, green and blue successively in the same side of the globe. But if the angle be Fig. 200.



made about 50 degrees, suppose by lifting up the globe at *G*, there will appear a red colour in that side of the globe which lyes toward the sun: and if the angle be made greater, suppose by lifting up the globe to *H*, the red will turn successively to other colours, yellow, green and blue. The same thing I have tried by letting a globe rest, and by raising or depressing the eye, or otherwise moving it to make the angle of a just magnitude. So far Sir *Isaac Newton*.

## LEMMA III.

266. *The tangent of the sum of two angles, is to the sum of their tangents, as the square of the radius, to the square of the radius diminished by the rectangle under the tangents: and the tangent of the difference of two angles, is to the difference of their tangents, as the square of the radius, to the square of the radius increased by the rectangle under the tangents.*

Fig. 203, 204.

Let *RA* and *RB* be tangents of two angles *ROA*, *ROB*. Then as *AB*, the sum or difference of the tangents, is to *AO*, the secant of either of the angles, so let *AO* be to *AC*, to be placed from *A* towards *B*. Again as *RC* is to *RO*, so let *RO* be to *RD*; and *RD* will be the tangent of the sum or difference of the two angles *ROA*, *ROB*. For joining *CO*, by the first of these proportions the triangles *AOB*, *ACO* will be equiangular<sup>a</sup>; and so the angle *AOB* is equal to *ACO*, or to *ROD* by the second proportion<sup>b</sup>. Hence in

<sup>a</sup> Euc. VI.6.

<sup>b</sup> Euc. VI.8.

fig. 203, because  $AC = \frac{AOq}{AB} = \frac{RAq + ROq}{RB + RA}$ , we have  $RC = (AC - AR) = \frac{RAq + ROq}{RB + RA} - RA = \frac{ROq - RB \times RA}{RB + RA}$ ; whence  $RD = \frac{RB + RA}{ROq - RB \times RA} \times ROq$ . By a like process fig. 204, we have  $AC = \frac{RAq + ROq}{RB - RA}$ ; whence  $RD = \frac{RB - RA}{ROq + RB \times RA} \times ROq$ . Q. E. D.

267. *Corol. 1.* Hence the tangent of the sum of any number of given angles, or the tangent of any multiple of a given angle, may be easily computed. Put  $RO = r$ ,  $RA = a$ ,  $RB = b$ , then the tangent of the sum of the angles whose tangents are *a* and *b*, that is  $RD = \frac{b + a}{rr - ab} \times rr$ ; call this tangent *x*; then for the same reason, the tangent of the sum of this last angle and of a third angle, whose



whose tangent is  $c$ , is  $\frac{x+c}{rr-xc} \times rr$  or (by substituting the value of

$x$ )  $\frac{rr \times a + b + c - abc}{rr - ab - ac - bc}$ , the tangent of the sum of three angles whose tangents are  $a, b, c$ ; and so on.

268. *Corol. 2.* Now put  $a = b = c$ ; and for the tangent of a double angle we have  $\frac{2a}{rr - aa} rr$ ; and for the tangent of a treble angle  $\frac{3arr - a^3}{rr - 3aa}$ ; and so on.

### PROPOSITION V.

269. *The apparent semidiameter of any rain-bow, or the greatest angle under an incident and emergent ray after any given number of successive reflections, being given; to find the ratio of refraction.*

Let  $m$  be the given number of successive reflections increased by one unite, and supposing the angles  $ABC, ABD$  to be the angles of incidence and refraction sought, let the angle  $ABE = m \times ABD$ , and the angle  $CBE$ , or  $m \times ABD - ABC$ , will be half the given angle under the incident and the emergent ray after  $m - 1$  reflections<sup>a</sup>. Put the common radius  $AB = r$ , the unknown tangent of<sup>a</sup> refraction  $AD = a$ , and the tangent of incidence  $AC = ma$ <sup>b</sup>, also  $AE = x$ , and  $t$  for the tangent of the given angle  $CBE$  answering to the radius  $r$ . Then by the lemma  $t : x - ma :: rr : rr + xma$ ; whence  $t = \frac{x - ma}{rr + xma} rr$ .

*Case 1.* In the first rain-bow  $m = 2$ ; whence  $t = \frac{x - 2a}{rr + 2xa} rr$ , and by art. 268,  $x = \frac{2a}{rr - aa} rr$  the tangent of  $2ABD$ . Substitute this value for  $x$  in the former equation and by reduction it becomes  $a^3 - \frac{3}{2}taa - \frac{1}{2}trr = 0$ . By resolving this equation the tangent  $a$  of the angle of refraction will be given, and the tangent of the angle of incidence  $AC = 2a$  by art. 251, whence the ratio of their sines is given by the tables.

*Case 2.* In the second rain-bow  $m = 3$ , whence  $t = \frac{x - 3a}{rr + 3xa} rr$ , and by art. 268,  $x = \frac{3arr - a^3}{rr - 3aa}$ , the tangent of  $3ABD$ . Substitute



stitute this value for  $x$  and you will find  $a^4 + \frac{2}{3} \frac{rr}{t} a^3 - 2rraa^* - \frac{1}{3} r^4 = 0$ ;

or putting  $T = \frac{rr}{t}$  the tangent of half the angle of

\* Art. 248. this bow<sup>2</sup>,  $a^4 + \frac{2}{3} T a^3 - 2rraa^* - \frac{1}{3} r^4 = 0$ . The same method serves for other bows to infinity.

270. *Corol.* In the first case putting  $T$  for  $2a$  or  $AC$  the tangent of the angle of incidence, and substituting  $\frac{1}{3} T$  for  $a$  in the former equation  $a^3 - \frac{2}{3} taa - \frac{1}{3} trr = 0$ , it is changed to this  $T^3 - 3tTT - 4rrt = 0$ , the same as Dr. *Halley's*, who proposed this problem as an expeditious method for finding the ratio of refraction in any fluid, by observing (when the sun is low and shines very bright) the angle under an incident and the emergent ray from a drop of any fluid hanging at the end of a capillary tube. See his examples *Phil. Trans.* N<sup>o</sup>. 267, and also the Rev<sup>d</sup>. Dr. *Morgan's* Dissertation upon the Rain-bow among the Notes upon *Robault's* Physics. P. 3. Ch. 17.

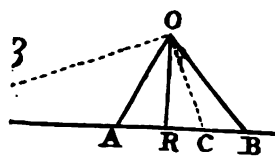
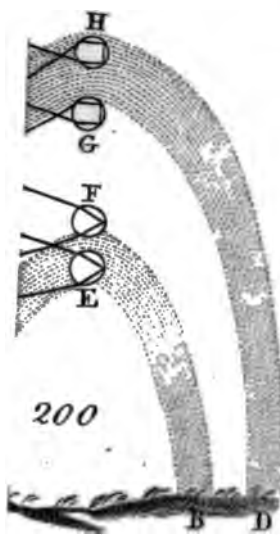
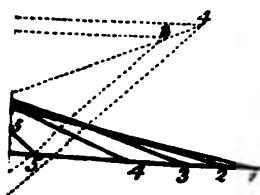
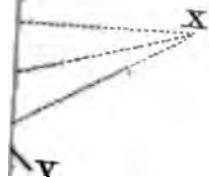
## CHAP. XIV.

### TELESCOPICAL DISCOVERIES IN THE FIXT STARS.

Multitudes  
of Telesco-  
pick Stars.

\* Art. 240. 271. **T**HAT the fixt stars have no sensible parallax, or, which is the same thing, that the earth's annual orbit (whose diameter a cannon ball could not describe in less than 50 years,) would appear of no sensible magnitude through a telescope placed at a fixt star, is such an amazing conclusion as could not be believed, were it not supported by undeniable evidence. But as this is the case, it is no longer a wonder why the best telescopes do not at all magnify the apparent diameters of the fixt stars, though they discover vast multitudes that are quite imperceptible to the naked eye; and the more of them as the aperture is more enlarged to take in more light, and the eye-glass made flatter to render it distinct\*. The Milky Way, which had puzzled the ancient philosophers for many ages, was found at last to be nothing else but a prodigious number of very minute stars, so close to one another that the naked eye can only perceive a whitish mixture of their faint lights. This was *Galileo's* discovery; who found also that those faint stars, which Astronomers call *Nebulosæ*, appeared through his telescope to be small clusters of very minute stars.

272. Hu-







272. *Hugenius* in the year 1656, looking by chance through a large telescope, at three small stars very close to one another in the middle of Orion's sword, saw several more as usual. But the three little stars very near one another (marked  $\theta$  by *Bayer*), together with four more, shone out as it were through a whitish cloud, much brighter than the ambient sky: which being very black and serene caused that lucid part to appear like an aperture, that gave a prospect into a brighter region. He viewed it many times, and found it continued in the very same place, and of the same shape as the figure represents, and called it *Portentum cui certe simile aliud nusquam apud reliquas fixas potuit animadvertere*<sup>a</sup>.

Lucid spots  
among the  
fixt Stars.

Fig. 206.

273. But in the Philosophical Transactions<sup>b</sup>, there is an account of a later discovery of five more such lucid spots, though less considerable than this of *Hugenius*; the middle of which, we are there told, is at present in  $\Pi. 19^{\circ}. 00'$ . with south latitude  $28^{\circ}. 45'$ ; and that it sends forth a radiant beam into the south east, as another in the girdle of *Andromeda* seems to do into the north east. It is also there remarked, that "though these spots are in appearance but small, and most of them but a few minutes in diameter, yet since they are among the fixt stars, as having no annual parallax, they cannot fail to occupy spaces immensely great, and perhaps not less than our whole solar system; in all which spaces it should seem, that there is a perpetual uninterrupted day."

<sup>a</sup> Systema  
Saturnium  
p. 8.

<sup>b</sup> N. 347.  
Jones's abr.  
Vol. 4. p. 224.

274. It is to the Author of these reflections, if I mistake not, that we owe another curious account of what is principally remarkable in the new stars that have appeared and disappeared for 150 years last past<sup>c</sup>. I will mention but one or two. "That in the chair of *Cassiopeia* was not seen by *Cornelius Gemma* on the eighth of November 1572, who says, he that night considered that part of the Heaven in a very serene sky, and saw it not: but that the next night, Novemb. 9; it appeared with a splendour exceeding all the fixt stars and scarce less bright than Venus. This was not seen by *Tycho Brahe* before the 11th of the same month; but from thence he assures us, that it gradually decreased and died away, so as in March 1574, after 16 months, to be no longer visible; and at this day not the least signs of it remain. The place thereof in the sphere of the fixt stars, by the accurate observations of the same *Tycho*, was  $0^{\circ}. 9^{\circ}. 17'$ . from the first star of Aries, with  $53^{\circ}. 45'$ . north latitude." To this account Sir *Isaac Newton* adds<sup>d</sup>, that in November, when it first appeared, it seemed equal to Venus in brightness, in December to Jupiter, in January 1573 less than Jupiter, but bigger than Sirius, and equal to him in February and March; in April and May equal to the

New Stars.

<sup>c</sup> Phil. Transf.  
N. 346.  
Jones's abr.  
Vol. 1. p. 222.

<sup>d</sup> Philos.  
Princip.  
p. 526.



the stars of the second magnitude, in June, July and August to those of the third, in September, October and November to those of the fourth, in December and January 1574 to those of the fifth, in February to those of the sixth, and in March it vanished. That its colour was at first clear, white and splendid, afterwards yellow, and in March 1573 red and fiery like Mars or Aldebaran, in May of a pale livid colour like that of Saturn, which grew fainter and fainter till it vanished.

275. " That such another star was seen and observed by the scholars of *Kepler* to begin to appear on Sept. 30. *St. Vet. Anno* 1604, which was not to be seen the day before; but it broke out at once with a lustre greater than that of Jupiter; and like the former, died away gradually, and in much about the same time disappeared totally, there remaining no footsteps thereof in January 1605. This star was near the ecliptick, following the right leg of *Serpentarius*; and by the observations of *Kepler* and others, was in  $7^{\circ} 20' 00''$  from the first star of *Aries*, with north latitude  $1^{\circ} 56'$ .

276. Lastly, that the sudden eruption of such another star, shining out more than usual, engaged *Hipparchus* to make the first catalogue of the fixt stars; that posterity might know what changes might happen among them.

The Origine  
of new Stars.

277. How minute soever the particles of light may be, the perpetual emission of them from the body of the sun must have caused, before this time, a sensible diminution in his magnitude, without some supply of new matter. But since a diminution of the sun's diameter has not yet been discovered by the most accurate observations, Sir *Isaac Newton* therefore imagined, that those comets which approach so near the sun as to pass through his atmosphere, may be so much resisted and retarded after several revolutions, as at last to fall upon the sun and so become a mean of keeping his magnitude nearly the same. And this opinion led him farther to conjecture, that the stars we have mentioned, which suddenly shine out with very great splendor and then decay gradually till they vanish out of sight, may now and then be stirred up and blaze out again by the shock of a comet falling down upon them. But those other new stars, which appear and disappear periodically, which increase by very slow degrees, and seldom exceed the stars of the third magnitude (several of which may be seen in the history I mentioned) he takes to be of another sort, or at least in another state; which revolving about their axes, like the sun, may expose their light and dark parts to us successively. For the fixt stars are undoubtedly self-shining bodies of the same kind as the sun, and therefore equally subject to large dark



dark spots or crusts upon their surfaces. Because the light of the sun propagated to those vast distances, and reflected back from opake bodies of no sensible apparent magnitudes, would be too much rarified to affect our senses; as *Galileo* collected from the fainter lights of the remoter planets from the sun, compared to the lustre of the fixt stars.

278. After several attempts by Dr. *Hook*, Mr. *Flamsteed*<sup>a</sup>, and others<sup>b</sup>, to determine the annual parallax of the fixt stars, the honourable *Samuel Molyneux* Esq; in the year 1725, erected at *Kew* a very accurate instrument, in order if possible to arrive at some degree of certainty in this difficult enquiry: in the prosecution of which he followed Dr. *Hook* in some respects, as in taking the zenith distances of the brightest star in the Dragon's head at the times of its transits over the meridian, and also in the form of his instrument, constructed almost upon the same principles with the Doctor's, but executed to a degree of exactness vastly greater, and chiefly owing to the care and contrivance of Mr. *George Graham*.

An enquiry into the annual parallax of the fixt stars. *Wallis's opera*, Vol. 3. p. 701. *Phil. Transf.* N. 364. *Abridg.* Vol. 6. p. 165.

279. The Rev. Mr. *Bradley*, Professor of Astronomy at Oxford, who all along assisted Mr. *Molyneux* in the prosecution of this noble design, has obliged the publick with a very accurate history of it, in a letter to D. *Halley*<sup>c</sup>; containing not only an account of several new and surprizing phænomena that attended the observations, (which he therefore continued and repeated after Mr. *Molyneux*'s decease,) but also a compleat discovery of the true cause of them; which at last enabled him to settle the point in question, and to draw from it some admirable consequences relating to the propagation of light. As I look upon these discoveries to be some of the finest that we have had since the invention of telescopes, I will endeavour to give the substance of them in as clear a manner as I can.

*Phil. Transf.* N. 406. *Abridg.* Vol. 6. p. 167.

280. The result of the observations upon the bright star in the dragon's head, marked  $\gamma$  by *Bayer*, was this;

Some account of the observations and phænomena.

Beginning from December 3, 1725, its distance from the zenith being taken several days, at the time of its transit over the meridian, there appeared no material difference in the observations.

281. On Decem. 17, it passed a little more southerly from the zenith than before, and still more on the 20th; which was matter of surprize, both because no sensible alteration of parallax could so soon be expected in this star at that time of the year, and because it was the contrary way to what it would have been, had it proceeded from an annual parallax.



282. About the beginning of March 1726, the star was found to be 20" more southerly than at the time of the first observation, and seemed to have arrived at its utmost limit southwards.

283. By the middle of April it appeared to be returning back again towards the north, and about the beginning of June it passed at the same distance from the zenith as it had done in December when it was first observed.

284. From that time it continued to move northwards till September following, when it again became stationary, being then near 20" more northerly than in June, and no less than 39" more northerly than it was in March.

285. From September it returned towards the south, till it arrived in December at the same situation it was in at that time twelve months, allowing for the difference of declination on account of the precession of the equinox.

286. By the like observations made upon a small star almost opposite in right ascension to  $\gamma$  *Draconis*, and at about the same distance from the north pole of the equator, it appeared to change its declination 19", that is about half as much as  $\gamma$  *Draconis* did in the same time. Which plainly proved, as Mr. *Bradley* observes, that these apparent changes were not owing to a nutation of the earth's axis, since the changes on this account would have been nearly equal in these stars, as lying near the solstitial colure.

287. Upon comparing the observations with each other it was discovered in both these stars, that the apparent difference of declination, reckoned from the limits above mentioned, was always nearly proportionable to the versed sine of the sun's distance from the equinoctial points.

288. And that the whole difference of declination in these stars, was as the sine of the latitude of each respectively.

Mr. *Bradley's*  
hypothesis to  
solve these  
phenomena.

289. After a year's observations upon many other stars, in different parts of the heavens, made with a new instrument set up at *Wansled* in 1727, Mr. *Bradley* found out some other properties of their apparent motions; and after examining and rejecting two or three hypotheses, by which he attempted to solve them, at last he conjectured that all these phenomena proceeded from the progressive motion of light and the earth's annual motion in her orbit. For he perceived, that if light was propagated in time, the apparent place of a fixt object would not be the same when the eye is at rest, as when it is moving in any other direction than that of a line passing through the eye and object; and that when the eye is moving



ing in different directions, the apparent place of the object would be different. I will first deduce some consequences from this hypothesis and then compare them with the phænomena.

290. If an eye moves uniformly in a straight line from  $a$  to  $b$  in the time that the light of a fixt star descends uniformly in a straight line from  $c$  to  $b$ , the star will appear in a direction constantly parallel to  $ac$ . Some consequences from the hypothesis. Fig. 207.

For conceiving the eye to carry the line  $ac$  parallel to itself, its intersection with the fixt line  $bc$  will move uniformly<sup>a</sup> from  $c$  to  $b$ , and will therefore accompany a particle of light descending uniformly from  $c$  to  $b$ ; and because this intersection is a moving point, not only in the fixt line  $bc$ , but also in the moving line  $ac$ , it is plain that the particle which accompanies the intersection  $c$ , moves relatively in the moving line  $ac$ . In like manner a particle of every other ray, parallel to  $cb$ , which the moving line  $ac$  successively meets with, moves also in the moving line  $ac$ ; and thus a succession of these particles, moving along  $ac$ , constitute a visual ray in whose direction the star appears. <sup>a</sup> Euc. VI. 2.

291. Hence supposing the earth's center  $b$  to move uniformly in a circular orbit  $r \propto Ab$ , round the sun in its center  $B$ ; if a line  $BC$  drawn towards a fixt star, supposed infinitely distant, you take a distance  $BC$  in proportion to  $Bb$  or  $BA$ , as the velocity of light to the velocity of the earth's center, an observer upon the earth at  $b$ , will constantly see that star in a direction very nearly parallel to a line  $AC$ , connecting the point  $C$  with a point  $A$  in the orbit constantly 90 degrees behind the earth. Fig. 208.

For drawing  $ba$  and  $bc$  parallel to  $BA$  and  $BC$  respectively, and tending the same ways from  $b$  and  $B$ , and also any line  $ac$  parallel to  $AC$ ; by the similar triangles  $bca$ ,  $BCA$ , we have  $bc : ba :: BC : BA$ , as the velocity of light to the velocity of the earth. Consequently if an eye be supposed to move along the tangent  $ab$  with this latter velocity, it will see the star in a direction constantly parallel to  $ac$ <sup>b</sup> or  $AC$ . But the eye in the orbit moves with that velocity, and passes by the point  $b$  in the direction of that tangent, and therefore at that passage it saw the star in the same direction in which the other eye in the tangent sees it constantly. The earth's diurnal motion alters this conclusion so little that I need not here consider it. <sup>b</sup> Art. 290.

292. Therefore the apparent parallax of the star to an observer at  $b$ , is constantly measured by the angle  $ACB$ , if the point  $A$  be always 90 degrees behind the earth at  $b$ , and consequently 90 degrees before the sun's apparent place  $\odot$  in the ecliptick.



293. Hence the apparent latitude of any star, supposed infinitely distant, will be least of all when the sun's place in the ecliptick is 90 degrees forwarder than the star's; and from that time it will increase for half a year, and then decrease for the next half year; and its increment reckoned from these limits will be constantly as the versed sine of the sun's longitude reckoned from his place before mentioned.

Fig. 209.

For drawing  $CD$  perpendicular to the plane of the orbit, join  $AD$ , and draw  $DB$  cutting the orbit in  $L$  and  $M$ . The point  $L$ , nearest to  $D$ , is the star's place in the ecliptick, and the point  $\odot$ , opposite to  $b$ , is the sun's place therein. Now when the point  $\odot$  was at  $N$ , 90 degrees forwarder than  $L$ , the point  $A$ , being always 90 degrees forwarder than  $\odot^a$ , was at  $M$ , the farthest point from the perpendicular  $CD$ ; and consequently the star's apparent latitude, always measured by the angle  $CAD$ , was then the least possible.

<sup>a</sup> Art. 292.

Draw  $AE$  perpendicular to  $LM$ , and joining  $CL$ ,  $CE$ ,  $CM$ , draw  $MF$  perpendicular to  $CE$  produced. Then conceiving the point  $A$  to move in the perpendicular  $AE$  towards  $E$ , the angle  $CAD$  will approach to a *maximum*  $CED$ , and therefore will increase very little, especially as the angular approach  $ACE$  is exceeding small. Therefore instead of the apparent latitude  $CAD$  we may take  $CED$ , and consequently the angle  $ECM$  for the increment of the least latitude  $CMD$ : now this small angle  $ECM$  is as its sine  $MF$  or (because the ratio of  $MF$  to  $ME$  varies very little) as  $ME$  the versed sine of the arch  $MA$  equal to  $N\odot$ , the sun's longitude from  $N$ , 90 degrees forwarder than the star's place  $L$ .

294. When a star is situated any where in the solstitial colure, the increments or decrements of its latitude and declination are the very same quantities; and therefore if the star be supposed infinitely distant, and its longitude be in the beginning of Capricorn, with north declination, its apparent declination will be the least at the time of the vernal equinox, and the greatest at the autumnal; and its increments and decrements reckoned from these limits, will be proportionable to the versed sine of the sun's longitude reckoned from the equinoctial points: which agrees with the phenomena in Art. 287.

Fig. 210.

295. The whole apparent parallax  $LP M$  of a star in the pole of the ecliptick, is to the whole apparent parallax  $LC M$  in the latitude of any other star, as the radius to the sine of the latitude  $CBL$ . For since  $BP$  equals  $BC$ , drawing  $BF$  and  $BG$  perpendiculars to  $CL$  and  $CM$ , the small angle  $BPL$  is to  $BCL$  as  $BL$  to  $BF$ , that is, as the radius to the sine of the angle  $BLF$  or of the latitude  $CBL^b$ . Again, the small angle  $BPM$  is to  $BCM$  as  $BM$  to  $BG$ , that is, as the

<sup>b</sup> Art. 68.



the radius to the sine of  $BMG$  or of the latitude  $CBL$ , as before. Therefore the whole angle  $LPM$  is to  $LCM$  as the radius to the sine of the latitude  $CBL$ .

Hence, from the observed parallax in latitude, or in declination of such stars as lye in or near the solstitial colure, we have the parallax that would belong to a star in the pole of the ecliptick, which is plainly the greatest of all. Thus in  $\gamma$  *Draconis* whose latitude  $CBL = 74^\circ. 58'. 20''$ , the observed parallax  $LCM$  was  $39''^a$ , and thence the greatest parallax  $LPM$  comes out  $40''.4$ . Likewise in the little star above mentioned, whose distance from the north pole of the equator is  $38^\circ. 28'. 35''^b$ , and consequently its latitude a little above  $28^\circ. 02'. 25''$ , as being almost opposite in right ascension to  $\gamma$  *Draconis*, the observed parallactick angle  $LCM$  was  $19''^c$ , and thence  $LPM$  comes out  $40''.4$  as before.

<sup>a</sup> Art. 284.

<sup>b</sup> The 36th Star *Camelopardis* Hevelii in *Flamsteed's* Catalogue.

<sup>c</sup> Art. 286.

296. Mr. *Bradley* having applied his observations upon the parallax in declination of stars in any situation whatever, to his theory farther pursued, assures us they all conspire to prove, that the greatest parallax  $LPM$  is about 40 or 41 seconds, and thinks the *medium*  $40''.\frac{1}{2}$  cannot differ so much as one second from the truth.

The greatest apparent parallax.

297. Hence the velocity of star-light comes out 10210 greater than the velocity of the earth's mean motion round the sun. For the former velocity is to the latter as  $BP$  to  $BL$  or  $BM^d$ , that is, as the radius to the tangent of  $BPL$  or  $BPM = 20''.\frac{1}{2}$  as above determined.

The velocity of star-light.

<sup>d</sup> Art. 290.

298. From what has been said Mr. *Bradley* infers, 1. That the lights of all those stars arrive at the earth with equal velocities. 2. That unless their distances from us are all equal, (which for other reasons besides that of their different lustre, is highly improbable) their lights are propagated uniformly to all distances from them. 3. That the velocity of star-light is such as carries it through a space equal to the sun's distance from us in  $8'. 13''$ , (this time being to the time in which the earth might describe that distance, with the velocity of her mean motion round the sun, as 1 to 10210, and this latter time, to half a year, as the diameter of a circle to its circumference.) 4. That the time so determined can scarce differ 5 or 10 seconds from the truth, which is such a degree of exactness as can never be expected from the eclipses of Jupiter's satellites. 5. That as this determination of the velocity of star-light, comes out a *medium* among several determinations of the velocity of the sun's light reflected from those satellites, we may reasonably conclude that the velocities of these lights are equal. And lastly, since it is highly probable that the velocity of the sun's emitted light is also equal to that of star-light,

Some properties of light...

light,



- \* Art. 277. light<sup>a</sup>, it is equally probable that its velocity is not altered by reflection into the same medium.

299. From Art. 291, 295, &c. it follows plainly, that a star placed in the pole of the ecliptick would appear in a year's time, to describe about the pole a little circle whose apparent semidiameter is  $20''\frac{1}{4}$ ; and that any other star will appear to describe, about its true place, an ellipsis whose long axis is at right angles to the circle of longitude passing through the star's true place, and equal to the diameter of the little circle just mentioned, and whose short axis is to the long one, as the sine of the star's latitude to the radius<sup>\*</sup>.

Real parallax  
of the stars  
insensible.

300. Upon this theory farther pursued Mr. *Bradley* proceeds synthetically, by assuming the *maximum* of apparent parallax as determined above, and calculating tables of the differences in declination of  $\gamma$  *Draconis* situated near the solstitial colure, and of  $\gamma$  *Ursæ Majoris* nearer to the equinoctial than the solstitial colure; and by com-

\* The following elegant proposition is reprinted from T. Simson's Mathematical Essays, as it affords an easy solution of the consequences mentioned in this article, and serves to confirm other parts of this Theory.

PROPOSITION. To find the path which the progressive motion of light and the motion of the Earth in its orbit make a Star appear to describe in one entire annual revolution of the Earth.

Fig. 211.

Let  $ATBA$  be the orbit of the earth;  $S$  the sun in one focus;  $F$  the other focus;  $T$  the earth moving in its orbit from  $A$  towards  $B$ ;  $DTn$  a tangent at  $T$ ; and  $SD$ ,  $FE$  perpendiculars thereto: Let  $\mathcal{Q}mKR\mathcal{Q}$  be part of an indefinite plane parallel to that of the ecliptick, passing through  $R$  the center of the given star; and take  $Tn$  to  $TR$ , as the velocity of the earth in its orbit at  $T$  to that of a particle of light coming from the said star: Let  $Tm$  be parallel to  $nR$ ;  $PnV$  perpendicular to  $AB$ ; and  $\mathcal{Q}RK$  parallel to  $PnV$ : Then from the 290th article it is manifest, that a ray of light, coming from  $R$  to the earth at  $T$ , will appear as if it proceeded from  $m$ , where the line  $Tm$  produced, intersects the said parallel plane, and therefore because  $Tm$  is parallel to  $Rn$ , and any parallelogram, intersecting two parallel planes, cuts them alike in every respect, it is evident that  $Rm$  must be equal to  $Tn$ , and  $\mathcal{Q}Rm$  to  $VnD$ ; wherefore since  $D$  and  $P$  are equal to two right angles,  $DSP$  and  $DnP$  must be equal, also, to two right angles<sup>a</sup>, and consequently  $\mathcal{Q}Rm (= VnD^b) = DSP = AFE$ . But  $Tn$  or  $Rm$ , expressing the celerity of the earth at  $T$ , is known to be inversely as  $SD^c$ ; or because  $SD \times FE$  is every where the same<sup>d</sup>, directly as  $FE$ ; wherefore the angles  $AFE$ ,  $\mathcal{Q}Rm$  being every where equal, and  $Rm$  in a constant proportion to  $FE$ , the curve  $\mathcal{Q}mK$  described by  $m$ , the apparent place of the star in the said parallel plane, will, it is manifest, be similar in all respects to  $AEB$  described by the point  $E$ : But this curve is known to be a circle<sup>e</sup>; therefore  $\mathcal{Q}mK$  must likewise be a circle, whose diameter  $\mathcal{Q}RK$  is divided by  $R$ , the true place of the star, in the same proportion as the transverse axis of the earth's orbit is divided by either of its foci. Wherefore, sofar as a small part of the circumjacent heavens may, in this case, be considered as a plane passing perpendicular to a line joining the eye and star, it follows from the principles of orthographic projection, that the star will be seen in the heavens as describing an ellipsis, whose center (as the excentricity of the orbit is but small) nearly coincides with the true place of the star, except the said place be in the pole or plane of the ecliptick; in the former of which cases the star will appear to describe a circle, and in the latter an arch of a great circle of the sphere, which by reason of its smallness may be considered as a right line.

<sup>a</sup> Eucl. I. 32.

Cor. 1.

<sup>b</sup> Eucl. XI. 10.

<sup>c</sup> Newt. prin-

cip. I. 1. Cor.

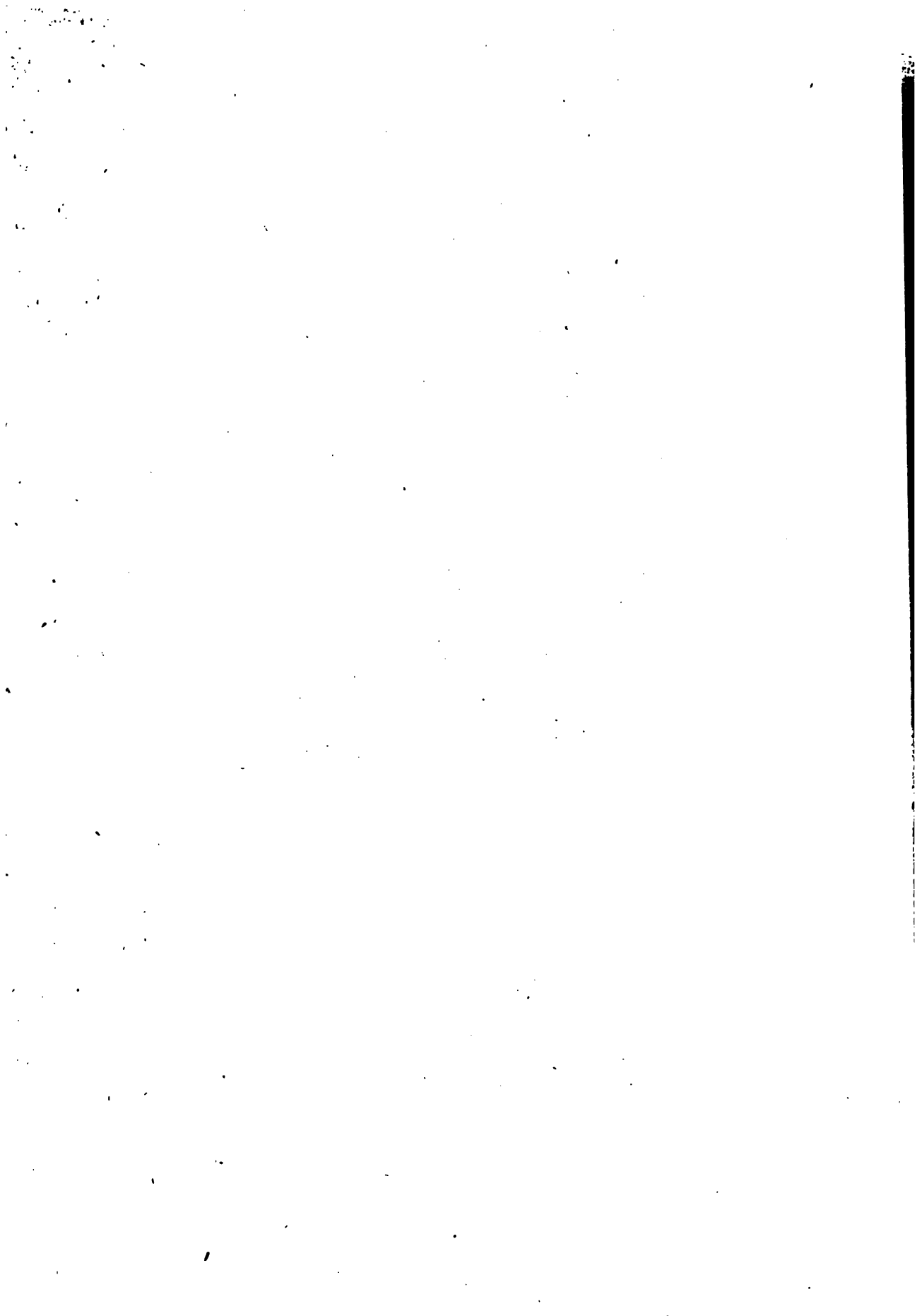
<sup>d</sup> Hamilton's

Conics, II. 21.

<sup>e</sup> Hamilton's

Conics, II. 20.

paring





paring the tables with his observations, he found they agreed together throughout the year, sometimes in the very same number of seconds, and that in 50 or 60 observations of each star, they never differed so much as two seconds; allowing for the variation of declination caused by the regression of the equinoctial points: which amounts to a physical demonstration of the truth of his theory, and does in consequence afford a very satisfactory answer to the point in question, concerning the real parallax and distance of the fixt stars. As to which he believes he may venture to say, that the real parallax in either of the stars abovementioned does not amount to 2", being of opinion that if it were 1" he should have perceived it in the great number of observations that he made especially upon  $\gamma$  *Draconis*; which agreeing with the theory, without allowing any thing for a real parallax<sup>a</sup>, nearly as when the sun was in conjunction with,<sup>a</sup> *Art. 290,* as in opposition to this star, it seemed to him very probable, that its<sup>291, &c.</sup> real parallax is not so great as one single second; and consequently that it is above 400000 times farther from us than the sun.

F I N I S.

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